

## Exercices

①  $\square A^r(x) = J^r(x) = \int dy \delta(x-y) J^r(y)$

$$A^r(x) = \int dy G_{ret}(x-y) J^r(y)$$

$$\Rightarrow J^r(x) = \int dy \square_x G_{ret}(x-y) J^r(y) \Rightarrow \square_x G_{ret}(x-y) = \delta(x-y)$$

$$G_{ret}(x-y) = \frac{1}{(2\pi)^4} \int dk G_{ret}(k) e^{-ik(x-y)}$$

$$\Rightarrow \square_x G_{ret}(x-y) = \frac{1}{(2\pi)^4} \int dk (-i)^2 k^2 G_{ret}(k) e^{-ik(x-y)} = \frac{1}{(2\pi)^4} \int dk e^{-ik(x-y)}$$

$$\Rightarrow -k^2 G_{ret}(k^2) = 1 \Rightarrow G_{ret}(k) = \frac{-1}{k^2}$$

$$G_{ret}(x-y) = \frac{1}{(2\pi)^4} \int dk G_{ret}(k) e^{-ik(x-y)} = \frac{1}{(2\pi)^4} \int dk \frac{-1}{k^2} e^{-ik(x-y)}$$

$$= \int \frac{dk}{(2\pi)^3} e^{+ik(\bar{x}-\bar{y})} \underbrace{\int \frac{dk^0}{(-2\pi)} \frac{e^{-ik^0(x^0-y^0)}}{(k^0)^2 - k^2}}_{\sim I \text{ par théorème des résidus...}}$$

$$I = \int_{-\infty}^{+\infty} dz \frac{e^{-iz(x^0-y^0)}}{z^2 - k^2} = \int_{-\infty}^{+\infty} dz \underbrace{\frac{e^{-iz(x^0-y^0)}}{(z-k_+)(z-k_-)}}_{F(z)}$$

$$I = -2\pi i \left[ \text{Res}(F, z_+) + \text{Res}(F, z_-) \right]$$

$$= -2\pi i \left[ \lim_{z \rightarrow z_+} (z-z_+) F(z) + \lim_{z \rightarrow z_-} (z-z_-) F(z) \right] =$$

$$= -2\pi i \left[ \lim_{z \rightarrow z_+} \frac{e^{-iz(x^0-y^0)}}{z-z_-} + \lim_{z \rightarrow z_-} \frac{e^{-iz(x^0-y^0)}}{z-z_+} \right] =$$

$$= -2\pi i \left[ \frac{e^{-i(k-i\varepsilon)(x^0-y^0)}}{2k} + \frac{e^{-i(-k-i\varepsilon)(x^0-y^0)}}{-2k} \right] \xrightarrow{\varepsilon \rightarrow 0} I = -2\pi i \left[ \frac{e^{-ik(x^0-y^0)}}{2k} - \frac{e^{ik(x^0-y^0)}}{2k} \right]$$

$$\begin{aligned}
\Rightarrow G_{xx}(x-y) &= \int \frac{dk}{(2\pi)^3} e^{+ik(x-y)} \cdot \frac{(-2\pi i)}{(-2\pi)} \left[ \frac{e^{-ik(x-y)}}{2k} - \frac{e^{+ik(x-y)}}{2k} \right] \Theta(x-y) = \\
&= \frac{-i}{(2\pi)^3} \int \frac{dk}{2k} e^{+ikr} [e^{+ikr} - e^{-ikr}] \Theta(r) \\
&= \frac{-i}{(2\pi)^3} \int \frac{|k|^2 d|k| - d\cos\theta d\phi}{2k^3} e^{+i|k|\cdot|r| \cos\theta} [e^{+ikr} - e^{-ikr}] \Theta(r) = \\
&= \frac{-i}{(2\pi)^3} \int \frac{d|k|}{2} |k| (2\pi) [e^{+ikr} - e^{-ikr}] \Theta(r) \int_0^1 d\cos\theta e^{+i|k|\cdot|r| \cos\theta} = \\
&= \frac{-i}{2(2\pi)^3} \int d|k| |k| [e^{+ikr} - e^{-ikr}] \Theta(r) - \frac{e^{-i|k||r|} + e^{+i|k||r|}}{i|k||r|} = \\
&= \frac{1}{2(2\pi)^2 |r|} \int dk^0 [e^{+ik^0 r} - e^{-ik^0 r}] [e^{+i|k||r|} - e^{-i|k||r|}] \Theta(r) = \\
&= \frac{1}{2(2\pi)^2 |r|} \int_0^{+\infty} dk^0 \left[ e^{+ik^0(r^0 + |r|)} - e^{+ik^0(r^0 - |r|)} - e^{+ik^0(-r^0 - |r|)} + e^{+ik^0(-r^0 + |r|)} \right] \Theta(r) = \\
&= \frac{1}{2(2\pi)^2 |r|} \int_0^{\infty} dk^0 \left[ e^{+ik^0(r^0 + |r|)} + e^{+ik^0(-r^0 + |r|)} \right] \Theta(r) + \frac{1}{2(2\pi)^2 |r|} \int_{-\infty}^0 dk^0 \left[ e^{+ik^0(r^0 + |r|)} + e^{+ik^0(-r^0 + |r|)} \right] \Theta(r) = \\
&= \frac{1}{4\pi |r|} \underbrace{\left[ -\delta(r^0 + |r|) + \delta(r^0 - |r|) \right]}_{=0} \Theta(r) = \frac{1}{4\pi |r|} \delta(r^0 - |r|) \\
\delta(x) &= \frac{1}{2\pi} \int e^{ikx} dk
\end{aligned}$$

$$\begin{aligned} \textcircled{2} \quad \mathcal{L} &= -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{\lambda}{2} (\partial_\mu A^\mu)^2 - \partial_\mu A^\mu \\ &= -\frac{1}{4} (2 \eta^{\mu\sigma} \eta^{\nu\rho} \partial_\mu A_\nu F_{\rho\sigma}) - \frac{\lambda}{2} (\eta^{\mu\nu} \partial_\mu A_\nu)^2 - \eta^{\mu\nu} \partial_\mu A_\nu \end{aligned}$$

$$\text{Équations d'Euler-Lagrange: } \partial_\beta \frac{\partial \mathcal{L}}{\partial (\partial_\beta A_\mu)} = \frac{\partial \mathcal{L}}{\partial A_\mu}$$

$$\frac{\partial \mathcal{L}}{\partial A_\mu} = -\eta^{\mu\nu} \partial_\nu \frac{\partial A_\mu}{\partial A_\mu} = -\eta^{\mu\nu} \partial_\nu \delta_\mu^\mu = -\partial^\mu$$

$$\begin{aligned} \frac{\partial \mathcal{L}}{\partial (\partial_\beta A_\mu)} &= -\frac{1}{2} \eta^{\mu\sigma} \eta^{\nu\rho} \left[ \underbrace{\frac{\partial (\partial_\mu A_\nu)}{\partial (\partial_\beta A_\mu)}}_{\delta_\mu^\beta \delta_\nu^\beta} \cdot F_{\rho\sigma} + \partial_\mu A_\nu \underbrace{\frac{\partial F_{\rho\sigma}}{\partial (\partial_\beta A_\mu)}}_{=0} \right] - \\ &\quad - \lambda \underbrace{\eta^{\mu\nu} \partial_\mu A_\nu}_{(\partial A)} \cdot \eta^{\mu\nu} \underbrace{\frac{\partial (\partial_\mu A_\nu)}{\partial (\partial_\beta A_\mu)}}_{\delta_\mu^\beta \delta_\nu^\beta} - \eta^{\mu\nu} \partial_\nu \underbrace{\frac{\partial A_\mu}{\partial (\partial_\beta A_\mu)}}_{=0} \end{aligned}$$

$$\begin{aligned} \Rightarrow \frac{\partial \mathcal{L}}{\partial (\partial_\beta A_\mu)} &= -\frac{1}{2} \eta^{\mu\sigma} \eta^{\nu\rho} \left[ \delta_\mu^\beta \delta_\nu^\alpha F_{\rho\sigma} + \partial_\mu A_\nu (\delta_\sigma^\beta \delta_s^\alpha - \delta_s^\beta \delta_\sigma^\alpha) \right] - \lambda (\partial A) \eta^{\mu\nu} \delta_\mu^\beta \delta_\nu^\alpha \\ &= -\frac{1}{2} [F^{\mu\alpha} + F^{\alpha\mu}] - \lambda \eta^{\mu\nu} (\partial A) = -F^{\mu\alpha} - \lambda \eta^{\mu\nu} (\partial A) \\ &= F^{\alpha\mu} + \lambda \eta^{\mu\alpha} (\partial A) \end{aligned}$$

$$\begin{aligned} \Rightarrow \partial_\beta \frac{\partial \mathcal{L}}{\partial (\partial_\beta A_\mu)} &= -\partial_\beta F^{\mu\alpha} - \lambda \partial_\beta \eta^{\mu\nu} (\partial A) \\ &= \partial_\beta \partial^\alpha A^\beta - \partial_\mu \partial^\mu A^\alpha - \lambda \partial^\alpha (\partial A) = \\ &= \partial^\alpha \partial_\mu A^\mu - \square A^\alpha - \lambda \partial^\alpha (\partial A) = \\ &= (1-\lambda) \partial^\alpha (\partial A) - \square A^\alpha \end{aligned}$$

$$\Rightarrow \text{Euler-Lagrange: } \square A^\alpha + (\lambda-1) \partial^\alpha (\partial A) = \partial^\alpha$$

$$\textcircled{3} \quad T^{\mu\nu} = -F^{\mu\alpha}F_\alpha^\nu + \frac{1}{4}\eta^{\mu\nu}(F \cdot F)$$

$$\begin{aligned} T^{00} &= -F^{\alpha\beta}F_\alpha^\beta + \frac{1}{4}\underbrace{\eta^{00}}_{=1}(F \cdot F) = \\ &= -\underbrace{F^{00}F_0^0}_{=0} + \underbrace{F^{0i}F_i^0}_{=E^i E_i} + \underbrace{\frac{1}{4}F^{\alpha\beta}F_{\alpha\beta}}_{=F^{00}F_{00}} - \underbrace{F^{0i}F_{0i}}_{=E^i E_i} - \underbrace{F^{i0}F_{i0}}_{=E^i E_i} + \underbrace{F^{ij}F_{ij}}_{\epsilon^{ijk}\epsilon_{ijl}g_k g^l} = 2g_k g^k \end{aligned}$$

$$\Rightarrow T^{00} = \frac{1}{2}E^i E_i + \frac{1}{2}g_k g^k = \frac{1}{2}\vec{E}^2 + \frac{1}{2}\vec{B}^2$$

$$\begin{aligned} T^{0i} &= -F^{\alpha\beta}F_\alpha^i + \frac{1}{4}\underbrace{\eta^{0i}}_{=0}(F \cdot F) = -\underbrace{F^{00}F_0^i}_{=0} + F^{0k}F_k^i = E_k (\epsilon_{k\ell}^i g^\ell) = \\ &= (\vec{E} \times \vec{B})^i \end{aligned}$$

$$\textcircled{4} \quad [A^r(\zeta, \bar{x}), \pi^v(\zeta, \bar{y})] = i\eta^{\mu\nu} \delta(\bar{x} - \bar{y})$$

$$\begin{aligned} &= [A^r(\zeta, \bar{x}), F^{v0}(\zeta, \bar{y}) - \lambda \eta^{v0}(\partial A)] = \\ &= [A^r(\zeta, \bar{x}), F^{v0}(\zeta, \bar{y})] - \lambda [A^r(\zeta, \bar{x}), \eta^{v0}(\partial A)] \stackrel{\lambda=1}{=} \\ &= [A^r(\zeta, \bar{x}), \underbrace{\partial^v A^0(\zeta, \bar{y})}_{=\dot{A}(\zeta, \bar{y})}] - [A^r(\zeta, \bar{x}), \underbrace{\eta^{v0}(\partial^v A_\alpha)}_{\rightarrow \partial^v A^0}] \\ \Rightarrow & -[A^r(\zeta, \bar{x}), \dot{A}(\zeta, \bar{y})] = i\eta^{\mu\nu} \delta(\bar{x} - \bar{y}) \end{aligned}$$

$$\textcircled{5} \quad G_F(k) = \frac{-1}{k^2} = \frac{-1}{(k^0)^2 - \vec{k}^2} = \frac{-1}{(k^0)^2 - \omega_k^2} \quad \vec{k}^2 = \omega_k^2 \quad (\text{car ici } m=0)$$

$$\Rightarrow G_F(x-y) = \int \frac{dk}{(2\pi)^4} G_F(k) e^{-i(x-y)k} = \\ = \underbrace{\int \frac{dk^0}{(2\pi)} \frac{-1}{(k^0)^2 - \omega_k^2} e^{-ik^0(x^0-y^0)} \int \frac{d\vec{k}}{(2\pi)^3} e^{+i\vec{k}(\vec{x}-\vec{y})}}_{=I}$$

Analogue au cas scalaire :

(i)  $x^0 > y^0$ : on boucle le contour "par le bas"

$$I = \int \frac{dk^0}{(2\pi)} \frac{-1}{(k^0 - \omega_k + i\epsilon)(k^0 + \omega_k - i\epsilon)} e^{-ik^0(x^0-y^0)} \stackrel{\text{Théorème des résidus}}{=} \\ = -\frac{2\pi i}{(2\pi)} \left[ e^{-i(\omega_k + i\epsilon)(x^0-y^0)} \frac{-1}{2\omega_k - 2i\epsilon} \right] \xrightarrow{\epsilon \rightarrow 0} \frac{i}{2\omega_k} e^{-i\omega_k(x^0-y^0)}$$

(ii)  $x^0 < y^0$ : on boucle "par le haut"

$$I = -\frac{2\pi i}{2\pi} \left[ e^{-i(-\omega_k + i\epsilon)(x^0-y^0)} \frac{1}{-2\omega_k + 2i\epsilon} \right] \xrightarrow{\epsilon \rightarrow 0} \frac{i}{2\omega_k} e^{+i\omega_k(x^0-y^0)}$$

$$\Rightarrow G_F(x-y) = i \int \frac{d\vec{k}}{(2\pi)^3} e^{+i\vec{k}(\vec{x}-\vec{y})} \left[ \Theta(x^0-y^0) \frac{i}{2\omega_k} e^{-i\omega_k(x^0-y^0)} + \Theta(y^0-x^0) \frac{i}{2\omega_k} e^{+i\omega_k(x^0-y^0)} \right] = \\ = i \int d\vec{k} \left[ \Theta(x^0-y^0) e^{-ik(x-y)} + \Theta(y^0-x^0) e^{+ik(x-y)} \right]$$

$$A_p(x) = \int d\vec{k} \sum_{\lambda} \left[ \underbrace{\alpha(\vec{k}, \lambda) e_p(\vec{k}, \lambda) e^{-ikx}}_{\rightarrow A_p^+(x)} + \underbrace{\alpha^*(\vec{k}, \lambda) e_p(\vec{k}, \lambda) e^{+ikx}}_{A_p^-(x)} \right]$$

$$\langle 0 | A_p(x) A_v(y) | 0 \rangle = \langle 0 | [A_p^+(x) + A_p^-(x)] [A_v^+(y) + A_v^-(y)] | 0 \rangle = \\ = \langle 0 | A_p^+(x) A_v^-(y) | 0 \rangle = \langle 0 | [A_p^+(x), A_v^-(y)] | 0 \rangle$$

$$[A_p^+(x), A_v^-(y)] = \int d\vec{k} \int d\vec{p} e^{-ikx} e^{-ipy} \sum_{\lambda} \sum_{\sigma} g_p(\vec{k}, \lambda) e_v(\vec{p}, \sigma) \underbrace{[\alpha(\vec{k}, \lambda) \alpha^{\dagger}(\vec{p}, \sigma)]}_{\text{relation de fermeture}} =$$

$$= (2\pi)^3 2\omega_k \delta(\vec{k} - \vec{p}) (-\eta^{\lambda\sigma})$$

$$= \int d\vec{k} e^{-ik(x-y)} (-\eta^{\lambda\sigma})$$

$$\Rightarrow \langle 0 | A_p(x) A_v(y) | 0 \rangle = -\eta_{\mu\nu} \int d\vec{k} e^{-ik(x-y)}$$

$$\Rightarrow \langle 0 | T\{A_p(x) A_v(y)\} | 0 \rangle = \Theta(x^0 - y^0) \langle 0 | A_p(x) A_v(y) | 0 \rangle +$$

$$+ \Theta(y^0 - x^0) \langle 0 | A_v(y) A_p(x) | 0 \rangle =$$

$$= -\eta_{\mu\nu} \left[ \Theta(x^0 - y^0) \int d\vec{k} e^{-ik(x-y)} + \Theta(y^0 - x^0) \int d\vec{k} e^{+ik(x-y)} \right]$$

$$= i\eta_{\mu\nu} G_F(x-y)$$