

# **Gamma Rays from Active Galactic Nuclei:**

- I. *Production sites, acceleration/radiation mechanisms*
- II. *Cosmological implications*

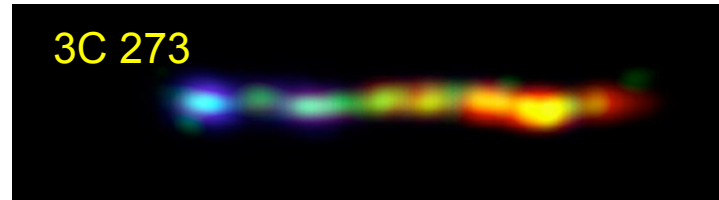
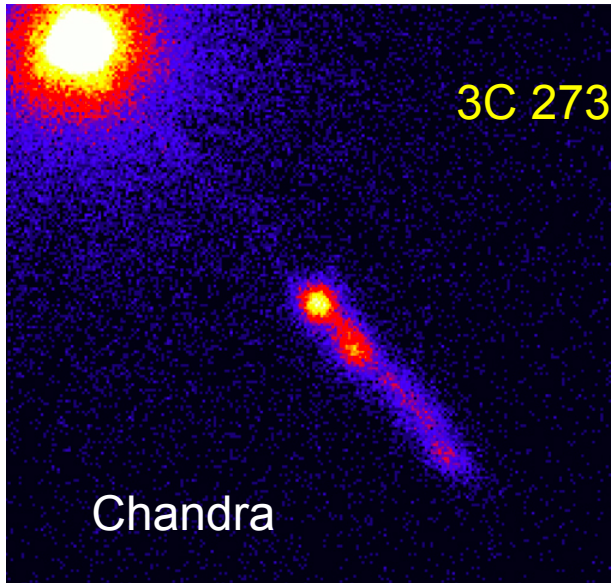
Felix Aharonian

Dublin Institute for Advanced Studies, Dublin

Max-Planck Institut f. Kernphysik, Heidelberg

Nice, September 14, 2012

# Production Sites, Scenarios, Mechanisms



Chandra X  
Hubble O  
Spitzer IR

jet – approximately 20 arcsec

1 arcsec = 3 kpc ( $z=0.158$ ,  $H_0=65$  km/s Mpc)

from R to X rays - nonthermal emission:  
one should expect also gamma-rays on all  
from sub-pc scales to multi-kpc scales!

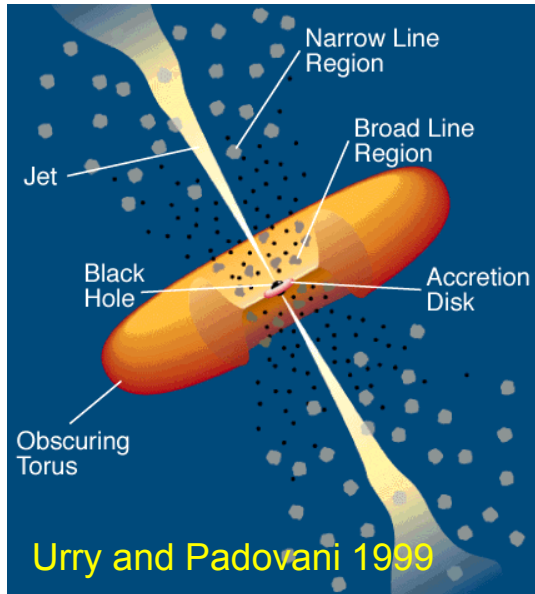
high (GeV) and very high (TeV) energy  $\gamma$ -rays have been reported from hundreds and tens of AGN, respectively !

almost all of them are **blazars** – objects with relativistic jets directed towards the observer

## sites

## acceleration

## radiation



(sub) relativistic shocks  
converter mechanism  
stochastic (Fermi II)  
magnetic reconnection,  
.....

Inverse Compton  
electron synchrotron  
Proton synchrotron  
photomeson processes  
(and subsequent cascades)  
 $\gamma$ - $\gamma$  pair production

plus very complex **magnetohydrodynamics**

BH  
magnetosphere  
sub-parsec jet  
pc-scale  
multi-pc jet  
radio lobes  
magnetized Mpc

⇒ broad range of possible realization (**scenarios**)

unfortunately,  $\gamma$ -ray images of AGN – not possible,  
information only through energy spectra and variability

# *Active Galactic Nuclei: AGN*

- Active Galaxies - galaxies with a compact, variable bright core (nucleus) the radiation of which is comparable or dominates over the emission of the host galaxy produced by stars (and partly reproduced by dust)
- AGN - central engines powered by Black Holes of  $10^6$  to  $10^{10} M_{\odot}$
- Radiation of AGN: thermal emission of the accretion flow and nonthermal emission of the relativistic jet
- several types of AGN - distinction between them sometimes arbitrary, but it is clear that we deal with several AGN populations - Seyfert galaxies, Quasars, Radio galaxies, BL Lac Objects,...

*difference between these classes?*

**spectral lines** e.g. Seyferts - rich line spectra, BL Lacs - featureless continuous  
**large scale morphology** e.g. radio galaxies - often with twin-lobed structures  
**luminosity** e.g. Seyferts -  $10^{43}$ - $10^{45}$  erg/s; QSO:  $> 10^{46}$  erg/s  
**polarization, variability**, etc.

## *classification and unification of AGN?*

**classification** of source populations is not a trivial task in astronomy, especially for AGN with very diverse spectral and temporal characteristics:

they can be biased because of different selection effects, e.g. because of strongly variable behavior of AGN. This is true especially for  **$\gamma$ -ray emitting compact, short-lived nonthermal structures** (solitary events) with highly variable parameters which contain large uncertainties (e.g. strength of the magnetic field, the speed and viewing angle of the moving emitter, etc.)

e.g. recently Giommi et al (2012) demonstrated, based on the WMAP, Planck and Fermi LAT data that the so called (very popular) “blazar sequence” (correlation between the luminosity **L** and the position of the synchrotron peak  **$\nu$** ) is a selection effect “arising from the comparison of shallow radio and X-ray surveys”.

# *classification and unification of AGN?*

## **Unification of AGN:**

- AGN basically are of two types – Seyferts and Quasars – and difference between them is in the luminosity ( $L_{\text{QSO}}/L_{\text{Sey}} > 100$ )
- Seyfert 1 and Seyfert 2 galaxies are intrinsically the same sources. In the case of Seyfert 2 galaxies we do not see the nuclear source because of absorption of optical lines due to the unfavorable viewing angle
- Blazars – radio loud AGN when jet axis is close to the line of sight

Unification approach based on **anisotropy** due to orientation of jets (nonthermal emission) and absorption of thermal emission in dusty torus, generally provides a useful tool for systematization of large data sets, but still it seems to be a simplification of the real picture

# AGN classifications: an example

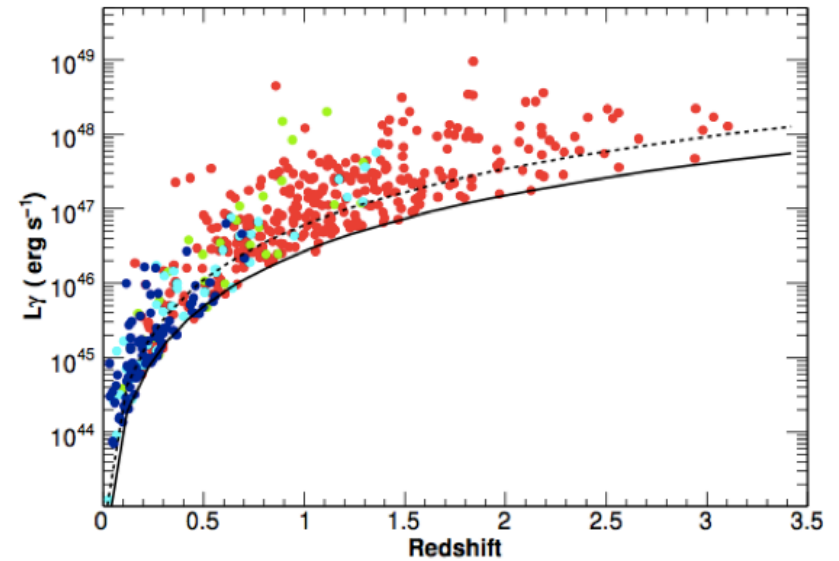
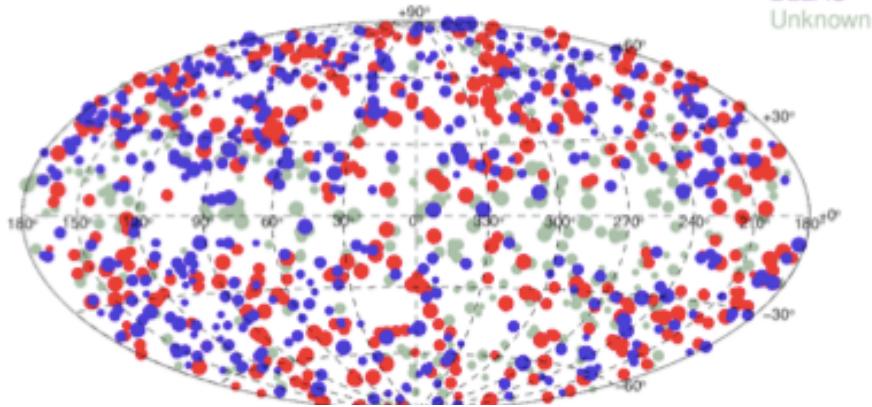
Radio quiet		Radio loud
Radio quiet quasar (RQQ) <i>Broad absorption line (BAL)</i>	Type 1	Radio loud quasar (RLQ) <i>Steep radio spectrum (SSRLQ)</i> <i>Flat radio spectrum (FSRLQ)</i>
Seyfert 1 <i>Sy 1.0.....1.9</i> <i>Narrow line Sy 1 (NLS1)</i>		Broad line radio galaxy (BLRG)
Seyfert 2 <i>NL X-ray galaxy (NLXG)</i>	Type 2	Narrow line radio galaxy (NLRG)
LINER	Type 3	Weak line radio galaxy (WLRG)
	Type 0	Blazar: BL Lac/OVV

from C. Tadhunter, New AR 2008

most of gamma-ray emitting AGN are BL Lac Objects at GeV and TeV energies or Optically Violently Variable (OVV) quasars at GeV (in a few cases at TeV) energies

# GeV AGN

Second LAT Catalogue (2LAC)  
TS>25, August 2008 – August 2010



only most powerful QSOs can be detected from large  $z$ :  $L_\gamma = 10^{46} (f_\gamma / 10^{-10} \text{ erg/cm}^2 \text{ s}) (d/3\text{Gpc})^2 \text{ erg/s}$

FERMI AGN: 310 FSRQs  
395 BL Lacs  
156 unknown type  
26 other AGN

these are apparent luminosities; “true” (intrinsic) luminosities cannot exceed  $10^{48} \text{ erg/s}$  even for  $10^{10} \text{ Mo}$  SBH; we see Doppler boosted radiation:

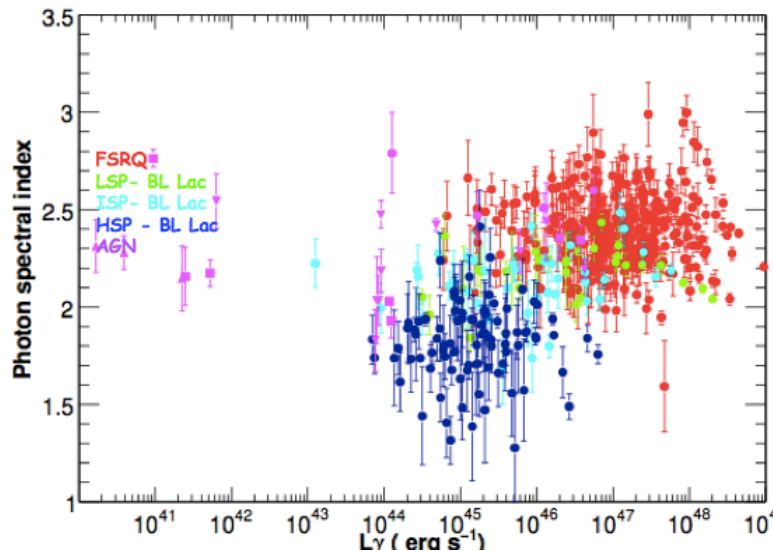
$$L_{\text{int}} \rightarrow \delta^{-4} L_{\text{app}}$$

$$E \rightarrow \delta^{-1} E_{\text{obs}}$$

$$\Delta t \rightarrow \delta \Delta t_{\text{obs}}$$

$\delta$  – Doppler factor: 10-100

# FERMI LAT: higher luminosity – harder spectrum

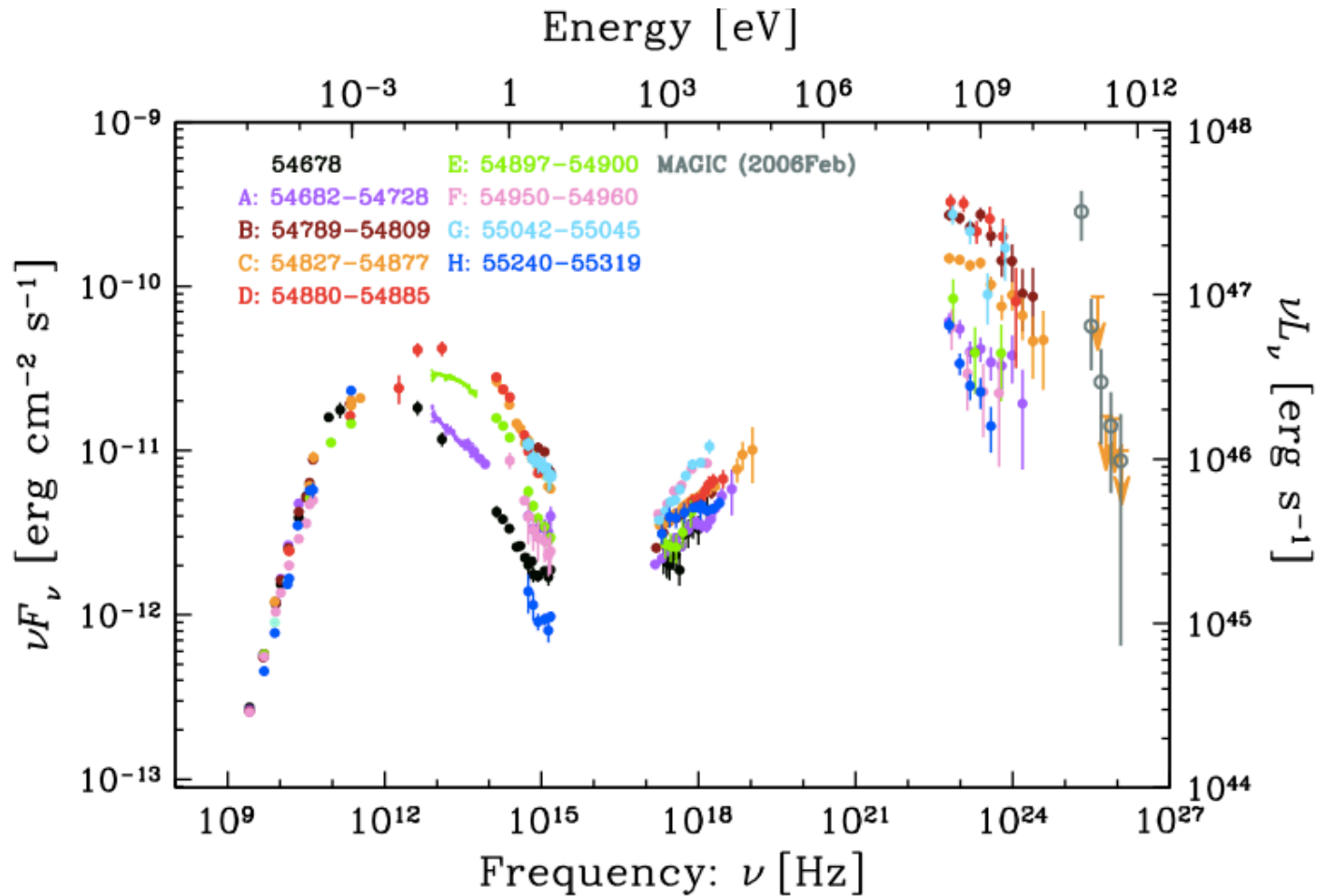


this confirms to some extent the fact that the TeV source population is dominated by less powerful BL Lacs

Reason? could be more than one reason, e.g.

- Quasars are distant sources, thus VHE  $\gamma$ -rays are absorbed due to interactions with EBL
- Quasars are very powerful sources - high density radiation fields may lead to
  - (1) “early” cutoffs in the electron spectra (due to enhanced Compton cooling) and correspondingly in the IC gamma-ray spectra
  - (2) internal photon-photon absorption
  - (3) could be both, but also something else

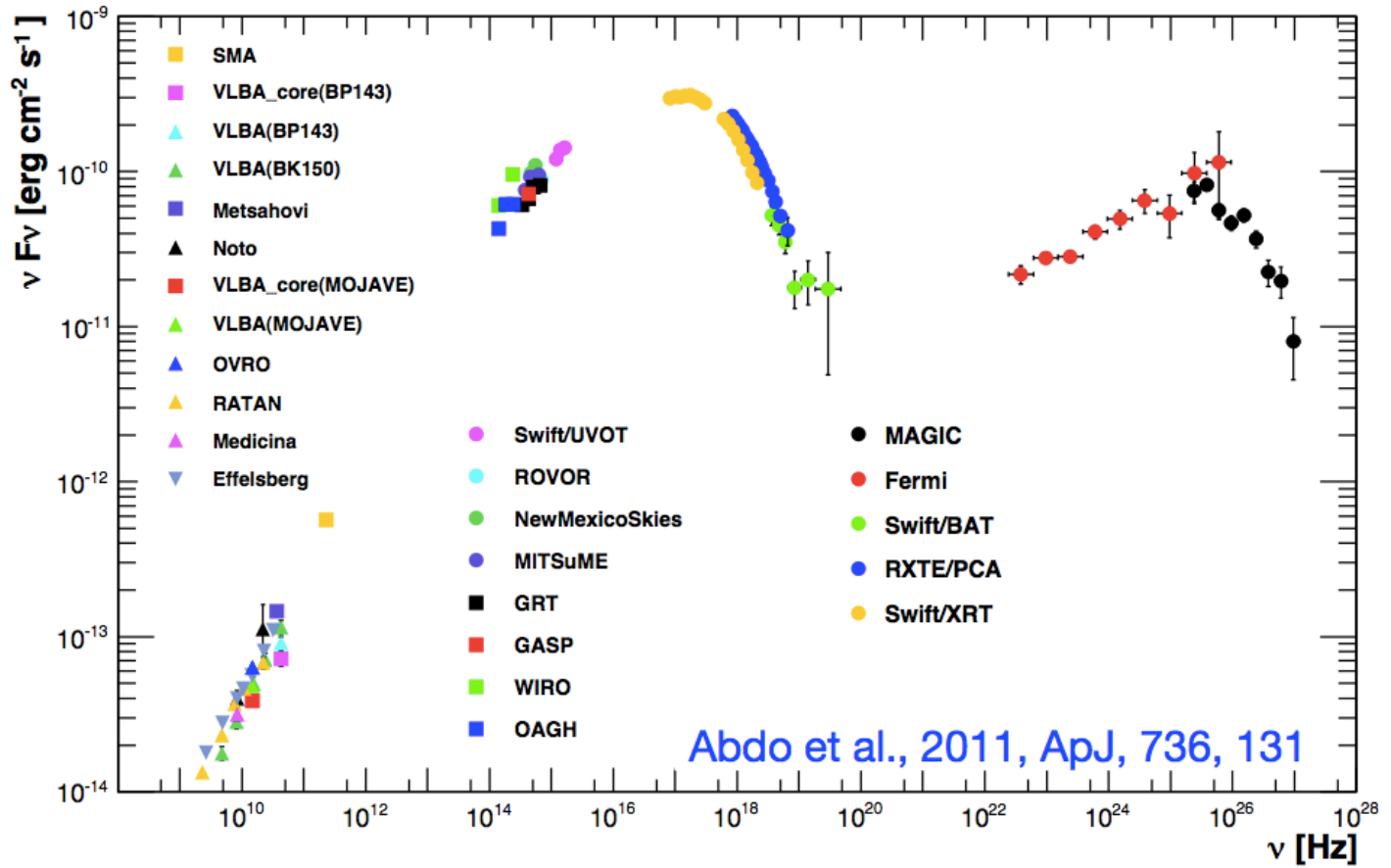
# SED of 3C 279 – a classical GeV blazar



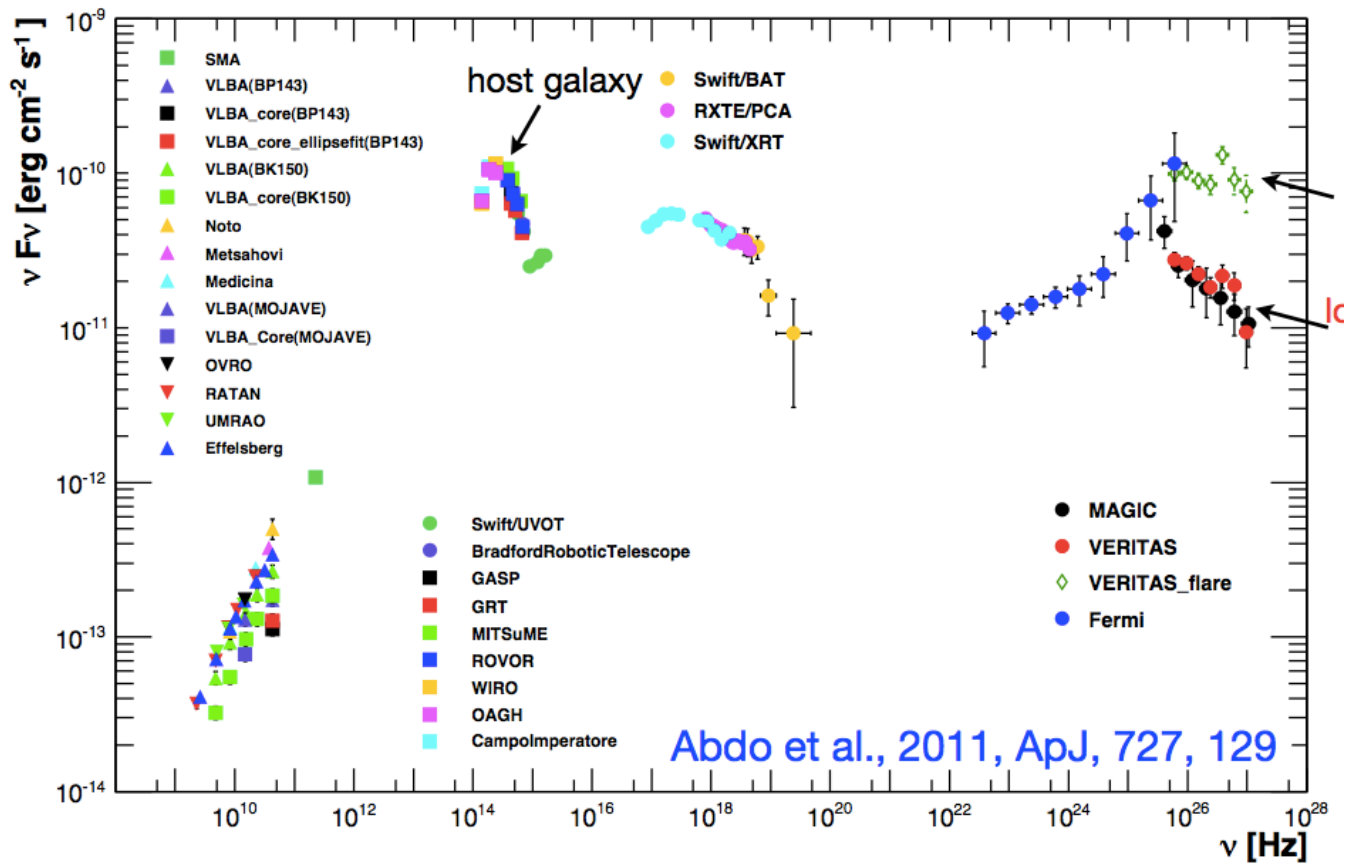
$L_\gamma/L_S > 10$  | Synchrotron peak at mm & MIR | X-rays of IC origin | variability - days

TeV emission?

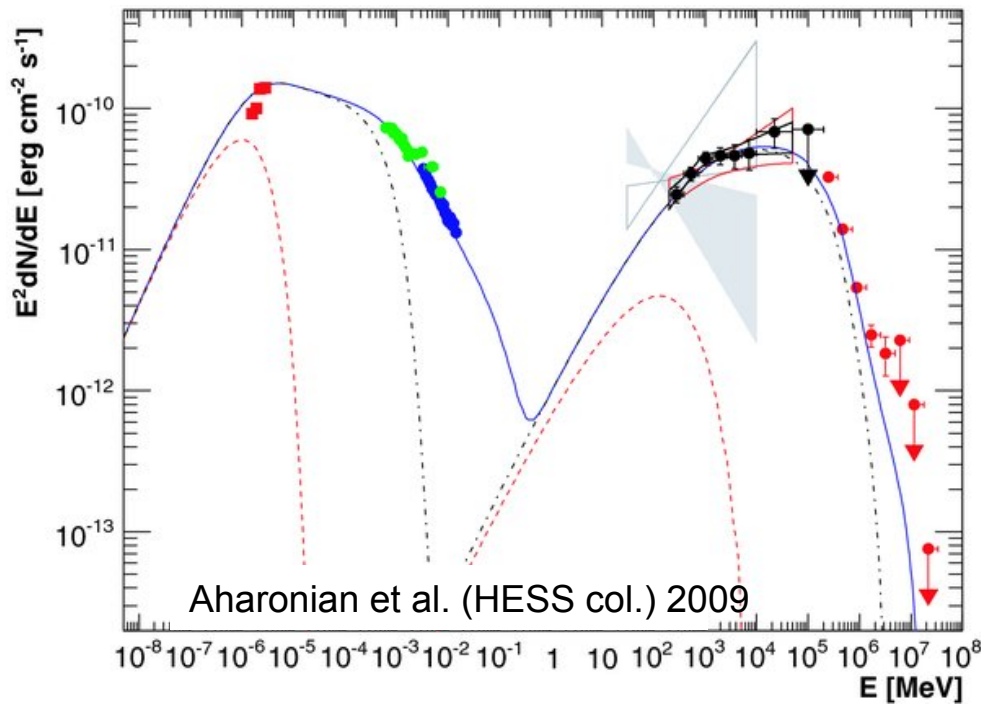
# *a typical TeV blazar: Mkn 421*



# *a typical TeV blazar: Mkn 501*



# TeV blazar PKS 2155-304 in a low state



## one zone SSC model

electron spectrum:

power-law with two breaks:

$p_1=1.3$ ,  $p_2=3.2$ ,  $p_2=4.3$

$\gamma_1=1.4 \times 10^4$ ,  $\gamma_2=2.3 \times 10^5$

and cutoff at  $\gamma_{\max}=3 \times 10^6$

$B=0.02$  G,  $R=2 \times 10^{17}$  cm,  $\delta=32$

with a **standard remark**:

“SSC is a simplification but it satisfactorily explains the SED and gives correct parameters”

**my opinion:** *I cannot agree with such a statement: it is poorly justified and misleading if the low-state SED in reality is a superposition of weak flares – the parameters will be different; also usually the role of multiple “breaks” is overestimated, while the shape at cutoff is underestimated. What to do? Better nothing, rather than ‘one-zone model’ “industry”.*

*When we deal with AGN we should remember that*

generally the phenomena relevant to compact relativistic objects proceed under extreme physical conditions in environments characterized with

- *huge gravitational, magnetic and electric fields,*
- *very dense background radiation,*
- *relativistic bulk motions (black-hole jets and pulsar winds)*
- *shock waves, highly excited (turbulent) media, etc.*

any coherent description and interpretation of phenomena related to high energy cosmic gamma-rays requires knowledge and deep understanding of many disciplines of experimental and theoretical physics, including

*nuclear and particle physics,  
quantum and classical electrodynamics,  
special and general relativity,  
plasma physics, (magneto) hydrodynamics, etc.*

and (of course) **Astronomy&Astrophysics**

## *gamma-ray emitting AGN*

*everything should proceed with an **extreme efficiency!***

conversion of gravitational, thermal, bulk motion, electromagnetic forms of energy to nonthermal relativistic particles, i.e. effective acceleration of particles coupled with favorable conditions for production of  $\gamma$ -rays

*SMBH and relativistic Doppler boosting* – **not sufficient:** we need extremely effective particle accelerators and effective emitters

## *Extreme Accelerators*

*machines where acceleration proceeds with efficiency close to 100%*

(i) fraction of available energy converted to nonthermal particles

*in PWNe and perhaps also in SNRs and AGN can be as large as 50 %*

(ii) maximum (theoretically) possible energy achieved by individual particles

*acceleration rate close to the maximum (theoretically) possible rate :*

$t_{\text{acc}} = \eta(\epsilon) E / ecB = \eta(\epsilon) r_L / c$  ;  $E_{\text{eff}} = \eta^{-1} B$  - projection of the electric field on the particle trajectory averaged when particle moves along trajectory; generally  $E \ll B$ , i.e.  $\eta \gg 1$   
 $\eta = 1$  extreme accelerator ( $\eta < 1 \Rightarrow E > B$  - cannot be excluded)

### analogy with X-ray Astronomy:

as cosmic plasmas are easily heated up to **keV temperatures** - almost everywhere, particles (electrons and protons) can be easily accelerated to **TeV energies** - almost everywhere, especially in objects containing relativistic outflows - jets & winds

# AGN and $10^{20}$ eV Cosmic Rays

the very fact of existence of such particles implies existence of extragalactic extreme accelerators...

the “Hillas condition” -  $l > R_L$  - an necessary but not sufficient condition...

- (i) maximum acceleration rate allowed by classical electrodynamics  
 $t^{-1} = \eta q B c$  or  $c/R_L$  with  $\eta \sim 1$  and  $\sim (v/c)^2$  in shock acceleration scenarios)
- (ii) B-field cannot be arbitrarily increased - the synchrotron and curvature radiation losses become a serious limiting factor, unless we assume perfect linear accelerators ...

only a few options survive from the original Hillas (“l-B”) plot:  
 $>10^9 M_\odot$  BH magnetospheres, small and large-scale AGN jets, GRBs

suspected sites of  $10^{20}$  eV cosmic rays based on the condition: source size  $>$  Larmor radius

$$(R/1\text{pc})(B/1\text{G}) > 0.1 (E/10^{20}\text{eV})$$

necessary but not sufficient condition: it implies

(1) minimum acceleration time

$$t_{\text{acc}} = R_L/c = E/eBc$$

acceleration in fact is slower:

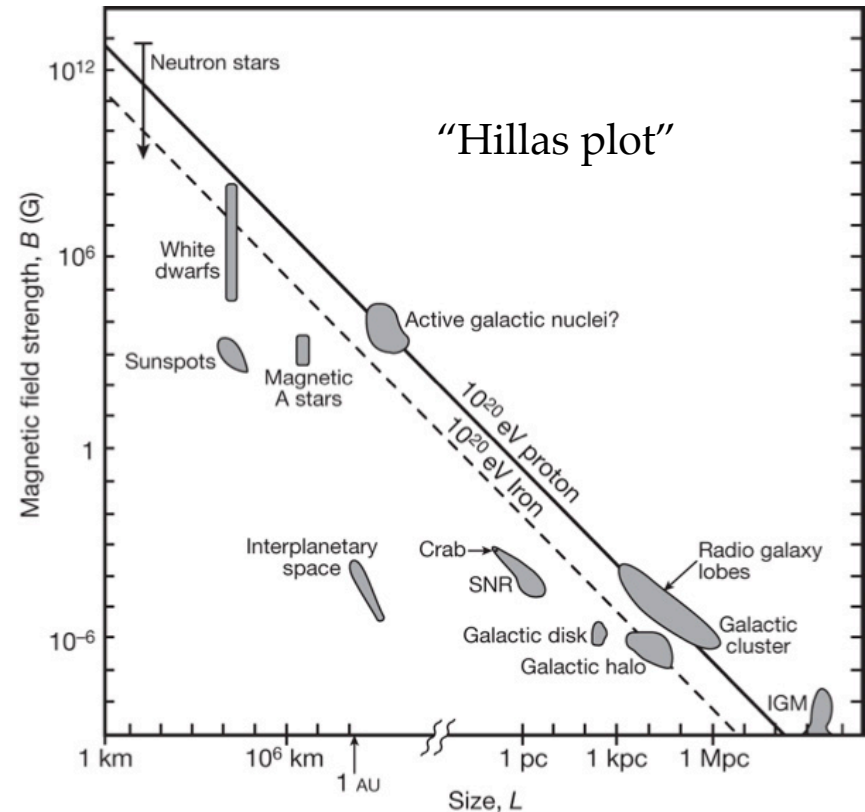
$$t_{\text{acc}} = (1-10)\eta R_L/c (c/v)^2$$

with  $\eta > 1$  and shock/bulk-motion speed  $v < c$  ( $\eta=1$  - Bohm diffusion)

Compact objects like AGN and GRBs are the best candidates

(2) no energy losses

but synchrotron/curvature losses in compact objects become severe limiting factor



PM Bauleo & JR Martino Nature 458, 847-851 (2009)

# acceleration sites of $10^{20}$ eV CRs ?

$$t_{\text{acc}} = \frac{R_L}{c} \eta^{-1}$$

signatures of extreme accelerators?

✓ **synchrotron self-regulated cutoff:**

$$h\nu_{\text{cut}} = \frac{9}{4} \alpha_f^{-1} mc^2 \eta :$$

$\simeq 300\text{GeV}$  proton synchrotron

$\simeq 150\text{MeV}$  electron synchrotron

✓ **neutrinos** (through “converter” mechanism) production of neutrons (through  $p\gamma$  interactions) which travel without losses and at large distances convert again to protons  $\Rightarrow \Gamma^2$  energy gain !

*Derishev, FA et al. 2003, Phys Rev D 68 043003*

✓ **observable off-axis radiation**

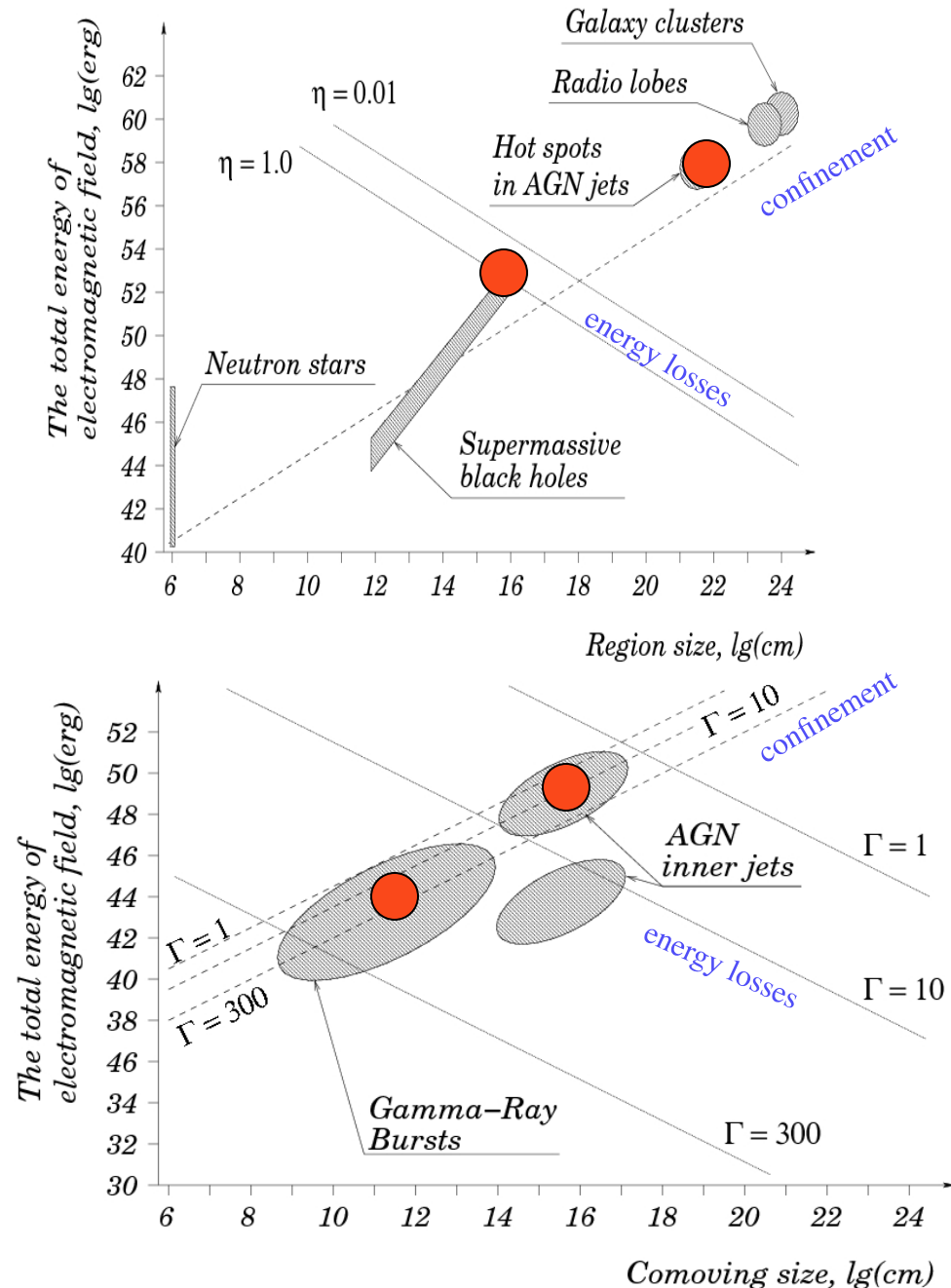
radiation pattern can be much broader than  $1/\Gamma$

*Derishev, FA et al. 2007, ApJ, 655, 980*

$$\eta \approx 0.1 (v_{\text{shock}}/c)^2$$

\*) in nonrelativistic shocks

*FA, Belyanin, Derishev 2002, Phys Rev D, 66, id. 023005*



## *radiation and absorption processes*

any interpretation of an astronomical observation requires

- ✓ unambiguous identification of radiation mechanisms and
- ✓ good knowledge of radiation and absorption processes

gamma-ray production and absorption processes:

*several but well studied*

## interactions with matter

E-M:

VHE

bremsstrahlung:	$e N(e) \Rightarrow e' \gamma N(e)$	*	$E_\gamma \sim 1/2 E_e$
pair production	$\gamma N(e) \Rightarrow e^+ e^- N(e)$	*	
e+e- annihilation	$e^+ e^- \Rightarrow \gamma \gamma$ (511 keV line)		

Strong/weak:	$pp (A) \Rightarrow \pi, K, \Lambda, \dots$	**	$E_\gamma \sim 1/10 E_p$
	$\pi, K, \Lambda \Rightarrow \gamma, \nu, e, \mu$		
	$\mu \Rightarrow \nu$		

also in the low energy region

Nuclear:	$p A \Rightarrow A^* \Rightarrow A' \gamma, n$		
	$n p \Rightarrow D \gamma$ (2.2 MeV line)		

# interactions with radiation and B-fields

## Radiation field

VHE

### E-M:

inverse Compton:  
 $\gamma\gamma$  pair production

$$e \gamma (B) \Rightarrow e' \gamma$$

$$\gamma \gamma (B) \Rightarrow e^+e^-$$

\*\*  $E\gamma \sim \epsilon(Ee/mc^2)^2$  (T) to  $\sim Ee$  (KN)

\*\*

### Strong/week

$$p \gamma \Rightarrow \pi, K, \Lambda, \dots$$

$$\pi, K, \Lambda \Rightarrow \gamma, \nu, e, \mu$$

$$\mu \Rightarrow \nu$$

\*

$$E\gamma \sim 1/10 E p$$

$$A \gamma \Rightarrow A^* \Rightarrow A' \gamma$$

\*

$$E\gamma \sim 1/1000 A E p$$

## B-field

synchrotron  
 pair production

$$e (p) B \Rightarrow \gamma$$

$$\gamma B \Rightarrow e^+e^-$$

\*

\*

$$E\gamma \sim B E_e^2; h\nu_{\max} \sim \alpha^{-1} mc^2$$

## *leptonic or hadronic?*

gamma-rays produced in interactions of electrons and protons/nuclei often are called  
**leptonic** and **hadronic** interactions

but it is more appropriate to call them as **E-M** (electromagnetic) and **S** (strong)

examples:

(i) *synchrotron radiation of protons - pure electromagnetic process*

*interaction of hadrons without production of neutrinos*

(ii) *photon-photon annihilation  $\Rightarrow \mu^+\mu^- \Rightarrow$  neutrons, antineutrinos*

*production of neutrinos by photons as parent particles*

**E-M** are calculated with high accuracy and confirmed experimentally

**S** are well studied experimentally and explained theoretically

*often several processes proceed together  $\Rightarrow$  cascades in matter, radiation and B-fields*

# *gamma-ray production efficiency*

many reported TeV gamma-ray sources require not only extreme particle accelerators but also effective production of gamma-rays

## effective gamma-ray production?

cooling time of the given gamma-ray production process is shorter than

- (1) timescales of radiative and non-radiative (e.g. adiabatic) losses
- (2) intrinsic dynamical (source age, acceleration time, particle escape time)

**Note:** high efficiency is an important but not sufficient/decisive condition for a gamma-ray sources to be detected. The detectability depends also on

- ✓ power and distance to the source ( $\sim W/d^2$ )
- ✓ beaming factor, e.g. Doppler boosting ( $\sim \delta^4$ )
- ✓ Sensitivity of the instrument in the given energy domain

*inefficient!*

## Nonthermal X-ray Bremsstrahlung

at first glance quite attractive (“*why should I invoke multi-TeV electrons to produce X-rays when can I use keV electrons to produce keV X photons?*”) in fact only less than  $10^{-5}$  fraction of the kinetic energy of electrons (protons) is released in X-rays; 99.99...% goes to the ionization and heating of the gas

$$L_e > 10^5 L_X = 10^{37} (f_X / 10^{-12} \text{ erg/s}) (d/1\text{kpc})^2 \text{ erg/s}$$

the same is true for gamma-ray line emission due to excitation of nuclei by sub-relativistic protons - both mechanisms “work” during Solar flares, otherwise it typically leads to unreasonably high requirements for production rate of sub-relativistic electrons - this makes the extremely interesting issues like detection of gamma-ray lines, in particular from ISM, SNRs, GMCs, etc (information about the sub-relativistic CRs !) observationally very difficult

we are unlucky with prompt gamma-ray line astronomy!

$pp \rightarrow \pi^0 \rightarrow 2\gamma$

*not very efficient*

no competing dissipation mechanisms - in “calorimetric scenarios”:  $L_\gamma \sim L_p/3$   
but the process itself is not very fast/relatively slow:  $t_\pi \sim 10^{15} (n/1\text{cm}^{-3})^{-1} \text{ s}$   
usually the source age or particle escape is a big issue !

SNRs: typical density:  $n \sim 1\text{cm}^{-3}$ , magnetic field  $B \sim 100\mu\text{G}$ , size  $R \sim 3 \text{ pc}$  assuming Bohm diffusion,  $D(E) = r_L c/3 = 10^{25} (E_p/10\text{TeV})^{-1} \text{ cm}^2/\text{s}$ , escape time of protons which produce 1 TeV gamma-rays:  $t_{\text{esc}} \sim R^2/D \sim 10^{13} \text{ s} \sim 0.01 t_\pi$

GMCs: typical densities  $n > 100\text{cm}^{-3}$ , size  $R > 10\text{pc}$ , but  $>1000$  times faster diffusion :  $t_{\text{esc}} \sim 10^{11} \text{ s} \sim 0.01 t_\pi \Rightarrow$  the same  $\sim 1\%$  efficiency

Galaxy densities  $n \sim 10^{-3} \text{ s}$ , size  $R > 1\text{Mpc}$  - full confinement!

Clusters:  $t_\pi < 10^{18} (n/1\text{cm}^{-3})^{-1} \text{ s}$  - comparable to the age (Hubble time) !

$\gamma$ Binaries: protons accelerated by the compact object and interacting with the dense stellar disk of companion:  $n \sim 10^{13} \text{ cm}^{-3}$  ; the cooling time could be shorter than escape time  $\Rightarrow$  potentially effective production of gamma-rays and  $\nu$ s

*higher efficiencies at MeV/GeV energies because of problem of confinement*

# Synchrotron radiation

*very efficient*

especially in extreme particle accelerators where acceleration proceeds at the maximum (theoretically possible) rate and the further acceleration is limited by synchrotron losses

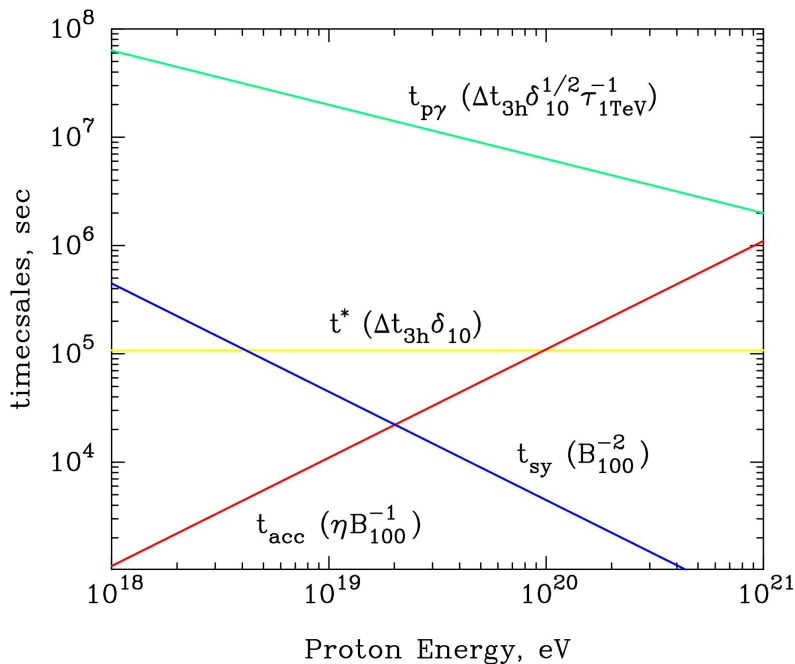
$$t_{\text{acc}} = \frac{R_L}{c} \eta^{-1} \quad \Rightarrow \quad \text{self regulated cutoff}$$

$$h\nu_{\text{cut}} = \frac{9}{4} \alpha_f^{-1} mc^2 \eta :$$

$$\simeq 300 \text{ GeV} \quad \text{proton synchrotron}$$

$$\simeq 150 \text{ MeV} \quad \text{electron synchrotron}$$

proton-synchrotron is effective in compact objects with large B-fields (when  $t_{\text{coo}} < R/c$ )



$$t_{\text{synch}} = 4.5 \times 10^4 (B/100 \text{ G})^{-2} (E/10^{19} \text{ eV})^{-1} \text{ s}$$

$$t_{\text{acc}} = 1.1 \times 10^4 (E/10^{19}) (B/100 \text{ G})^{-1} \text{ s}$$

$$E_{\text{max}} \sim B^{-1/2}, \text{ but } h\nu_{\text{cut}} - \text{independent of } B$$

$$t(h\nu_{\text{cut}}) = 2.4 \times 10^4 (B/100 \text{ G})^{-3/2} \eta^{1/2} \text{ s} < R/c$$

$$B > 100 (R/10^{15} \text{ cm})^{-2/3} \eta^{1/3} \text{ G}$$

## Synchrotron radiation

*very efficient*

*do we have evidence for signatures of extreme accelerators?*

**electron synchrotron** - most likely in the spectrum of the Crab Nebulae  
**protons synchrotron** - in some blazars, GRBs ?

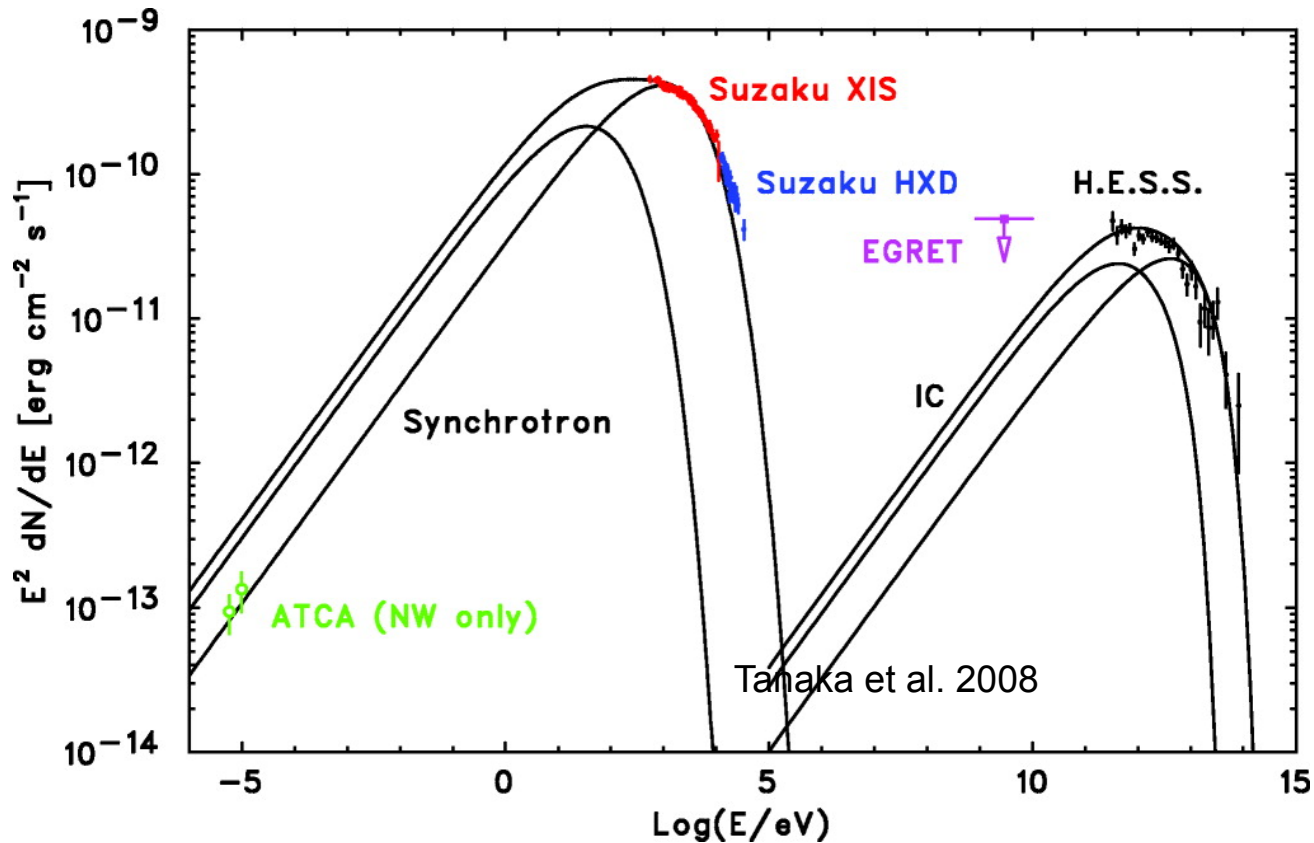
factors reducing the maximum energies of the synchrotron cutoff?

- *radiative losses in the case of electrons-synchrotron (in binaries)*
- *not sufficiently strong B-fields in the case of proton-synchrotron*

**position of the synchrotron peak as indicator of acceleration efficiency**

electron synchrotron efficiency could be close 100% even in non-extreme accelerators though the radiation at lower energies, e.g. in young SNRs

# Sy-IC SED of SNR RXJ 1713-4639

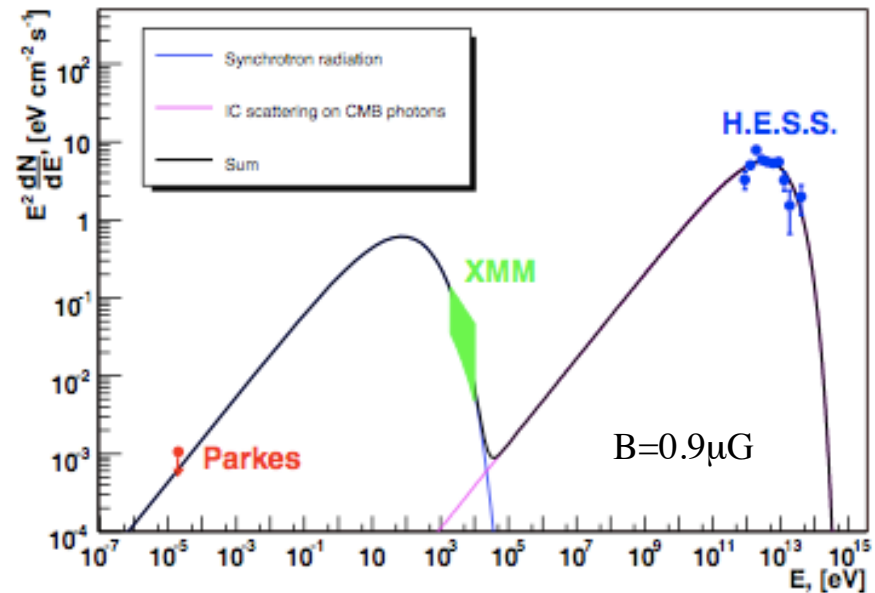
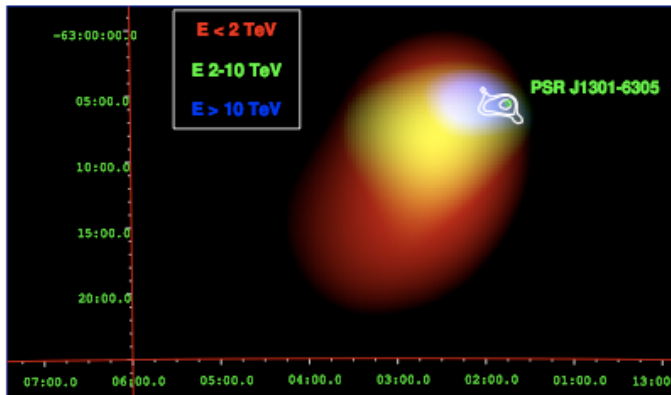


two characteristic Synchrotron and IC peaks;  $h\nu_{\text{max}} \sim 1 (v/3000\text{km/s})^2 \kappa \text{ keV}$ ;  
 $\kappa = 1$  implies maximum effective Bohm diffusion regime for DSA

PWNe - *perfect electron accelerators and perfect  $\gamma$ -ray emitters!*

- (1) rot. energy  $\Rightarrow$  (2) Poynting flux  $\Rightarrow$  (3) cold ultrarelativistic wind  $\Rightarrow$   
 (4) termination of the wind/acceleration of electrons  $\Rightarrow$  gamma-radiation:  
*efficiency at each stage  $>50\%$  ! but synchrotron peak below 1keV*

HESS J 13030-62 = PSR J1301-6305?

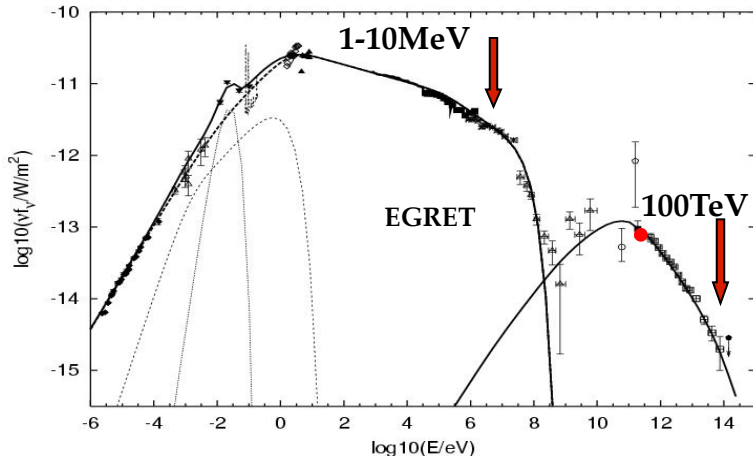


dramatic reduction of the angular size with energy: strong argument in favor of the IC origin of the  $\gamma$ -ray nebula

very small average B-field; for  $d=12.6\text{kpc}$   
 $L_\gamma/L_{SD} = 0.07$ ;  $3\text{arcmin} \sim 10\text{ pc}$

because of small B-field we see “relic” electrons produced at early epochs of the pulsar

# Crab Nebula – a perfect electron PeVatron and an extreme accelerator



standard MHD theory (Kennel&Coroniti)

cold ultrarelativistic pulsar wind terminates by reverse shock resulting in acceleration of multi-TeV electrons

synchrotron radiation => **nonthermal optical/X** nebula

Inverse Compton => **high energy gamma-ray** nebula

Crab Nebula – a powerful  $L_e = 1/5 L_{\text{rot}} \sim 10^{38}$  erg/s

and extreme accelerator:  $E_e \gg 100$  TeV

$$E_{\text{max}} = 60 (B/1\text{G})^{-1/2} \eta^{-1/2} \text{ TeV} \text{ and } h\nu_{\text{cut}} \sim 150\eta^{-1} \text{ MeV}$$

Cutoff at  $h\nu_{\text{cut}} \sim 100$  MeV =>  $\eta \sim 1$  – acceleration at the maximum rate

Flares! maximum of SED beyond 100 MeV =>  $h < 1$  (or  $E > B$ ) or Doppler boosting?

(see lecture of M. Tavani)

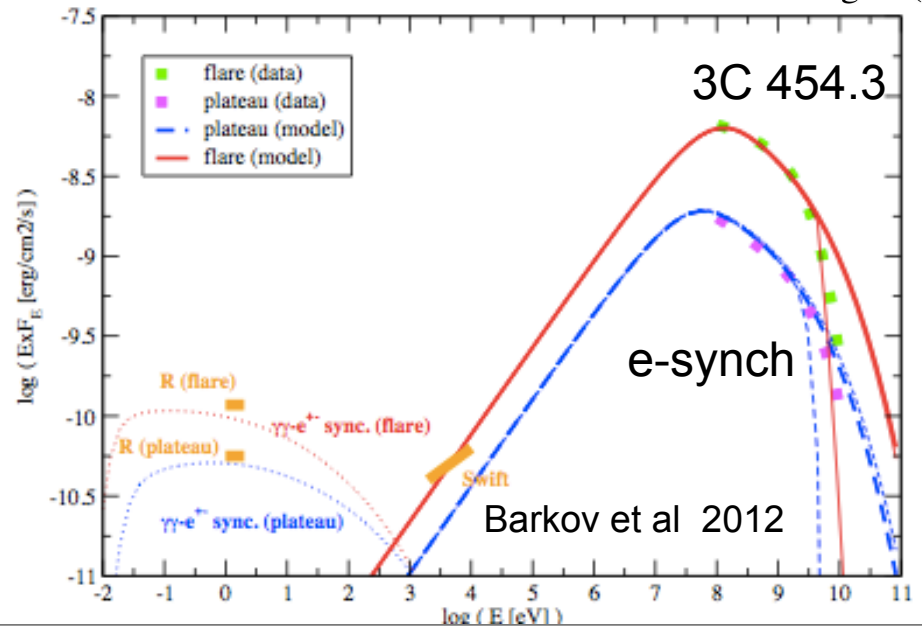
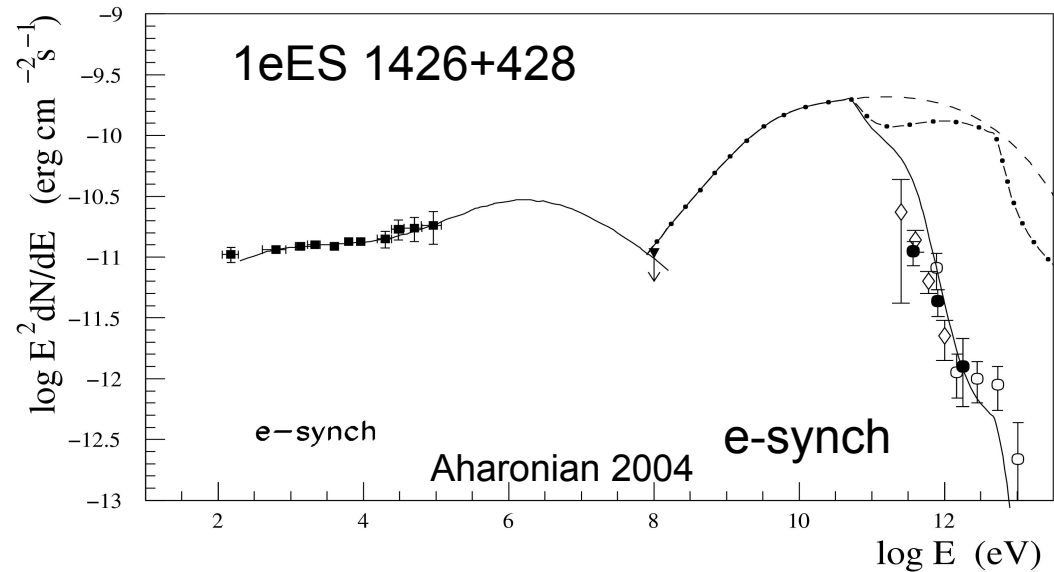


# Can AGN operate like the Crab Nebula

proton synchrotron as an alternative for explanation of the TeV emission spectra of BL LACs ?

boosted electron synchrotron can explain the GeV emission of FSRQ?

Internal  $\gamma$ - $\gamma$  absorptions provides Better fit, and explains the low-energy radiation by synchrotron radiation of secondary electrons



SSC or EXC models: while the production of nonthermal radiation is efficient, location of the synchrotron peak below X-ray band implies very low,  $\eta < 10^{-6}$  (!?) acceleration rate?

IC:  $e\gamma \rightarrow e+\gamma'$

*very efficient*

- ✓ compact objects - binaries, AGN... -  
*very effective with some exceptions*
- ✓ PWNe with very small B-field:  $L_{\text{IC}}=L_e(w_{2.7\text{K}}/w_B)=L_e(B/3\text{mG})^{-1} \sim L_e$   
if  $B < 3\text{mG}$ ; thanks to very effective (relativistic shock?) acceleration  
*electrons still can be accelerated to 100 TeV or beyond*
- ✓ Clusters of Galaxies - despite small B-field ( $\sim 1\mu\text{G}$ ) and limited shock  
speed ( $\sim 2000\text{ km/s}$ ), thanks to the large size and age of these cosmological  
structures, protons can be accelerated to  $10^{18}\text{-}10^{19}\text{ eV}$ , produce secondary  
(Bethe-Heiliter) pairs at interactions with 2.7K CMBR, and the secondary  
electrons can produce effective IC gamma-rays upscattering 2.7K CMBR
- ✓ many other realizations....

*when in the Klein-Nishina regime, IC is accompanied by  $\gamma\text{-}\gamma$  pair-production*

# Hadronic vs. Electronic models of TeV Blazars

SSC or external Compton – *currently most favoured models:*

- easy to accelerate electrons to TeV energies
  - easy to produce synchrotron and IC gamma-rays
- recent results require more sophisticated leptonic models*

## Hadronic Models:

- **protons interacting with ambient plasma** neutrinos
- **protons interacting with photon fields** neutrinos\*  
low efficiency + severe absorption of TeV  $\gamma$ -rays
- **proton synchrotron** no neutrinos  
very large magnetic field  $B=100\text{ G}$  + acceleration rate  $c/r_g$

*“extreme accelerator” (of EHE CRs) Poynting flux dominated flow*



\*detectable neutrinos from EGRET AGN but not from TeV blazars

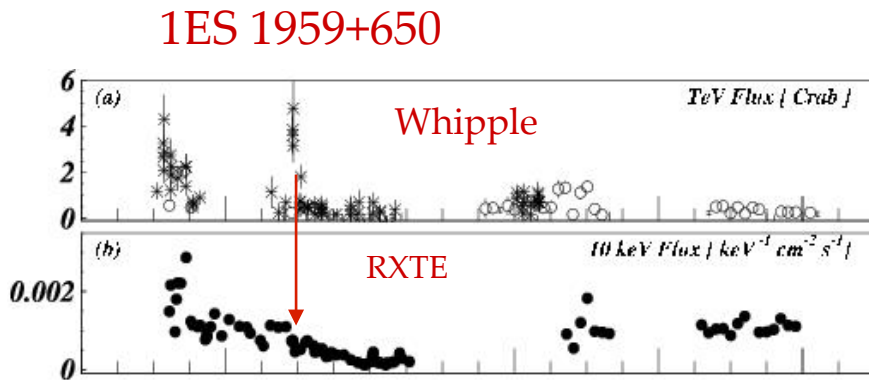
## Gamma-ray emission of Blazars

large Doppler factors: make more comfortable the interpretation of variability timescales (larger source size, and longer acceleration and radiation times), reduces (by orders of magnitude) the energy requirements, allow escape of GeV and TeV  $\gamma$ -rays ( $t_{\gamma\gamma} \sim \delta_j^6$ )

uniqueness: Only TeV radiation tells us unambiguously that particles are accelerated to high energies (one needs at least a TeV electron to produce a TeV photon) in the jets with Doppler factors  $> 10$  otherwise gamma-rays cannot escape the source due to severe internal photon-photon pair production

combined with synchrotron: derivation of several basic parameters like B-field, total energy budget in accelerated particles, thus to develop a quantitative theory of MHD, particle acceleration and radiation in relativistic jets, although yet with many conditions, assumptions, caveats...

## deviations from standard concepts

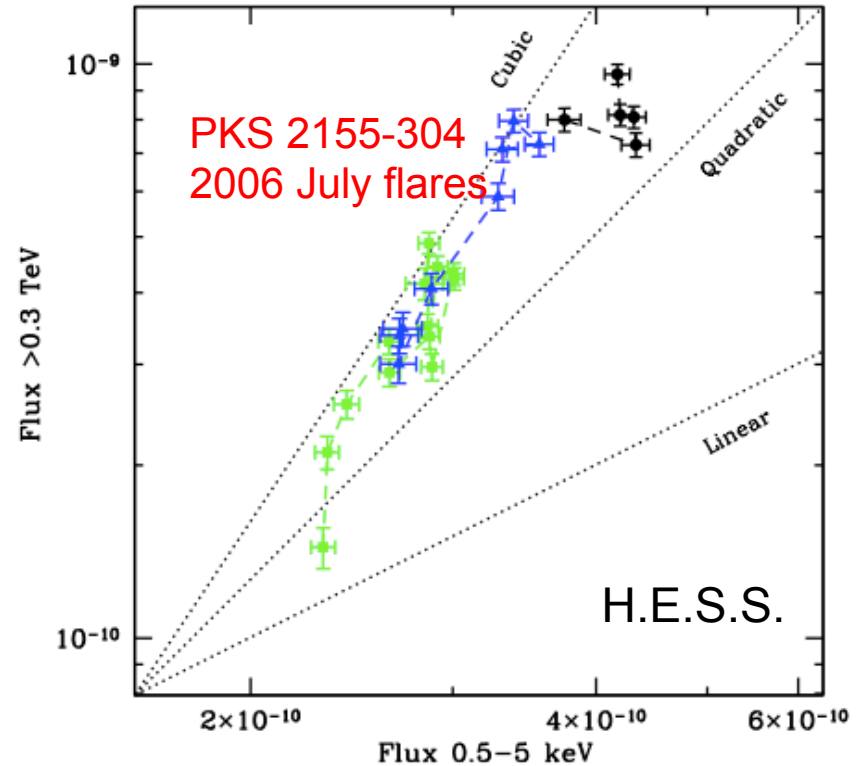


“orphan“  $TeV$  flare:  
no  $X$ - $TeV$  correlations

contradicts to the concept of  
the Compton origin of  $\gamma$ -rays ?

not really... there could be several natural  
explanations within the leptonic models

generally in IC scenarios one can expect  
quite unusual/nonstandard  $X$ - $TeV$  relation



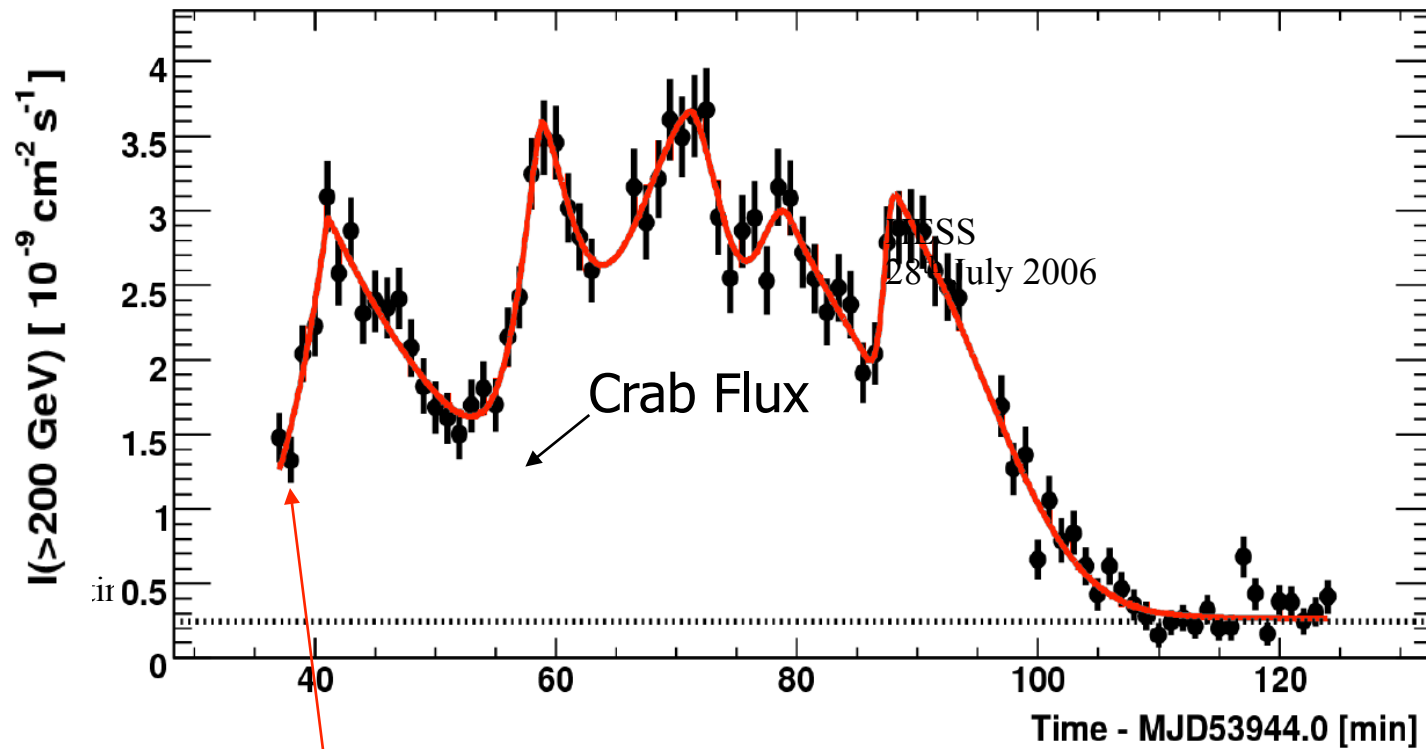
increase of  $TeV$  flux by a factor of 20,  
while  $X$ -ray and optical fluxes have  
been increased only by a factor of  
2 and 15%, respectively

$TeV$ - $X$  cubic dependence against  
the simple SSC model

## *most exciting results of recent years*

- ultra short time variability (on min scales)
- Jet power exceeds Eddington luminosity
- extremely hard energy spectra

several min (200s) variability timescale  $\Rightarrow R=c \Delta t_{\text{var}} \delta_j=10^{14}\delta_{10}$  cm  
 for a  $10^9\text{Mo}$  BH with  $3R_g = 10^{15}$  cm  $\Rightarrow \delta_j > 100$ , i.e. close to the  
 accretion disk (the base of the jet), the bulk motion  $\Gamma > 50$



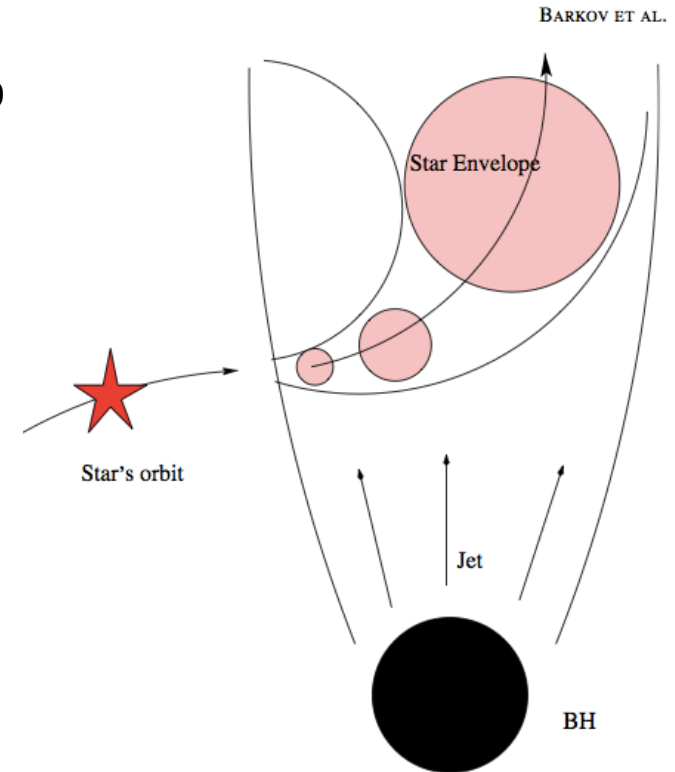
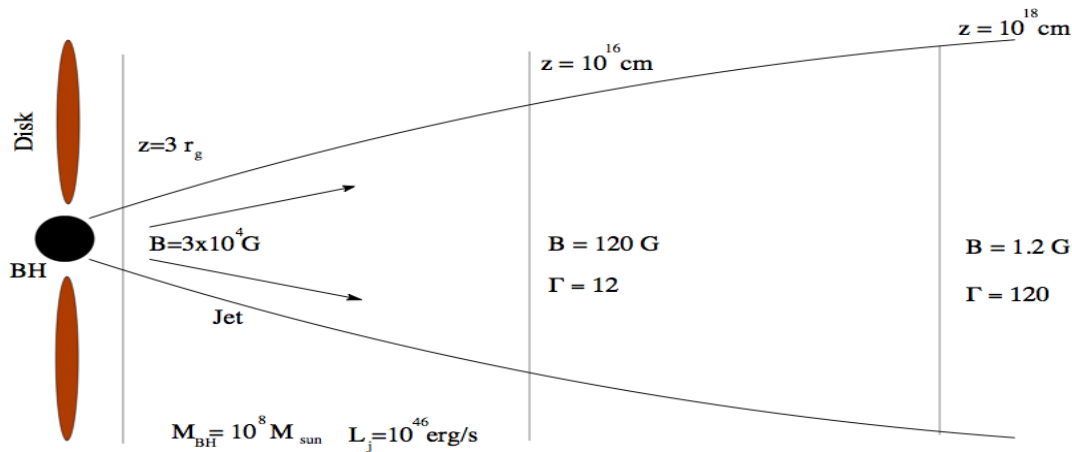
rise time <200s

## on the Doppler boosting and mass of BH in PKS2155-309

- several min variability timescale  $\Rightarrow R = ct_{\text{var}} \delta_j \sim 10^{13} \delta_j$  cm for a  $10^9 \text{Mo}$  BH with  $3R_g \sim 10^{15}$  cm  $\Rightarrow \delta_j > 100$ , i.e. close to accretion disk (the base of the jet), the Lorentz factor of the jet  $\Gamma > 50$  - this hardly can be realized close to  $R_g$ !
- the (internal) shock scenario: shock would develop at  $R = R_g \Gamma^2$ , i.e. minimum  $\gamma$ -ray variability would be  $R_g/c = 10^4 (M/10^9 \text{Mo})$  sec, although  $\gamma$ -ray production region is located at  $R_g \sim ct_{\text{var}} \Gamma^2$  (e.g. Chelotti et al. 1998) - this is true for any other scenario with a “signal-perturbation” originating from the central BH
- thus for the observed  $t_{\text{var}} < 200$  s, the mass of BH cannot significantly exceed  $10^7 \text{Mo}$ . On the other hand the “BH mass–host galaxy bulge luminosity” relation for PKS2155-304 gives  $M > 10^9 \text{Mo}$ .

*Solution?* perturbations are caused by external sources, e.g. by magnetized condensations (“blobs”) that do not have direct links to the central BH; do we deal with the scenario “**star crosses the relativistic e<sup>+</sup>e<sup>-</sup> jet**” ?

# “star crosses the relativistic $e^+e^-$ jet”



B-field: very large or very small?

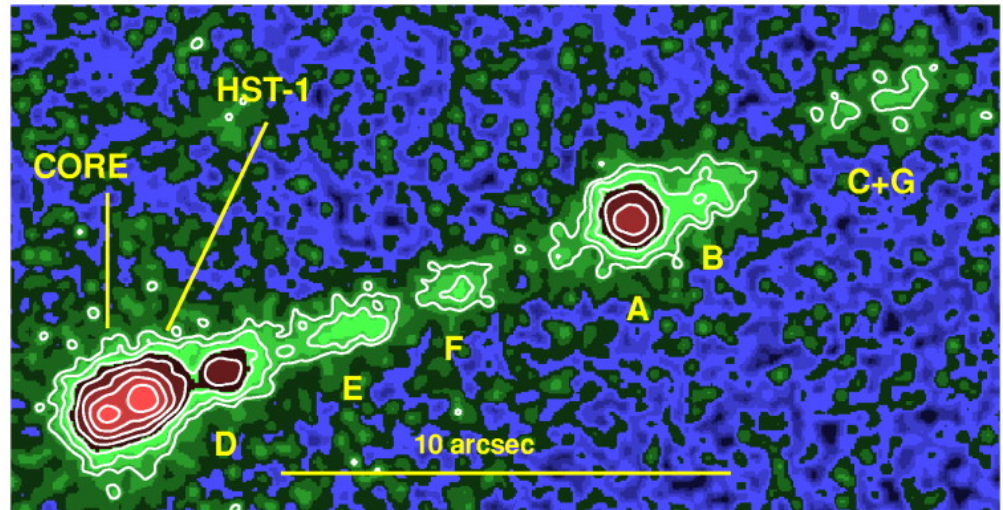
in powerful blazars at subparsec scales B-field cannot be smaller than 1G, a serious constraint for the simplified one-zone “leptonic models,

## M 87 – evidence for production of TeV gamma-rays close to BH ?

- Distance:  $\sim 16$  Mpc
- central BH:  $3 \times 10^9 M_{\odot}$  \*)
- Jet angle:  $\sim 30^\circ$   
=> *not a blazar!*

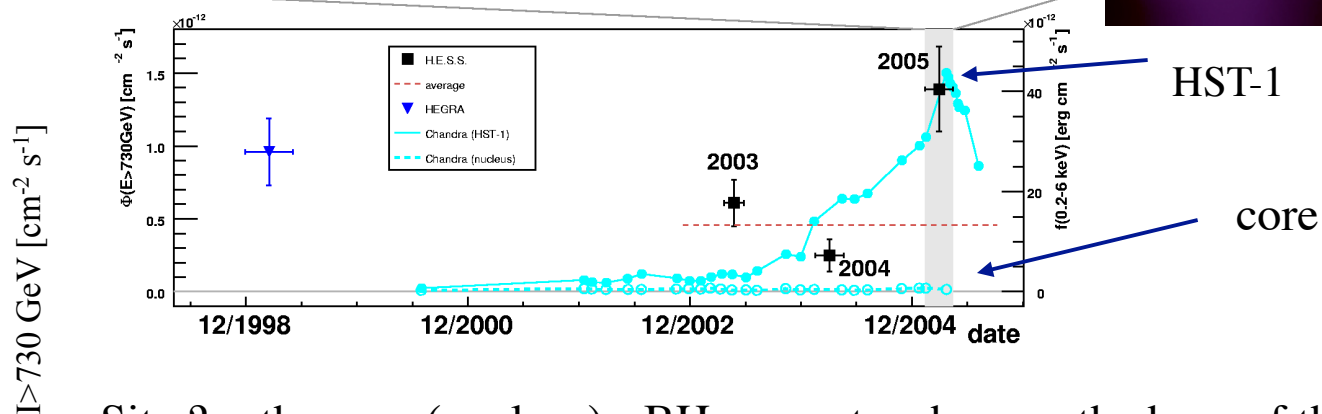
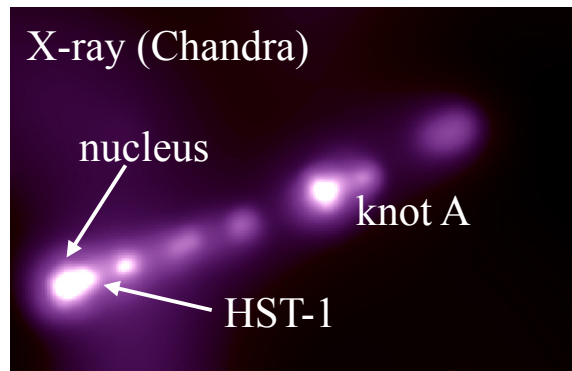
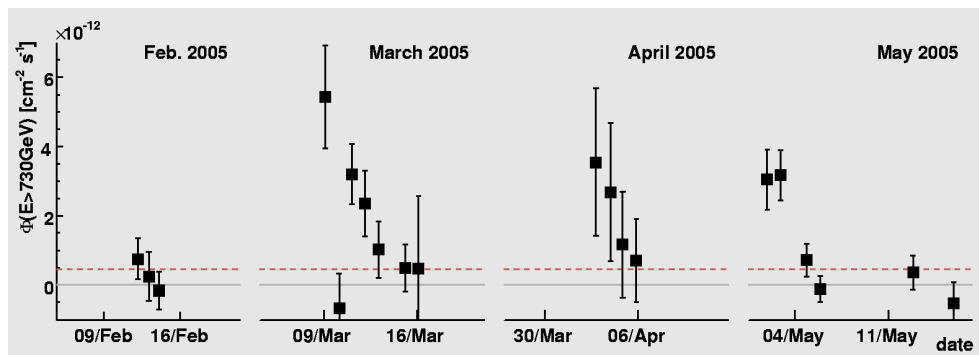
discovery ( $>4\sigma$ ) of TeV  $\gamma$ -rays  
by [HEGRA](#) (1998) and confirmed  
recently by [HESS/VERITAS](#), [MAGIC](#)

\*) recently  $6.4 \times 10^9 M_{\odot}$   
arXiv: 0906.1492 (2009)



M87: light curve and variability

HESS Collaboration 2006, Science, 314,1427



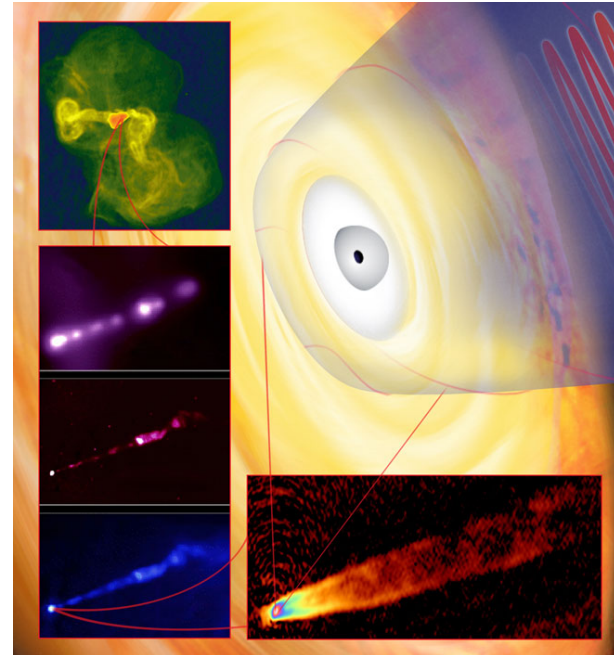
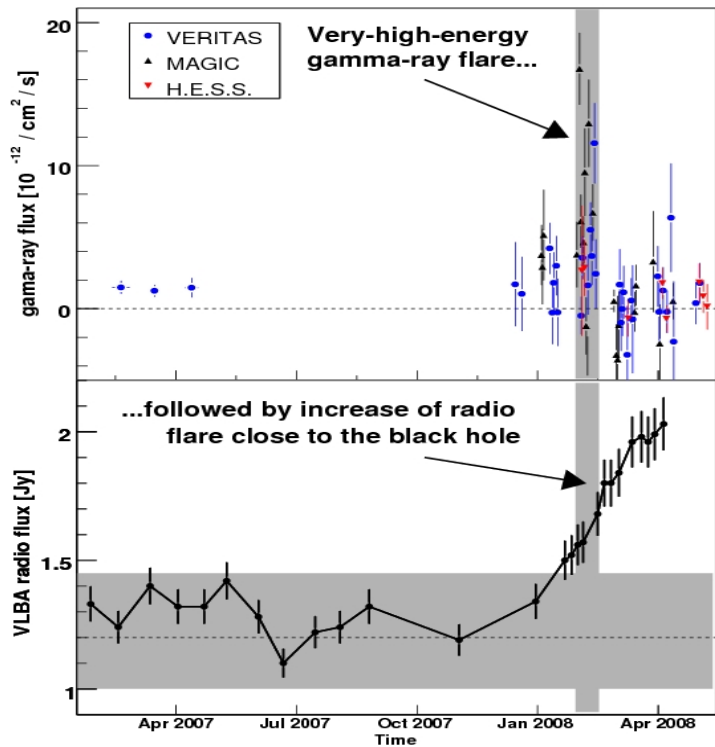
Site ? - the core (nucleus) - BH magnetosphere or the base of the jet

short-term variability on 1-2 day scales => emission region  $R \sim 5 \times 10^{15} \delta_j \text{ cm}$   
 => production of gamma-rays very close to the 'event horizon' of BH?

*because of very low luminosity of the core in O/IR:  
 TeV gamma-rays can escape the production region*

$$L_{\text{IR}} \approx 10^{-8} L_{\text{Edd}}$$

**New!** NRAO and VERITAS/MAGIC/HESS: *Science*, July 2, 2009  
Simultaneous TeV and radio observations allow localization of  
gamma-ray production region within  $50 R_s$



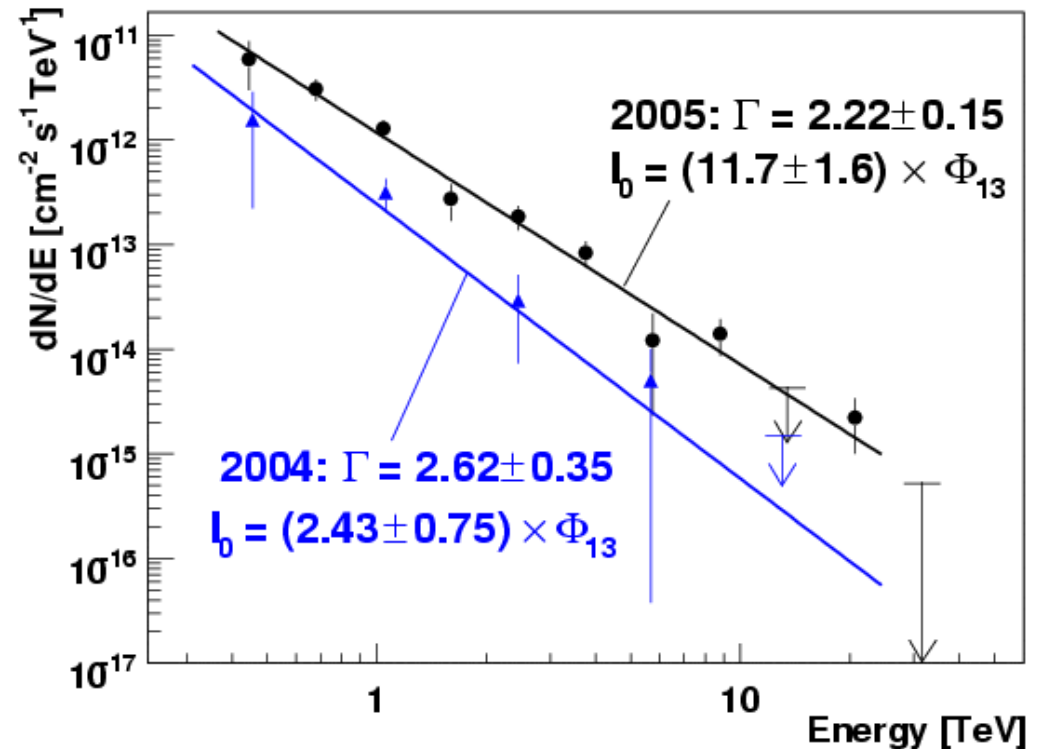
monitoring of the M87 inner jet with VLBA at 43 GHz (ang. res.  $0.21 \times 0.43$  mas) revealed increase of the radio flux by 30 to 50% correlated with the increase in TeV gamma-ray flux in Feb 2008

**conclusion?** *TeV gamma-rays are produced in the jet collimation region within  $50 R_s$  around BH*

# energy spectra

energy spectra for 2004 ( $\sim 5\sigma$ )  
and 2005 ( $\sim 10\sigma$ )

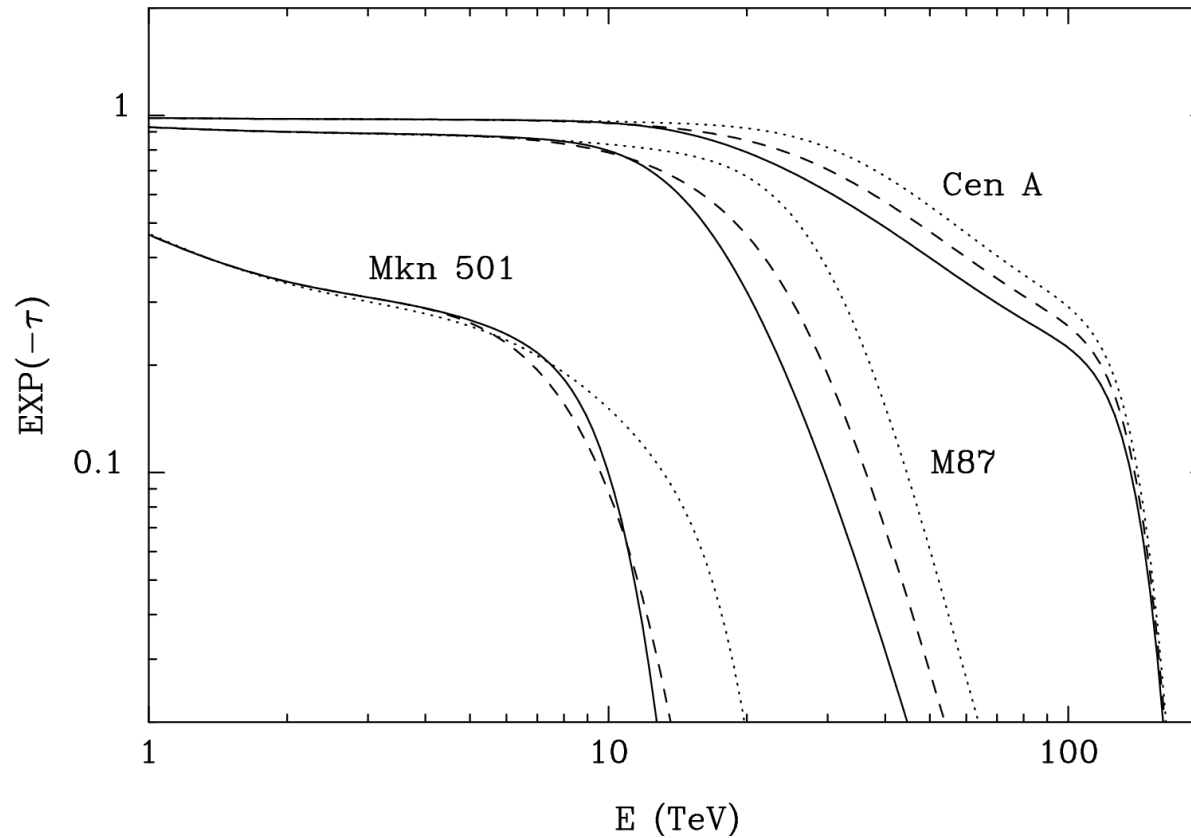
Differential spectra well  
described by power-laws:



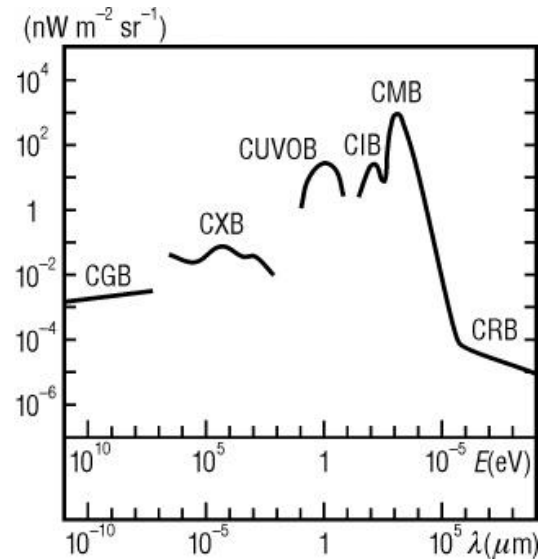
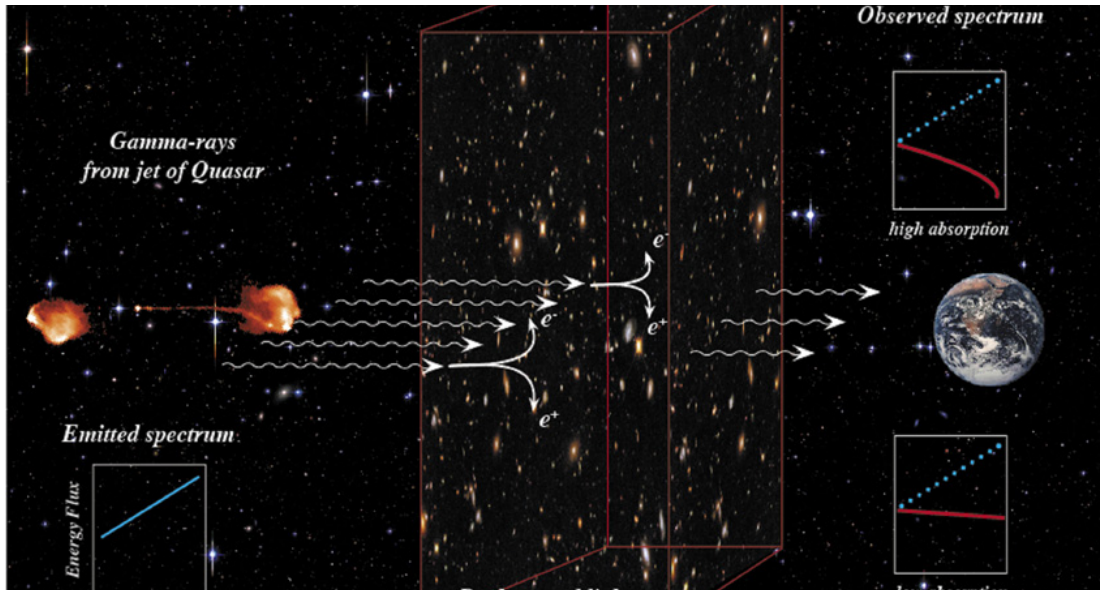
$$\Phi_{13} = 10^{-13} \text{ cm}^{-2} \text{ s}^{-1} \text{ TeV}^{-1}$$

2004 vs. 2005:  
Photon indices compatible, but different flux levels

# Probing DEBRA at MIR /FIR with $E_\gamma > 10$ TeV $\gamma$ -rays from nearby extragalactic sources ( $d < 100$ Mpc)



# *gamma-ray blazars and EBL*



energy-dependent gamma-ray absorption  $\Rightarrow$  information about EBL

cascades in CMB/EBL  $\Rightarrow$  total high energy luminosity of the Universe

$\Rightarrow$  EBL and total energy in relativistic electrons

$\Rightarrow$  evidence of extremely low IGMFs

# Blazars and EBL - two different but tightly connected topics

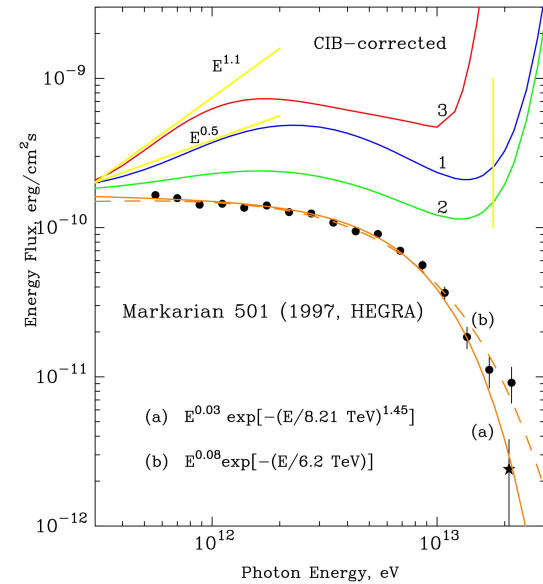
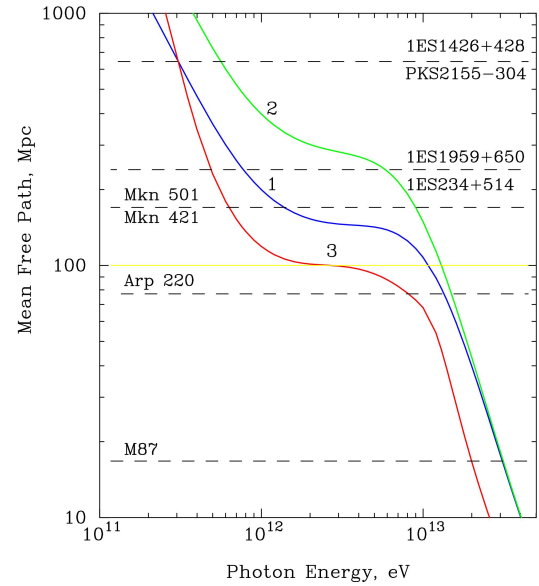
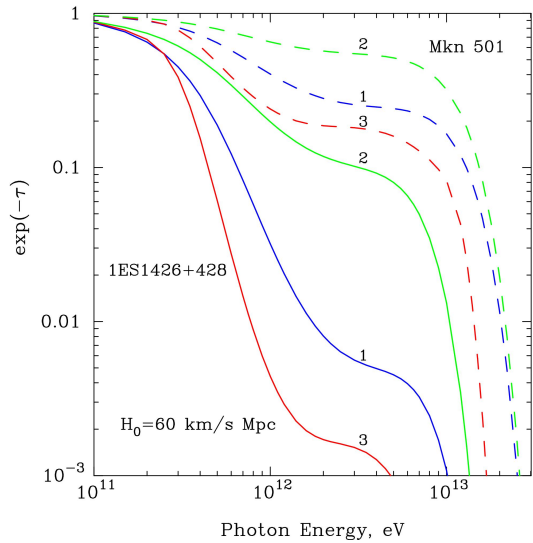
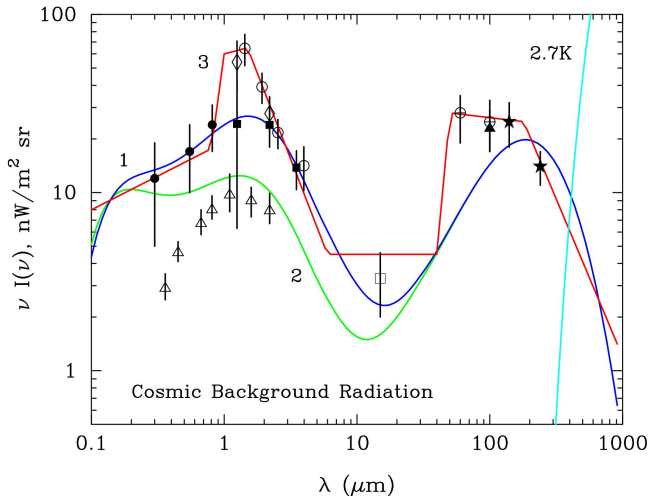
GeV/TeV gamma-ray observations:  
strong impact on

- ✓ Blazar physics and astrophysics
- ✓ Diffuse Extragalactic Background (EBL)  
Intergalactic Magnetic fields (IGMF)

most exciting results of recent 2-3 years

- ◆ *variability on 2-3 min timescales*
- ◆ *jet power exceeds Eddington luminosity?*
- ◆ *unusually hard gamma-ray spectra*

# intergalactic absorption of gamma-rays



unusuall gamma-ray spectra? two options:

claim that EBL is “detected” between O/NIR and MIR

or

propose *extreme* hypotheses, e.g.

violation of Lorentz invariance, non-cosmological origin of z ...

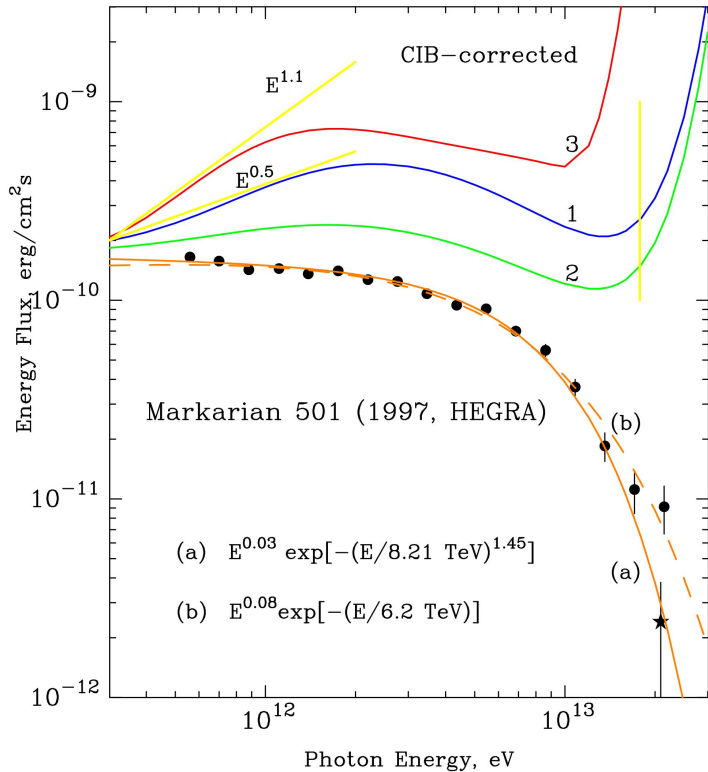
or

propose *less dramatic (more reasonable)* ideas, e.g.

- very specific spectrum of electrons  $\rightarrow \nu F_\nu \sim E_\gamma^{1.33}$
- TeV emission from blazars due to comptonization of cold relativistic winds with bulk Lorentz factor  $\Gamma > 10^6$
- internal gamma-ray absorption

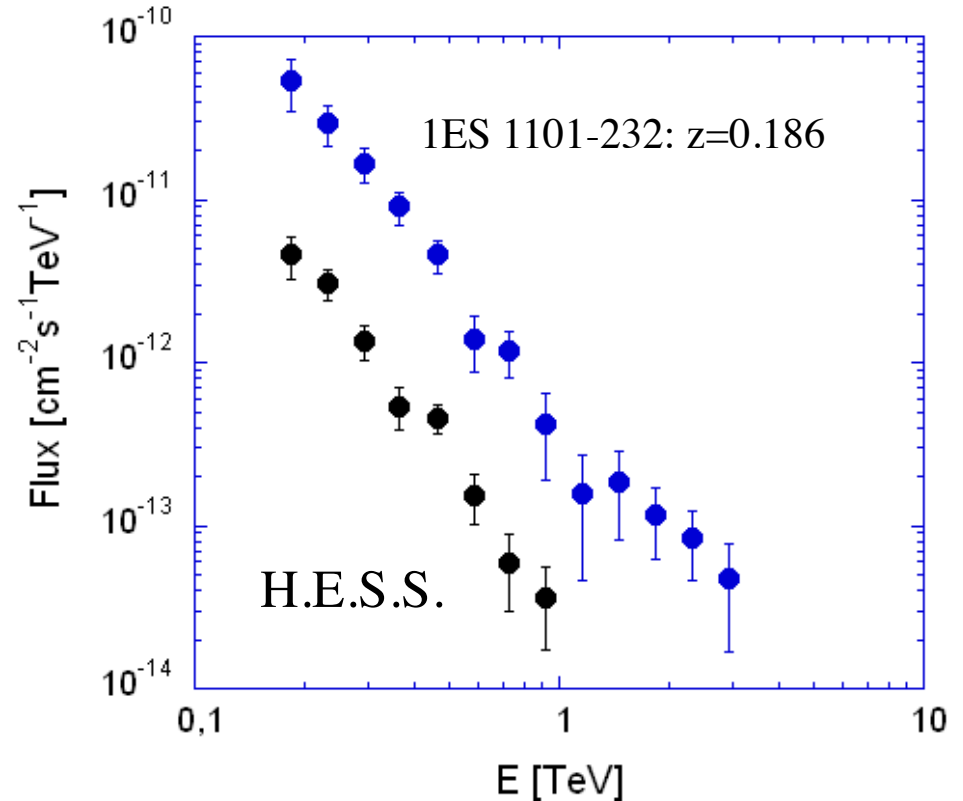
# Blazars and EBL

Mkn 501:  $z=0.031$ : an “infrared crisis”, but with a happy end...



reported EBL flux at FIR  
have not been confirmed

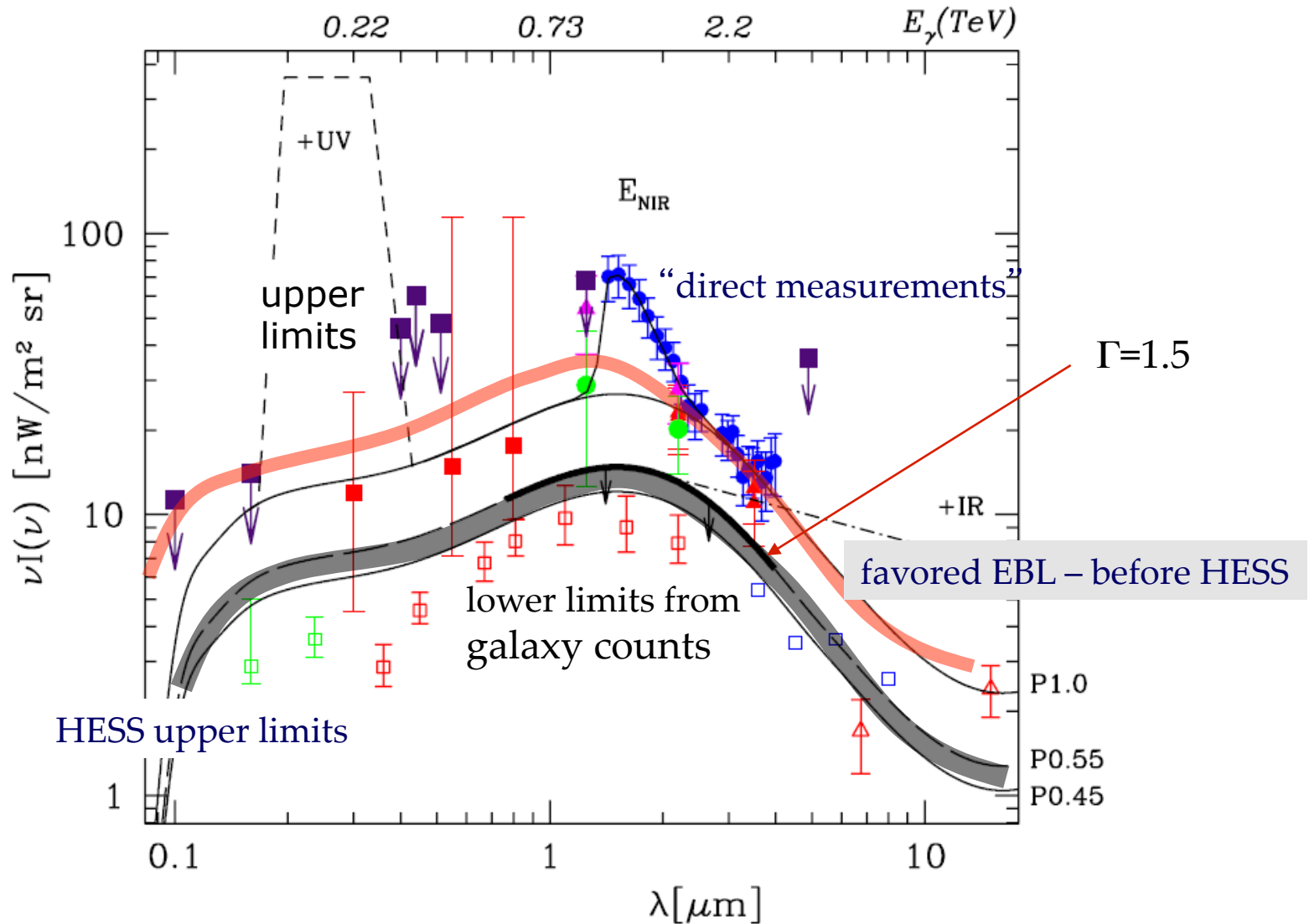
TeV blazars detected by HESS at  $z > 0.15$  !



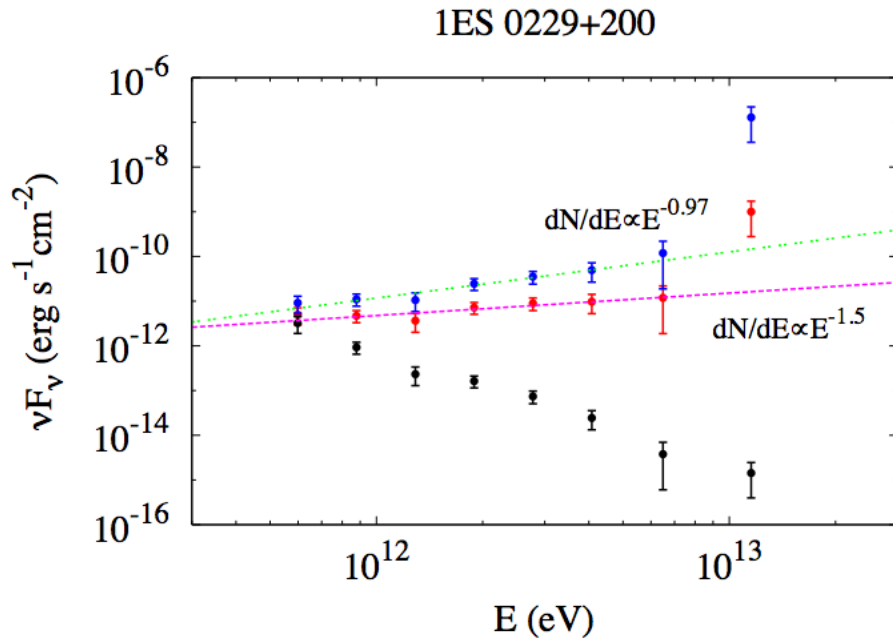
corrected for EBL absorption  
 $\gamma$ -ray spectrum not harder  
than  $E^{-\Gamma}$  ( $\Gamma=1.5$ )  $\Rightarrow$  **u.l. EBL**

# HESS upper limits on EBL - good agreement with recent EBL studies

EBL (almost) resolved at NIR ?



*new “problematic” sources*

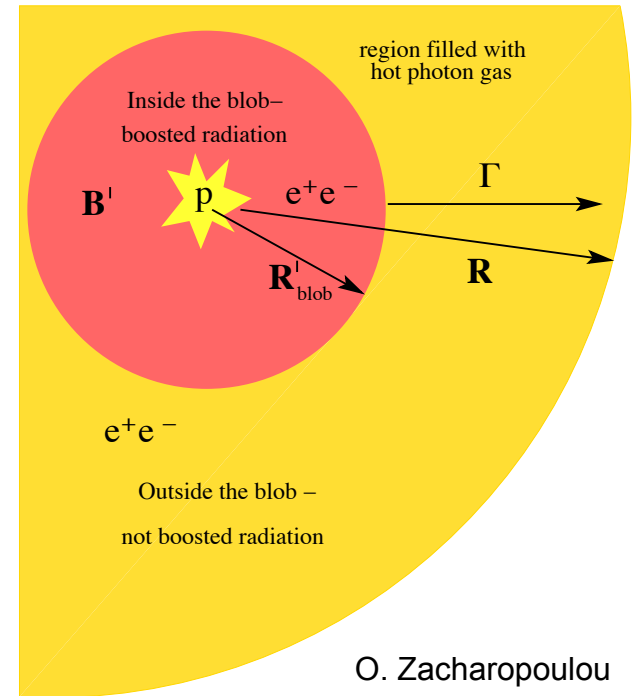
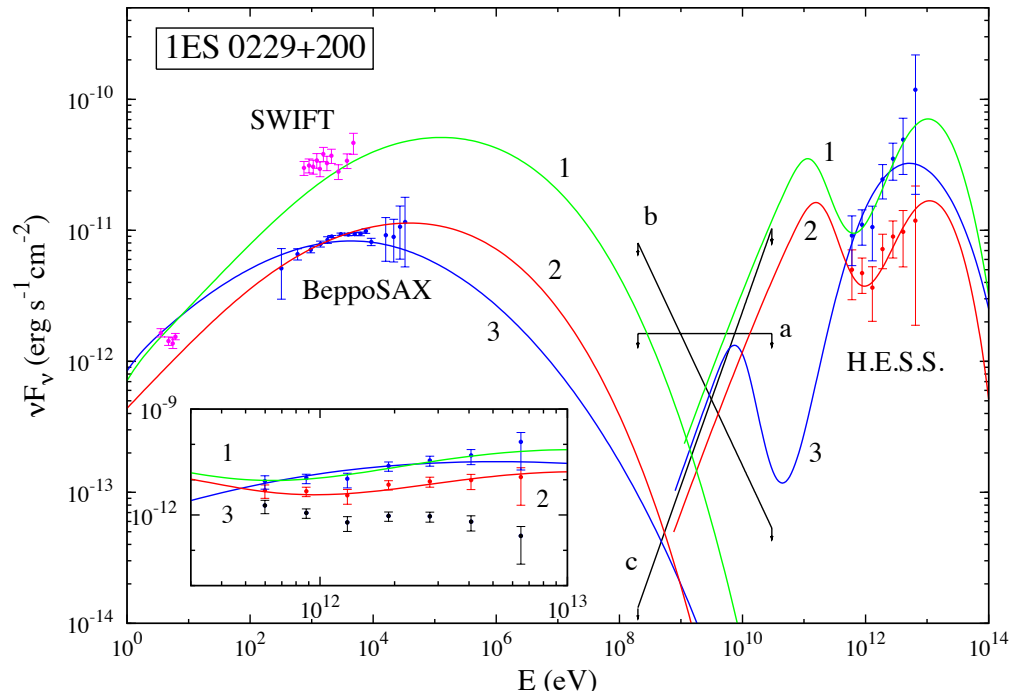


*possible explanations:*

- ✓ **very narrow electron distribution** - no significant radiative energy losses => typically very small B-field: 0.001G introduce adiabatic losses or assume stochastic (Fermi II type) acceleration with Maxwellian type distribution
- ✓ **internal  $\gamma$ - $\gamma$  absorption** => very strong magnetic field,  $B > 10$  G mechanism: proton synchrotron

1ES 0229+200:  $z = 0.14$  spectrum extends to  $> 5$  TeV ! even slight deviation from the “standard” EBL => extremely hard  $\gamma$ -ray spectra with  $\Gamma \sim 1$

# Proton synchrotron and internal $\gamma$ - $\gamma$ absorption

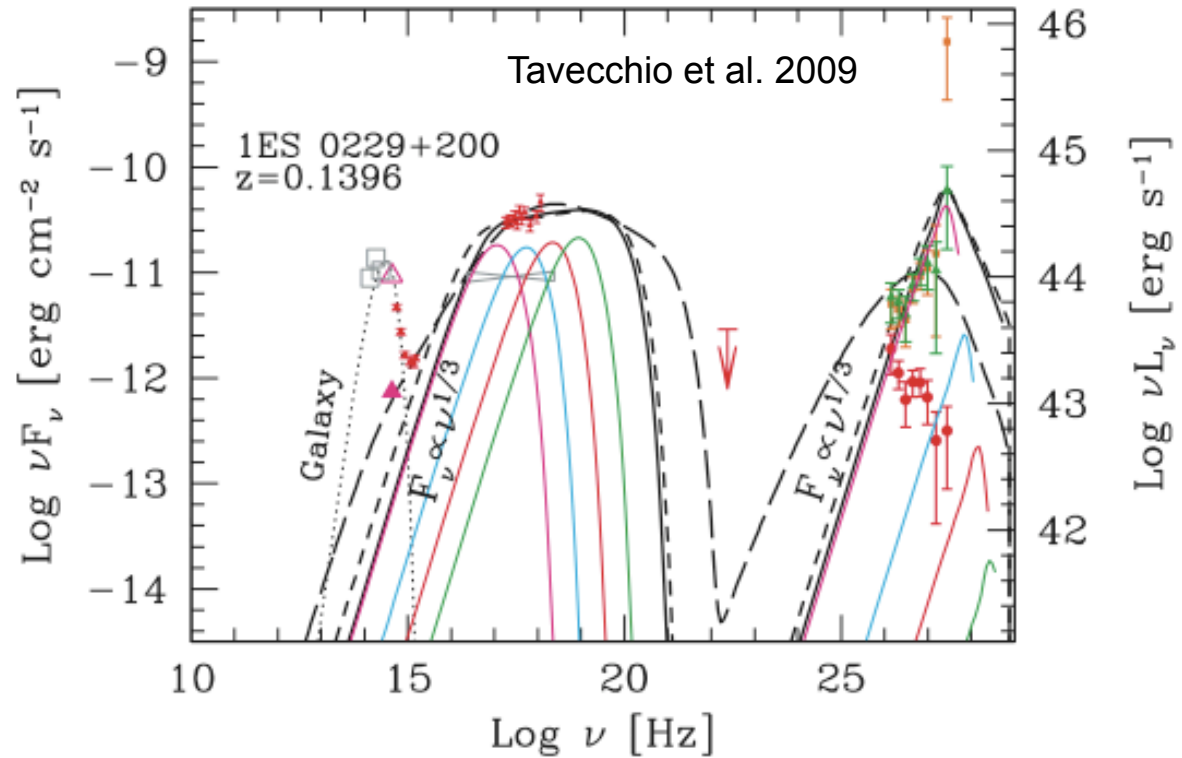


O. Zacharopoulou

very strong magnetic field:  $B > 10 \text{ G} !$

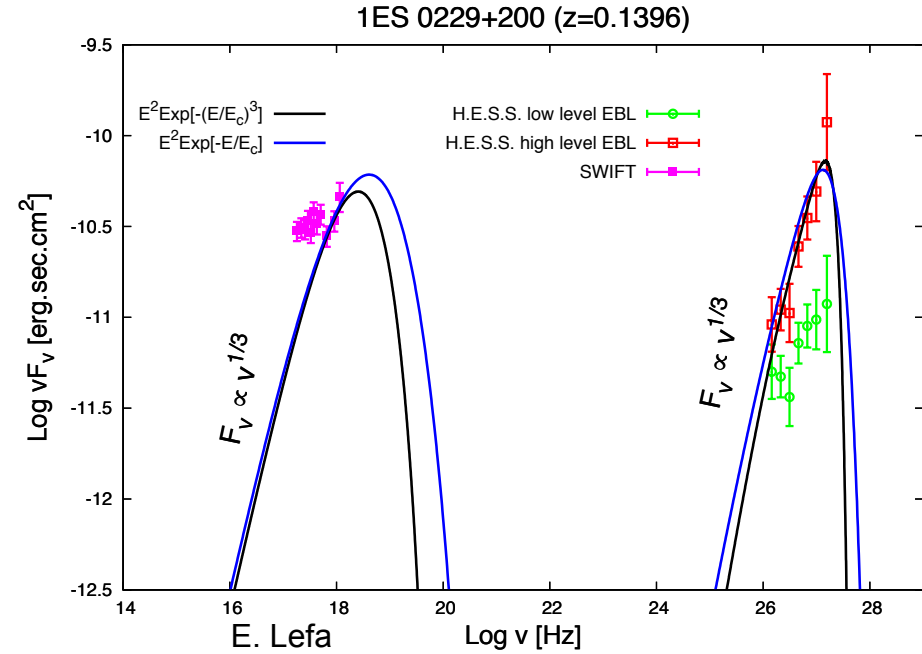
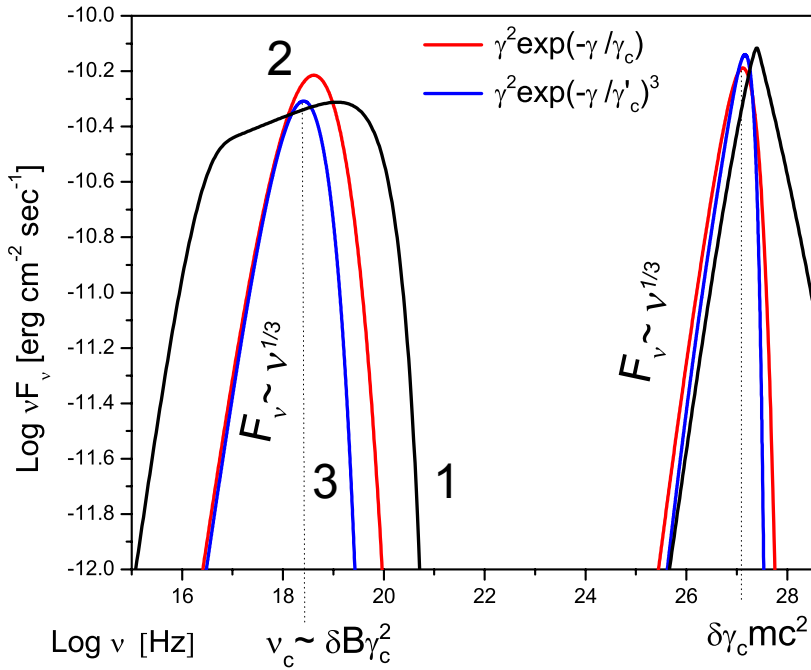
**Table 2.** Input parameters for the models shown in Fig.1 (Model 1: long dashed line; Model 2: solid line; Model 3: short dashed line). See text for definitions.

	$\gamma_{\min}$	$\gamma_b$	$\gamma_{\max}$	$n$	$n_2$	$B$ (G)	$K$ (cm $^{-3}$ )	$R$ (cm)	$\delta$
1	$10^4$	$6 \times 10^5$	$3 \times 10^7$	1.5	3.4	$8.5 \times 10^{-3}$	6	$10^{16}$	50
2	$8.5 \times 10^5$	–	$4 \times 10^7$	2.85	–	$5 \times 10^{-4}$	$3.5 \times 10^9$	$5.4 \times 10^{16}$	30
3	$5 \times 10^5$	–	$4 \times 10^7$	2.85	–	$4 \times 10^{-4}$	$6.7 \times 10^8$	$5.4 \times 10^{16}$	50



$B \sim 10^{-3}$  G: deviation from equipartition by many orders of magnitude!

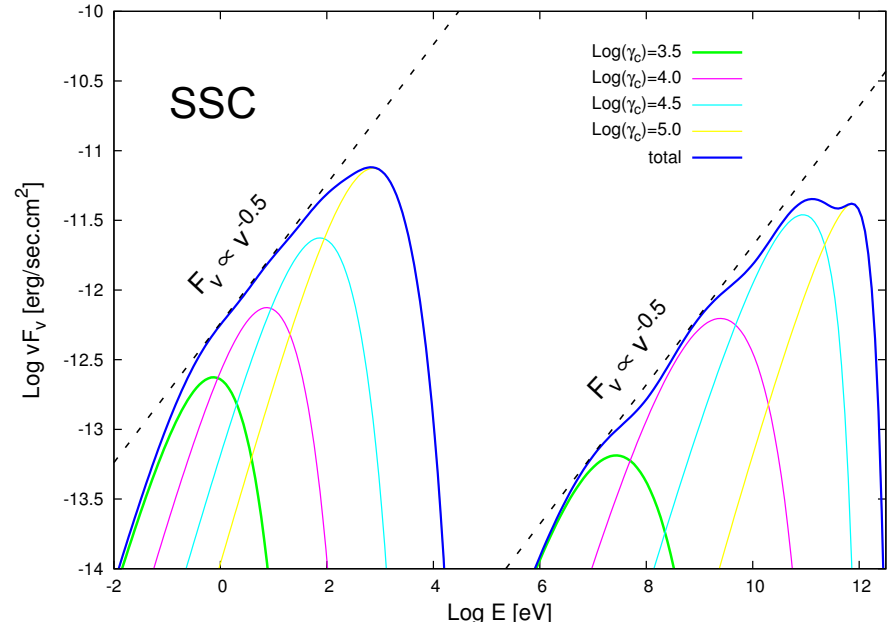
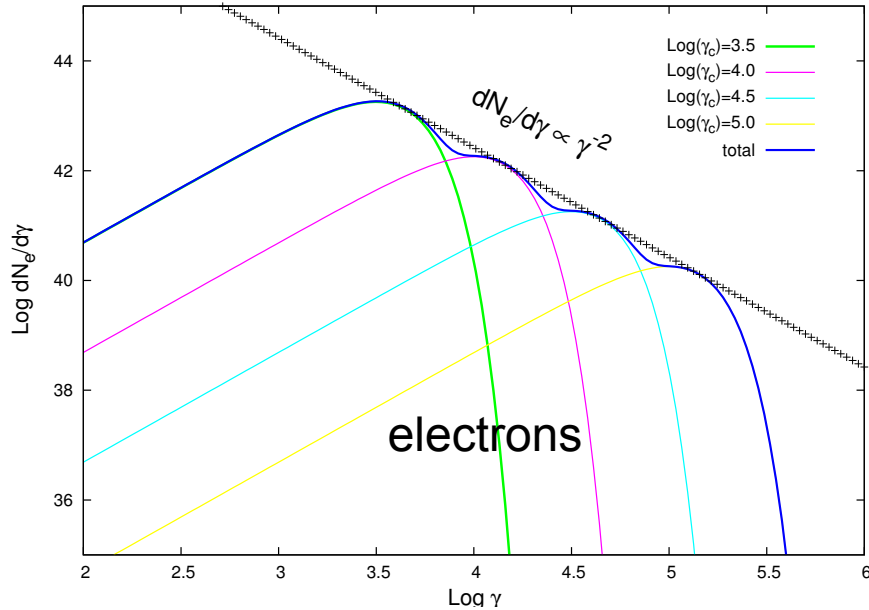
# Synchrotron Self Compton: narrow distribution of electrons



1.  $\gamma_{\min} = 5 \cdot 10^5$ ;  $\gamma_{\max} = 4 \cdot 10^7$ ;  $B = 0.4 \text{ mG}$ ;  $\delta = 50$
2.  $\gamma_c = 1.5 \cdot 10^5$ ;  $B = 70 \text{ mG}$ ;  $\delta = 33$
3.  $\gamma_c = 5.3 \cdot 10^5$ ;  $B = 0.4 \text{ mG}$ ;  $\delta = 33$

small or very small B-field!

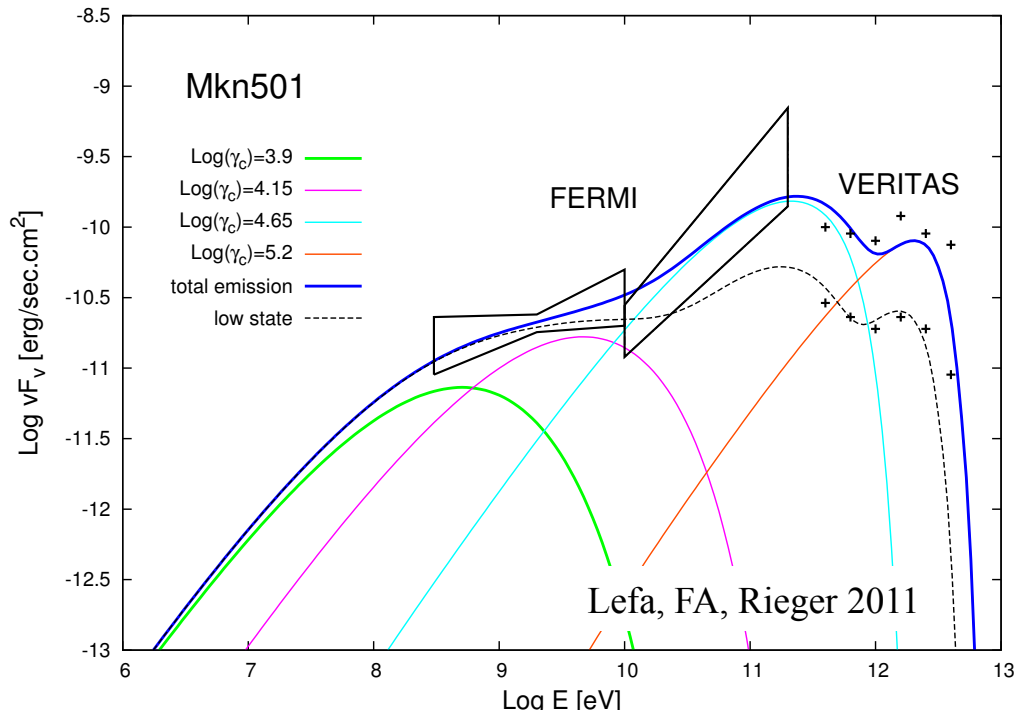
# multi-zone (multi-blob) concept



$$\frac{dN}{d\gamma} = A\gamma^2 \exp[-(\gamma/\gamma_0)^2] \quad E_i = 2 \times 10^{44} \text{ erg}, \quad B = 0.1 \text{ G}, \quad \delta = 30, \quad R = 3 \times 10^{14} \text{ cm}$$

depending on  $E_i$  and  $\gamma_{0i} \Rightarrow$  arbitrary total electron spectrum  
for  $E_i = \text{const}$ , but different  $\gamma_0$  and  $i \gg 1$  almost ideal  $\gamma^{-2}$  spectrum

# Very hard spectrum of Mkn 5011 during 2009 flare



*Fermi* LAT: flat spectrum in a low state and very hard  $dN/dE \sim E^{-1}$  type during 2009 flare (Abdo et al. 2010 and Neronov et al 2011)

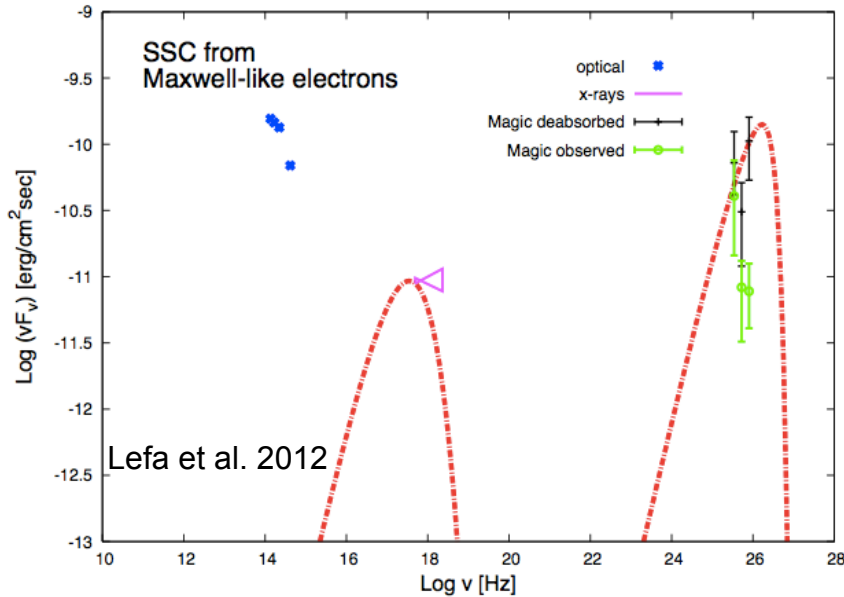
can be explained by change  $\delta=30$  to  $40$  of two “hottest” blobs;  $B=0.1G$ ,  $R=10^{14}$  cm

we can expect g-ray spectrum of arbitrary form; in flaring state as hard as  $E^{-1}$

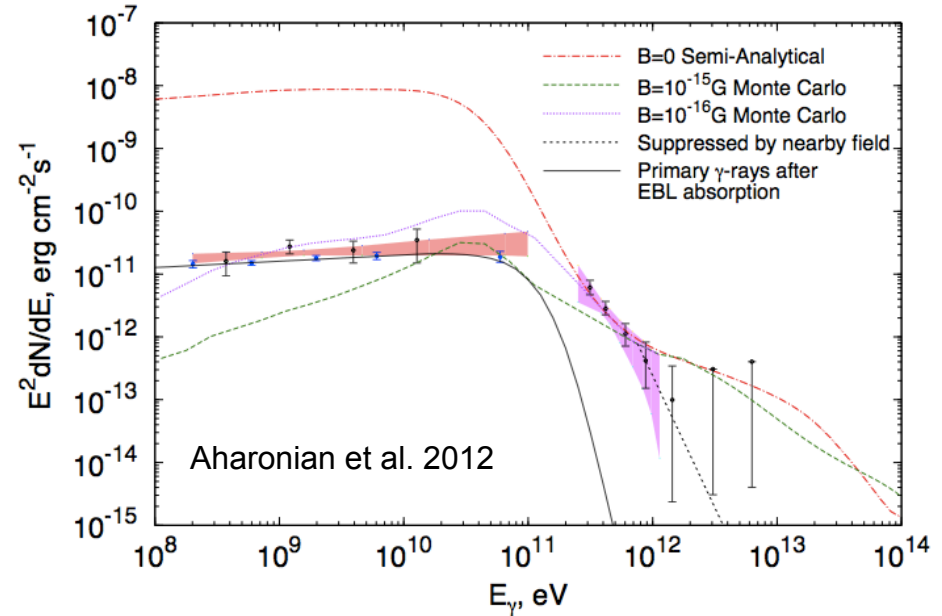
**conclusions:** do not try to get ‘smooth’ spectral fits, especially in low-states  
do not overestimate the potential of “single-zone” models  
do not overestimate the potential of  $\gamma$ -rays for derivation of EBL

# approaching to the threshold of “trouble”?

3C279:  $z=0.536$



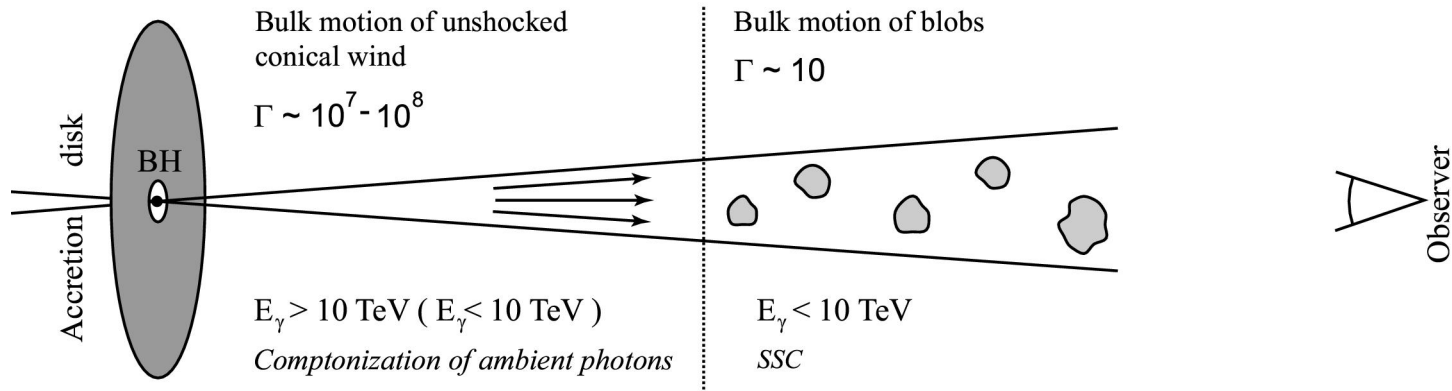
PKS 0447-439:  $z=1.26?$



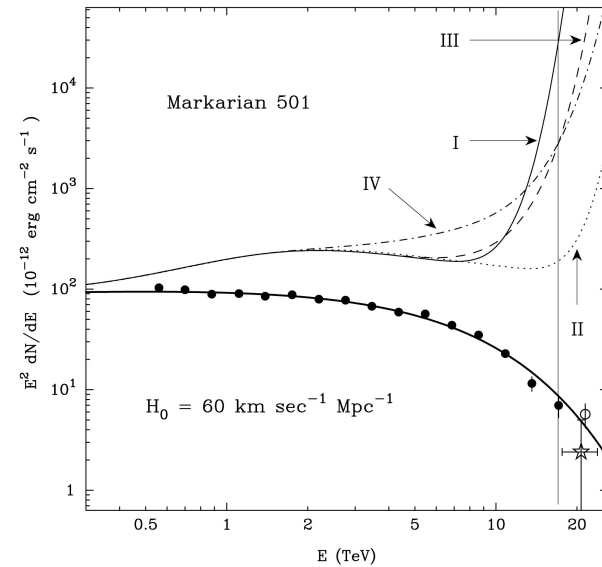
SSC with Maxwellian type electron distribution:  $B=0.3 \text{ G}$ ,  $\delta=40$ ;  $\gamma_0=7 \times 10^4$

possible explanations:  $\gamma$ -rays due to cascades induced by  $>10^{18} \text{ eV}$  protons in IGM with  $B < 10^{-15} \text{ G}$ ? or the claimed redshift is wrong

# Gamma Rays from a cold ultrarelativistic wind ?



in fact not a very exotic scenario ...



# Pair Halos

TeV Gamma-rays from distant extragalactic sources,  $d > 100$  Mpc interact effectively with Extragalactic Background Radiation (EBL; (0.1-100  $\mu\text{m}$ ))

when a gamma-ray is absorbed its energy is not lost !  
absorption in EBL leads to E-M cascades supported by

- Inverse Compton scattering on 2.7 K CMBR photons
- photon-photon pair production on EBL photons

if the intergalactic field is sufficiently strong,  $B > 10^{-11}$  G,  
the cascade  $e^+e^-$  pairs are promptly isotropised

➔ formation of extended structures - Pair Halos

# how it works ?

energy of primary gamma-ray

$$E_{\gamma,0} \simeq 10(E_{\gamma}/100\text{GeV})^{1/2} \text{ TeV}$$

mean free path of parent photons

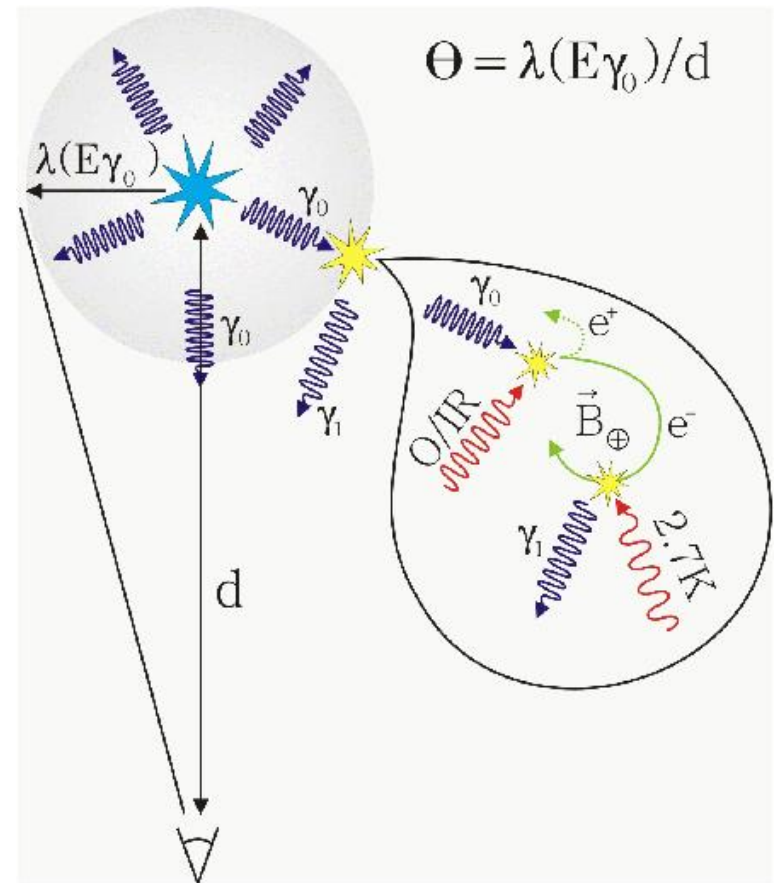
$$\lambda(E_{\gamma,0}) \sim d \times \Theta$$

information about EBL flux at

$$\lambda \simeq 10(E_{\gamma}/100\text{GeV})^{1/2} \mu\text{m}$$

gamma-radiation of pair halos can be recognized by its distinct variation in spectrum and intensity with angle, and depends rather weakly (!) on the features of the central VHE source

two observables – angular and energy distributions allow to disentangle two variables  $u_{\text{EBL}}(\lambda, z)$  and  $d(H_0)$  !



## Pair Halos as Cosmological Candles

- ❑ information about EBL density at fixed cosmological epochs given by the redshift of the central source unique !
- ❑ estimate of the total energy release of AGN during the active phase
- ❑ objects with jets at large angles - many more g-ray emitting AGN

but the advantage of the large Doppler boosting of blazars disappears: beam  $\Rightarrow$  isotropic source

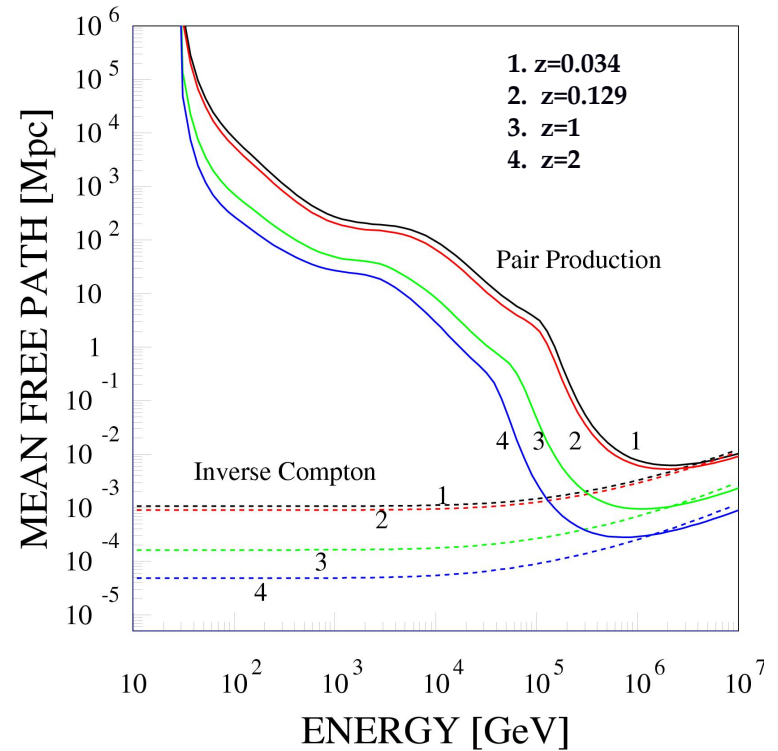
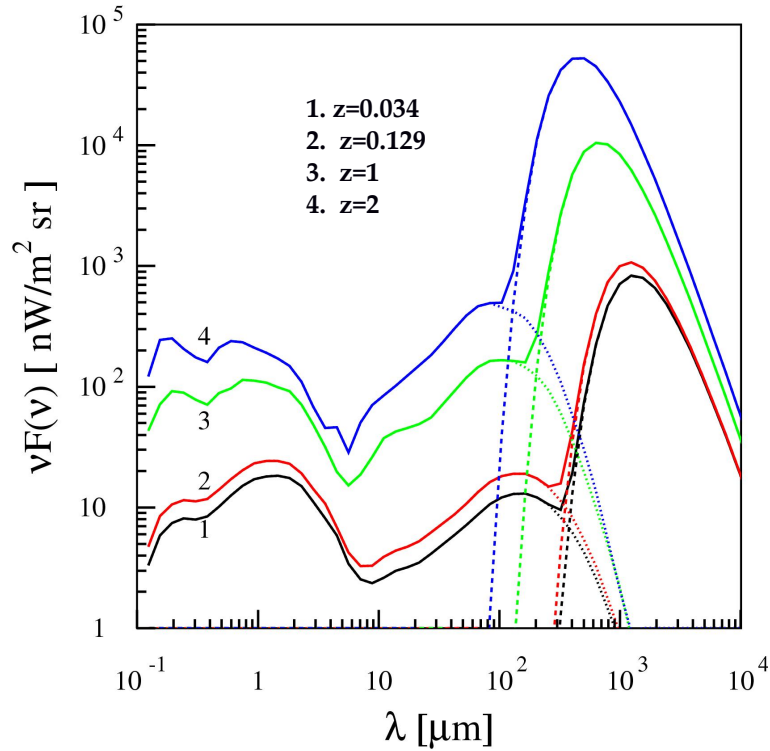
therefore very powerful central objects needed

QSOs and Radiogalaxies (sources of EHE CRS ?)

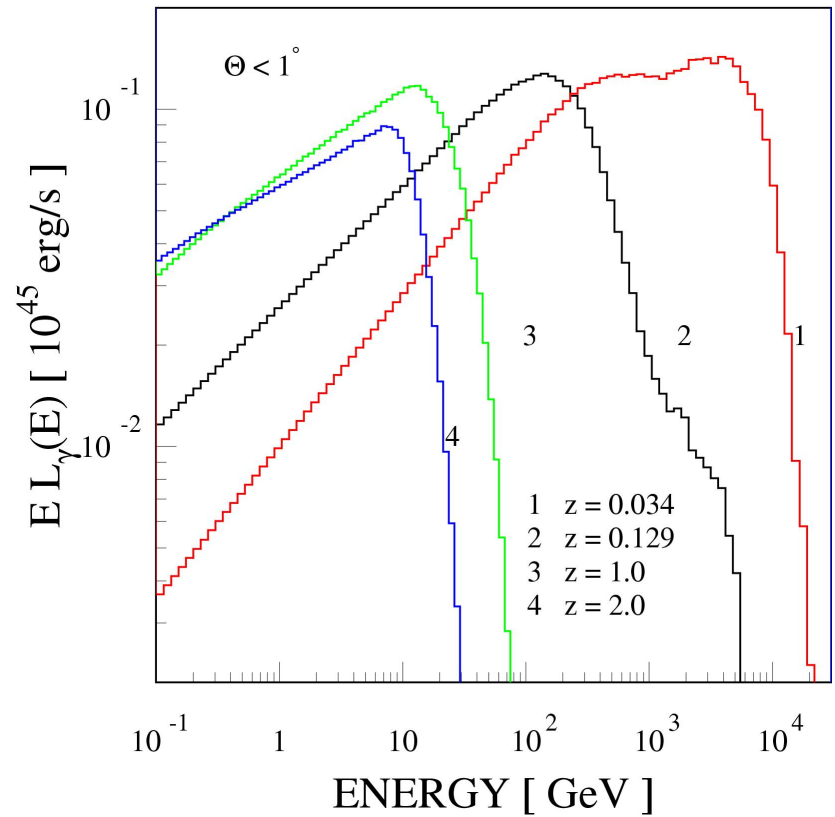
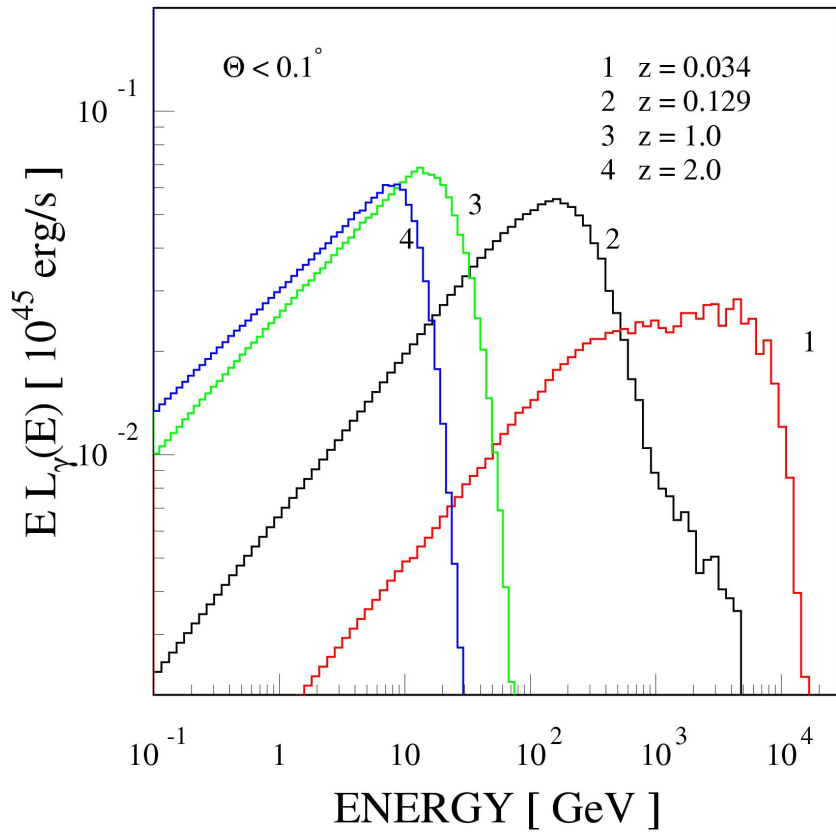
as better candidates for Pair Halos

this requires low-energy threshold detectors

# EBL at different $z$ and corresponding mean freepaths



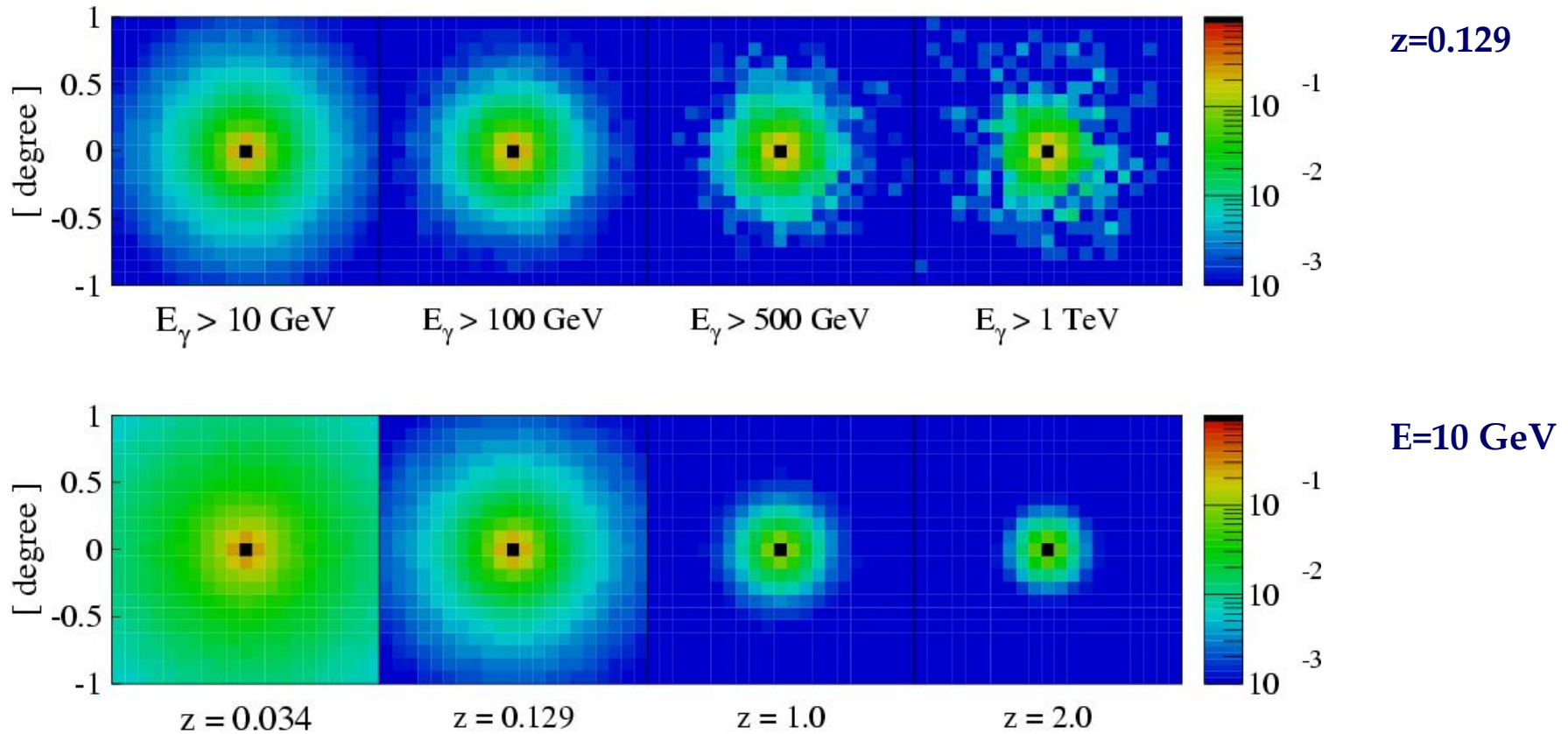
## SEDs for different $z$ within $0.1^\circ$ and $1^\circ$

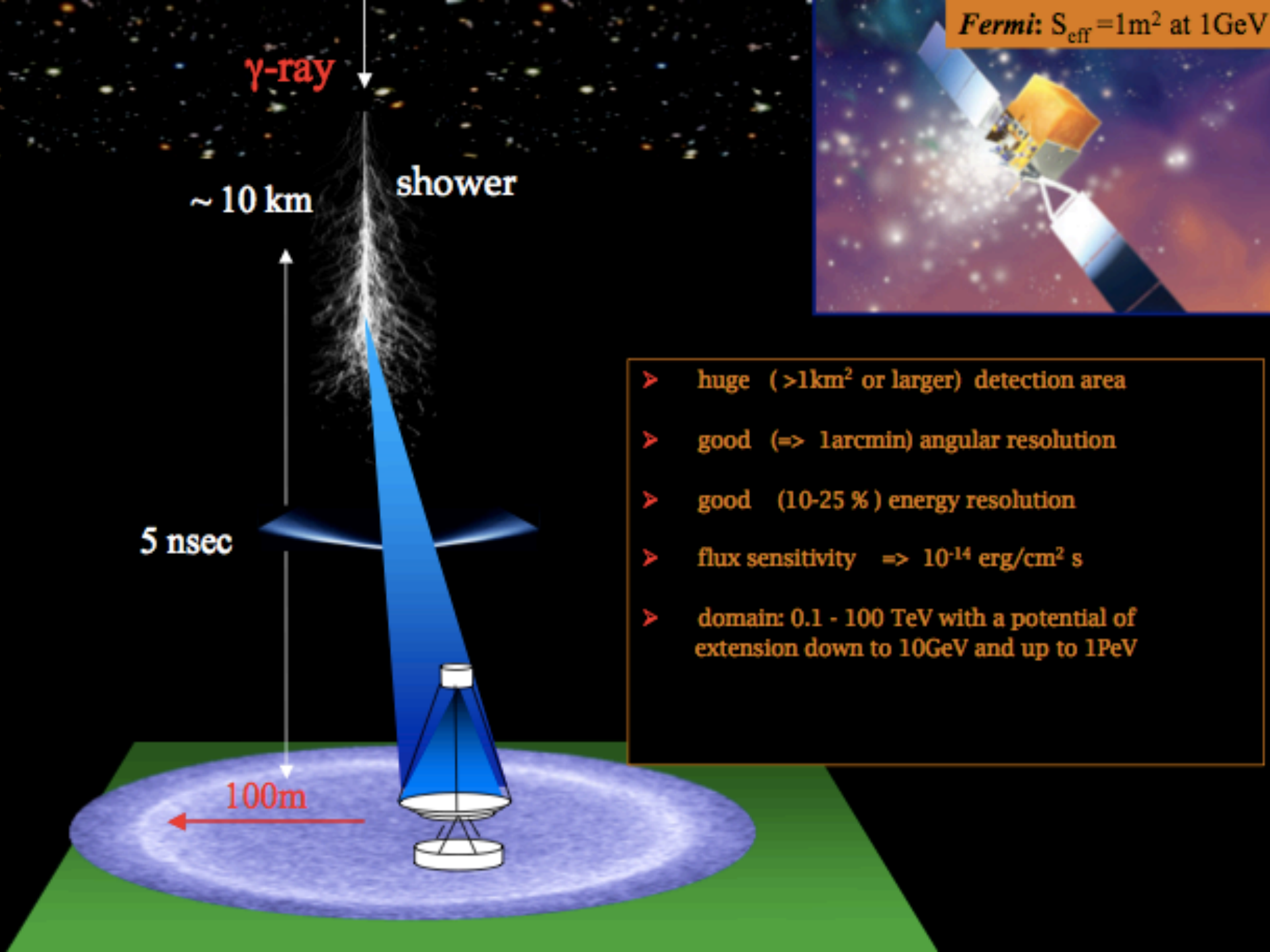


EBL model – Primack et al. 2000

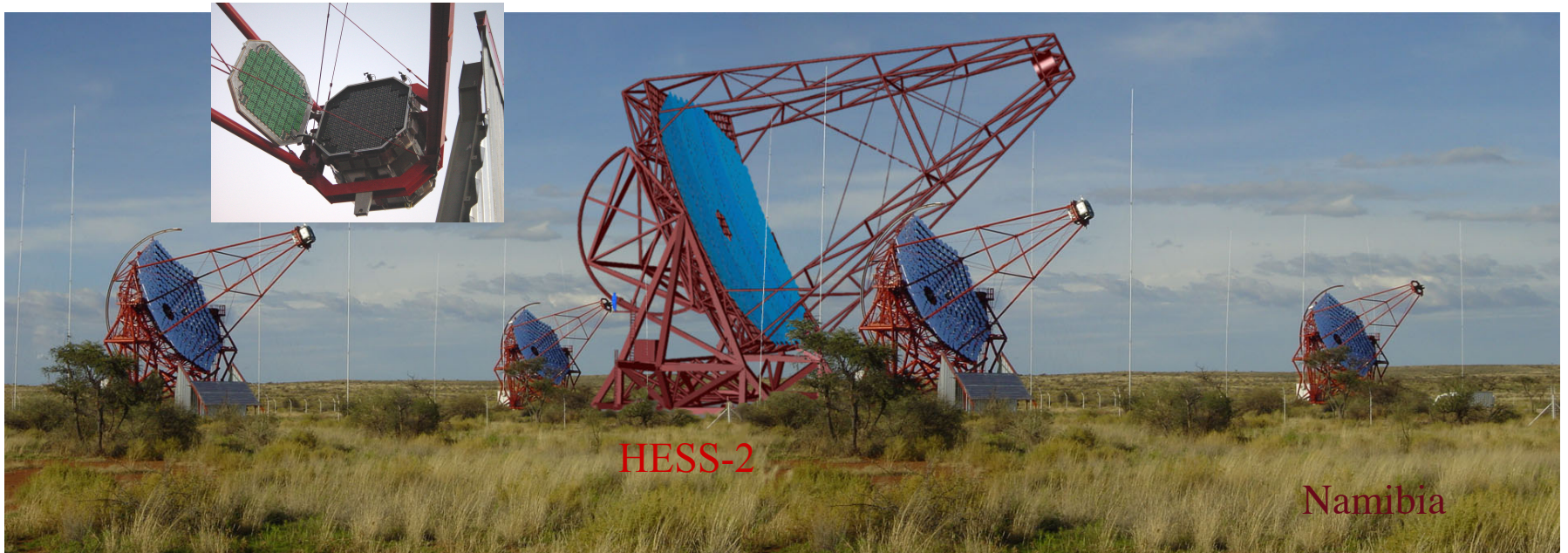
$L_o = 10^{45}$  erg/s

## Brightness distributions of Pair Halos





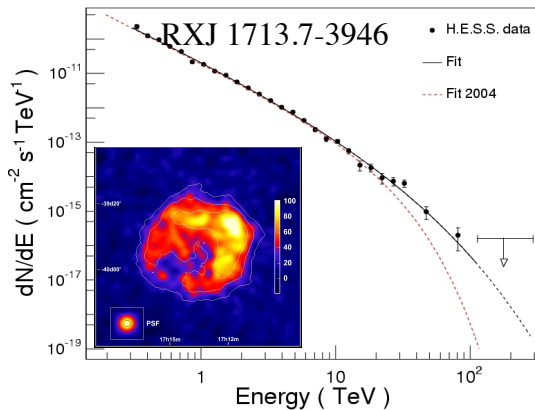
*H.E.S.S. - High Energy Stereoscopic System*



one of the current 3 (HESS, MAGIC, VERITAS) major IACT arrays

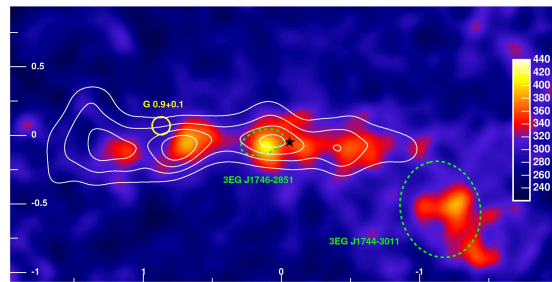
good performance => high quality data => solid basis for physics studies

spectrometry



TeV image and energy spectrum of a SNR

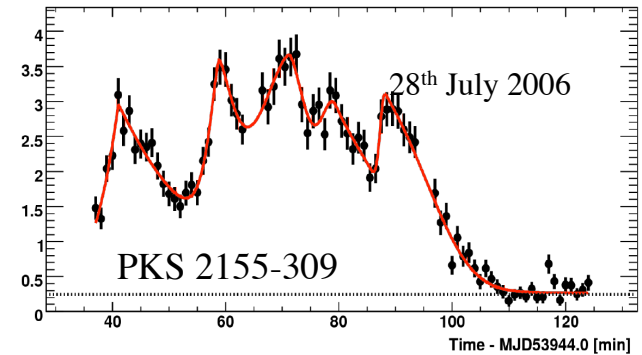
imaging



Galactic Center

resolving GMCs in the Galactic Center 100pc region

timing



variability of TeV flux of a blazar on minute timescales

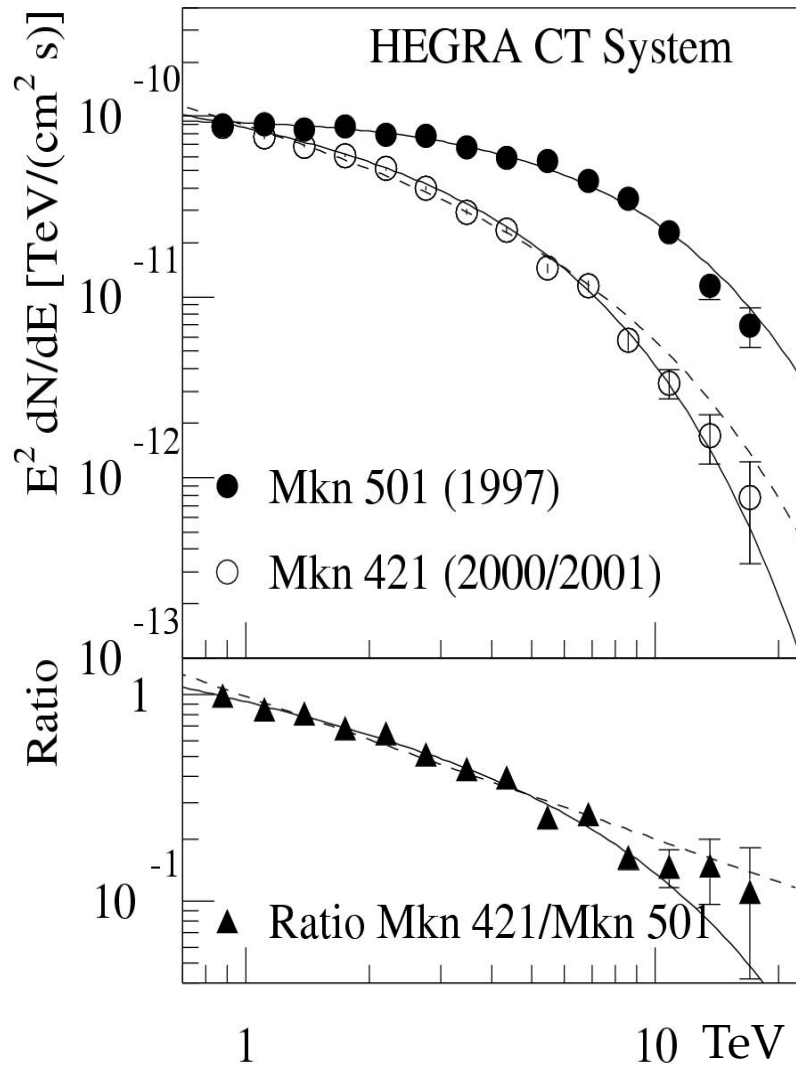
multi-functional tools: *spectrometry temporal studies morphology*

✓ extended sources: *from SNRs to Clusters of Galaxies*

✓ transient phenomena  *$\mu$ QSOs, AGN, GRBs, ...*

*Galactic Astronomy | Extragalactic Astronomy | Observational Cosmology*

# Spectrometry with HEGRA beyond $3E_{\text{cutoff}}$ !



## Unprecedented photon statistics

Mkn 421 - 60,000 TeV photons  
detected in 2001

Mkn 501 - 40,000 TeV photons  
detected in 1997

spectra: canonical power-law  
with exponential cutoff

**Cutoff = 6.2 TeV and 3.8 TeV  
for Mkn 501 and Mkn 421**

time average spectra of  
Mkn 421 and Mkn 501

# THE NEXT BIG STEP: THE CHERENKOV TELESCOPE ARRAY

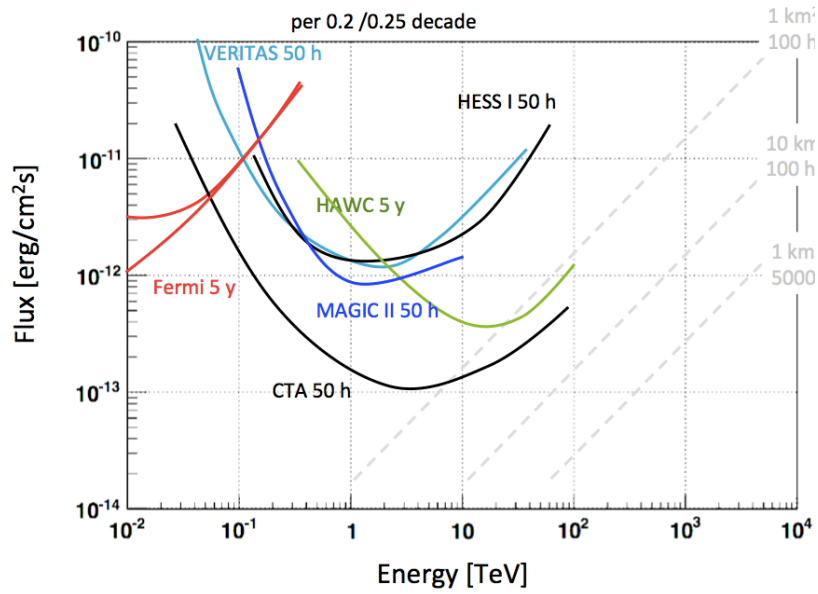
W. Hofmann GAMMA2012

**10 fold improvement in sensitivity**  
**10 fold improvement in usable energy range**  
**much larger field of view**  
**strongly improved angular resolution**

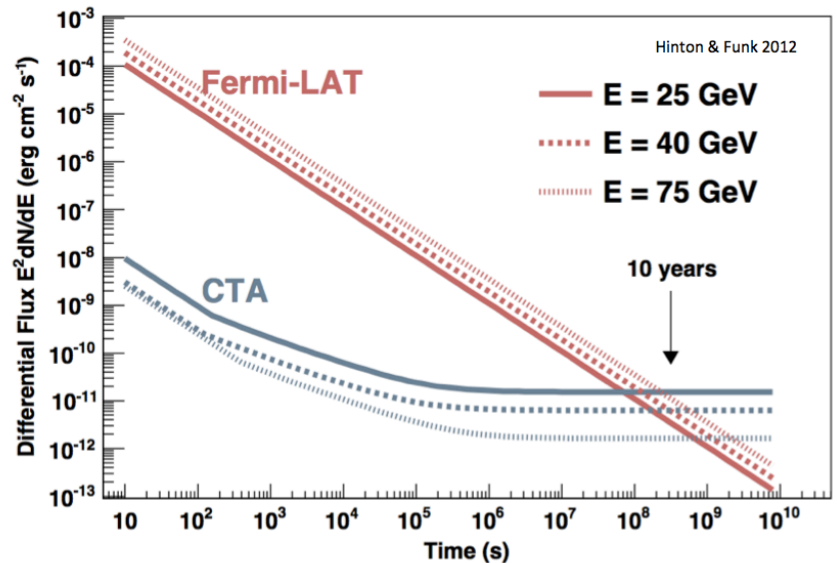


# CTA - a powerful tool for AGN studies

## DIFFERENTIAL SENSITIVITY



## CTA VERSUS FERMI – TRANSIENT SOURCES

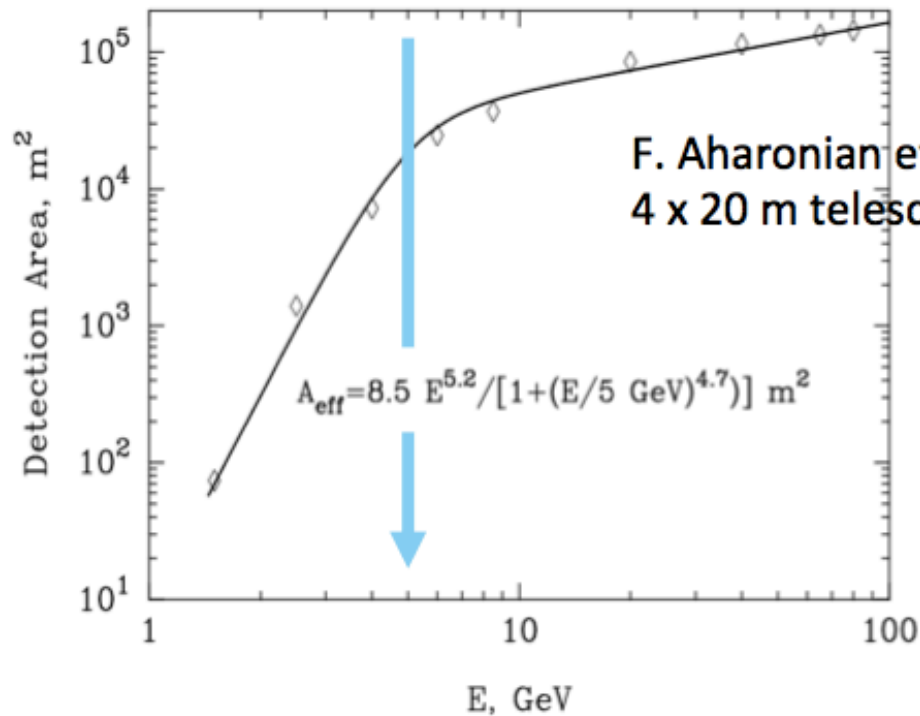


W. Hofmann GAMMA 2012

Detection of 'nominal' (Fermi/AGILE) AGN for just 1 min, but above several tens of GeV the emission could be strongly suppressed at tens of GeV – low threshold is critical

# HIGH-ALTITUDE CHERENKOV TELESCOPES

10-12 km



higher light intensity (5000 m: x2)

→ lower threshold

smaller light pool area (5000 m: /2)

Reference height ~2000 m

## *current status of AGN studies in gamma-rays*

almost one thousand GeV and several dozens TeV  $\gamma$ -ray emitting AGN!

`GeV-to-TeV' ratio is not too large given the continuous monitoring of AGN by Fermi LAT and AGILE of a significant fraction of the sky

GeV - large *source* statistics - good for population studies

TeV - large *photon* statistics - good for physics studies

the field is not yet saturated:

- ✓ *Fermi/AGILE and HESS/MAGIC/VERITAS will bring more results*
- ✓ *CTA will elevate the status of the field to a new level*
- ✓ *10 GeV (or so) threshold IACT arrays - new surprises and challenges*