

Elliptic deformations of quantum Virasoro and W_n algebras

Work in collaboration with L. Frappat and E. Ragoucy (LAPTH Annecy)
Extension of work by J.A, L.F., M. Rossi, P. Sorba, 1997-99

References:

Deformed Virasoro algebras from elliptic quantum algebras
Jean Avan, Luc Frappat, Eric Ragoucy
[arXiv: 1607-05050](https://arxiv.org/abs/1607.05050)

Plan :

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I-1 INTRODUCTION: THE FIRST DEFORMED VIRASORO ALGEBRA

Reference : [J. Shiraishi, H.Kubo, H. Awata, S. Odake ; Lett. Math. Phys. 38 \(1996\), 33](#)

Notion discussed in e.g. [Curtright-Zachos 1990](#). More precisely derived from construction of Virasoro Poisson algebra on extended center of affine A_1 algebra at critical center $c=-2$

Extended center of affine q -deformed A_1 algebra again at $c=-2$: [Reshetikhin Semenov-Tjan-Shanskii 1990](#) ; Poisson structure : [Frenkel Reshetikhin 1996](#) ; Quantized by [SKAO](#) and [Feigin-Frenkel 1996](#) ;

Extension to q - W_n algebras by same authors and [Awata-Kubo-Odake-Shiraishi 1996](#) : extended center at $c=-n$ for affine q -deformed $sl(n)$ algebra.

Proposed structure :

Algebra generated by the generating functional $T(z)$ satisfying the relation

$$f(z/w)T(z)T(w) - T(w)T(z)f(z/w) = -\frac{(1-q)(1-t^{-1})}{1-p} \left[\delta\left(\frac{pw}{z}\right) - \delta\left(\frac{w}{pz}\right) \right]$$

Structure function is given by

$$f(x) = \exp \left(\sum_{n=1}^{\infty} \frac{(1-q^n)(1-t^{-n})}{1+p^n} \frac{x^n}{n} \right)$$

$t=p/q$

Warning : notations here are « old reckoning » : p, q (AKOS) = q^2, p^{-1} (our soon-to-come parameters)

Classical limit $\ln p = \ln q$; Virasoro limit $p, q \rightarrow 1 + o(\epsilon)$, $\ln p / \ln q = b$; $T(z) = 2 + \epsilon^2 (t(z) + \dots)$

Occurs as e.g.

Symmetries for restricted SOS models ([Lukyanov 1996](#)) ;

Natural operators acting on eigenvectors for Ruijsenaar Schneider models ; hence connection with Macdonald and Koornwinder polynomials ([SKAO 1996](#)) ;

Natural algebraic structure for partner of 5D gauge field theory in extension of AGT conjecture (Awata-Yamada 2010, Nieri 2015). Hence connection also to q -Painlevé : **See new work :**

[Bershtein-Shchekkin 1608.02566](#)

Quantization by construction of vertex operators using current algebra construction of Virasoro/ W_n and q -deforming it ([SKAO, FF](#)).

Alternative : Direct embedding into larger algebraic structure ? Started in previous papers [AFRS 97-99](#).

General idea : [SKAO](#) formula has 2 parameters p, q and constituent blocks of elliptic functions : Ratio of structure functions $f(x)/f(x^{-1})$ is ratio of elliptic Jacobi Theta functions.

Hence suggests quantization of classical DVA naturally inserted into *elliptic affine algebra* instead of quantum affine algebra.

Consider elliptic $gl(N)$ algebra with generic N . A priori leads to deformation of W . But restrict here to spin 1 generator, or $N=2$. Extension to full W_N in project. Partially realized in [AFRS 97-99](#).
Notion of quantum « powers » to be refined.

I -2 THE ELLIPTIC ALGEBRA $A_{qp}(gl(N)_c)$

Original proposition for $gl(2)$ by [Foda-Iohara-Jimbo-Kedem-Miwa-Yan 1994](#). Goes to $sl(2)$ by factoring out **q-determinant**. To be commented later on.

Extended to $gl(N)$ by [Jimbo-Kono-Otake-Shiraishi 1999](#). Justification of quasi Hopf structure in [Arnaudon-Buffenoir-Ragoucy-Roche 1998](#) = identification of Drinfel'd twist. $sl(N)$??? q-determinant ???

Lax matrix encapsulates generators as :

$$L(z) = \begin{pmatrix} L_{11}(z) & \cdots & L_{1N}(z) \\ \vdots & & \vdots \\ L_{N1}(z) & \cdots & L_{NN}(z) \end{pmatrix}.$$

Exchange relation

$$\widehat{R}_{12}(z/w) L_1(z) L_2(w) = L_2(w) L_1(z) \widehat{R}_{12}^*(z/w),$$

R-matrix ([Baxter 1981](#), [Chudnovskii-Chudnovskii 1981](#)), in term of Jacobi theta functions with rational characteristics. g, h related with periods of elliptic functions ;

$$\mathcal{Z}(z, q, p) = z^{2/N-2} \frac{1}{\kappa(z^2)} \frac{\vartheta \left[\begin{smallmatrix} \frac{1}{2} \\ \frac{1}{2} \end{smallmatrix} \right] (\zeta, \tau)}{\vartheta \left[\begin{smallmatrix} \frac{1}{2} \\ \frac{1}{2} \end{smallmatrix} \right] (\xi + \zeta, \tau)} \sum_{(\alpha_1, \alpha_2) \in \mathbb{Z}_N \times \mathbb{Z}_N} W_{(\alpha_1, \alpha_2)}(\xi, \zeta, \tau) I_{(\alpha_1, \alpha_2)} \otimes I_{(\alpha_1, \alpha_2)}^{-1}$$

$$z = e^{i\pi\xi}, \quad q = e^{i\pi\zeta}, \quad p = e^{2i\pi\tau}.$$

$$W_{(\alpha_1, \alpha_2)}(\xi, \zeta, \tau) = \frac{\vartheta \left[\begin{smallmatrix} \frac{1}{2} + \alpha_1/N \\ \frac{1}{2} + \alpha_2/N \end{smallmatrix} \right] (\xi + \zeta/N, \tau)}{N \vartheta \left[\begin{smallmatrix} \frac{1}{2} + \alpha_1/N \\ \frac{1}{2} + \alpha_2/N \end{smallmatrix} \right] (\zeta/N, \tau)}$$

$$R(z, q, p) = (g^{\frac{1}{2}} \otimes g^{\frac{1}{2}}) \mathcal{Z}(z, q, p) (g^{-\frac{1}{2}} \otimes g^{-\frac{1}{2}}); \quad g_{ij} = \omega^i \delta_{ij}$$

$$\widehat{R}_{12}(z) \equiv \widehat{R}_{12}(z, q, p) = \tau_N(q^{\frac{1}{2}} z^{-1}) R_{12}(z, q, p), \quad \tau_N(z) = z^{\frac{2}{N}-2} \frac{\Theta_{q^{2N}}(qz^2)}{\Theta_{q^{2N}}(qz^{-2})}.$$

$$\widehat{R}_{12}^*(z, q, p) = \widehat{R}_{12}(z, q, p^* = pq^{-2c})$$

$\Theta_p(z) = (z; p)_\infty (pz^{-1}; p)_\infty (p; p)_\infty$ Here arises *quasi-Hopf* structure (recall e.g. dynamical quantum algebra).

Warning : R is unitary solution to Yang-Baxter equation, but **does not enter** into definition of elliptic algebra from quasi-Hopf structure. (Drinfel'd - twisted R matrix is non-unitary one! See [Arnaudon-Buffenoir-Ragoucy-Roche 1998](#)). Seems not relevant but ...

Relation with (untwisted) quantum group structure obtained by redefining :

$$\begin{aligned} L^+(z) &= L(q^{\frac{c}{2}} z), \\ L^-(z) &= (g^{\frac{1}{2}} h g^{\frac{1}{2}}) L(-p^{\frac{1}{2}} z) (g^{\frac{1}{2}} h g^{\frac{1}{2}})^{-1}. \quad \text{and reads :} \\ \widehat{R}_{12}(z/w) L_1^\pm(z) L_2^\pm(w) &= L_2^\pm(w) L_1^\pm(z) \widehat{R}_{12}^*(z/w), \\ \widehat{R}_{12}(q^{\frac{c}{2}} z/w) L_1^+(z) L_2^-(w) &= L_2^-(w) L_1^+(z) \widehat{R}_{12}^*(q^{-\frac{c}{2}} z/w). \end{aligned}$$

Use in this formulation of non-unitary R matrix instead of R **NOW** leads to different structure due to different normalization. (to be kept in mind).

I -3 : THE QUANTUM DETERMINANT FOR N=2

FOR N = 2 Elliptic R-matrix evaluated at $-1/q$ degenerates to **I +P₁₂** (permutation operator). Equal to antisymmetrizer . Hence possible to define q-determinant :

$$\begin{aligned} \mathfrak{q}\text{-det } L(z) &= L_{11}(q^{-1}z)L_{22}(z) - L_{21}(q^{-1}z)L_{12}(z) \\ &= L_{22}(q^{-1}z)L_{11}(z) - L_{12}(q^{-1}z)L_{21}(z) \\ &= L_{11}(z)L_{22}(q^{-1}z) - L_{12}(z)L_{21}(q^{-1}z) \\ &= L_{22}(z)L_{11}(q^{-1}z) - L_{21}(z)L_{12}(q^{-1}z) \end{aligned}$$

Lies in center of quantum elliptic algebra. Allows to define inverse $\widehat{L}(q^{-1}z)L(z) = \mathfrak{q}\text{-det } L(z) \mathbb{I}_2$.

and comatrix :

$$\widehat{L}(z) = \begin{pmatrix} L_{22}(z) & -L_{12}(z) \\ -L_{21}(z) & L_{11}(z) \end{pmatrix}$$

I -4 THE STRATEGY

First step : define quadratic functional of Lax matrix, and conditions on parameters p,q,c such that quadratic functional close exchange algebra : **CLOSURE CONDITIONS**

Second step : define second (analytical) set of conditions on p,q,c such that exchange algebra become abelian : **ABELIANITY CONDITIONS**.

Then expand around set of conditions to get **Poisson structure**. Closure algebra then automatically yields quantization of Poisson structure.

Does one get classical DVA Poisson ? Does it quantize to DVA ?

II-1 THE CLOSURE RELATIONS

General result :

Theorem 3.1 *In the three-dimensional parameter space generated by q, p, c , consider the two-dimensional surfaces \mathcal{S}_{mn} defined by the relation $(-p^{\frac{1}{2}})^m (-p^{*\frac{1}{2}})^n = q^{-N}$ for any integers $m, n \in \mathbb{Z}$, $n \neq 0$. On any surface \mathcal{S}_{mn} , the generators*

$$t_{mn}(z) = \text{tr} \left((g^{\frac{1}{2}} h g^{\frac{1}{2}})^{-m} L((-p^{*\frac{1}{2}})^n z) (g^{\frac{1}{2}} h g^{\frac{1}{2}})^{-n} L(z)^{-1} \right)$$

realize an exchange algebra with the generators $L(w)$ of $\mathcal{A}_{q,p}(\widehat{gl}(N)_c)$:

$$t_{mn}(z) L(w) = \frac{\mathcal{F}_{-m}(z/w)}{\mathcal{F}_n^*(z/w)} L(w) t_{mn}(z)$$

where

$$\mathcal{F}_a(x) = \begin{cases} \prod_{k=0}^{a-1} \mathcal{U}((-p^{\frac{1}{2}})^k x) & \text{for } a > 0 \\ \prod_{k=1}^{|a|} \mathcal{U}((-p^{\frac{1}{2}})^{-k} x)^{-1} & \text{for } a < 0 \end{cases}$$

and $\mathcal{F}_a^*(x) = \mathcal{F}_a(x)|_{p \rightarrow p^*}$. One sets $\mathcal{F}_0(x) = 1$.

Remark 1 : the previous results (AFRS 97-99) correspond to the case $n=-1$, $m = kN - 1$, k integer, where the contribution of the automorphism $M = g^{1/2} h g^{1/2}$ essentially vanish due to $M^N = 1$, the balance is actually reabsorbed in the redefinition $L \rightarrow L_-$.

Remark 2 : when $n = -m = \pm 1$ closure relation yields $c = \pm N$ and leads directly to extended center, i.e. commutation of $t_{nm}(w)$ with $L(z)$. $c=-N$ known in affine algebras. $c=N$ not yet known ; seems specific to elliptic algebras.

Remark 3 : this extends to $n=-m$, n odd, $c=N/n$ and $-p^{1/2} = q^{-Nk/n}$, k not equal to $n-1$, coprime with n with Bezout coefficients resp. b, b' , $bk + b'n = 1$, and $b+1$ coprime with n : « weak extended center » (requires two conditions).

Implies commutation of generators t with themselves, seen later in abelianity section.

Remark 4 : Alternative closure relation constructed with automorphism $M \rightarrow g^a M$, $M' \rightarrow g^b M'$, denoting M and M' as resp. left and right automorphism in definition of t_{nm} . Uses antisymmetry of R matrix. Adds extra $(-)^{a+b}$ sign to closure relation and phase $e^{i(a+b)\pi/N}$ to exchange algebra.

Remark 5 : Liouville formula at $N=2$

$t_{0,2}(z) = \frac{1}{2} \text{tr} (L(q^{-2}z)L(z)^{-1})$ lies in the center of the elliptic quantum algebra.

No closure condition required ! Related to the q -det through

$$t_{0,2}(z) = \frac{q\text{-det } L(q^{-1}z)}{q\text{-det } L(z)}.$$

Remark 6 : Alternative construction :

Theorem 3.3 *On the surface \mathcal{S}_{mn} , the generators*

$$t_{-n,-m}^*(z) = \text{tr} \left((g^{\frac{1}{2}} h g^{\frac{1}{2}})^n L((-p^{\frac{1}{2}})^{-m} z)^{-1} (g^{\frac{1}{2}} h g^{\frac{1}{2}})^m L(z) \right)$$

realize an exchange algebra with the generators $L(w)$ of $\mathcal{A}_{q,p}(\widehat{gl}(N)_c)$:

$$t_{-n,-m}^*(z) L(w) = \frac{\mathcal{F}_n^*(z/w)}{\mathcal{F}_{-m}(z/w)} L(w) t_{-n,-m}^*(z)$$

From which it is immediate to get :

Proposition 3.4 *On the surface \mathcal{S}_{mn} , the generators $t_{mn}(z)t_{-n,-m}^*(z)$ lie in the center of $\mathcal{A}_{q,p}(\widehat{gl}(N)_c)$.*

In addition $t^*(z)$, $t^*(w)$ realize same exchange algebra as $t(z)$, $t(w)$.

Connection between t and t^* generators suggested by above results, seems to be $t^{-1} \sim t^*$.

However only possible to establish assuming closure relation **and $N=2$** as :

$$t_{-n,-m}^*(z) = \frac{q\text{-det } L(qz)}{q\text{-det } L((-p^{\frac{1}{2}})^{-m} z)} t_{mn}(qz)$$

Inverse structure function in fact realized through modular relation $U(z) \sim 1/U(qz)$. Quite peculiar to $N=2$ as degeneracy between $N/2$ and 2 .

$N > 2$?

Exchange relations for $t(z)$, $t(w)$ follow immediately :

Corollary 3.5 *On the surface \mathcal{S}_{mn} , the generators*

$$t_{mn}(z) = \text{tr} \left((g^{\frac{1}{2}} h g^{\frac{1}{2}})^{-m} L((-p^{*\frac{1}{2}})^n z) (g^{\frac{1}{2}} h g^{\frac{1}{2}})^{-n} L(z)^{-1} \right)$$

close a quadratic subalgebra in $\mathcal{A}_{q,p}(\widehat{gl}(N)_c)$:

$$t_{mn}(z) t_{mn}(w) = \mathcal{Y}_{mn}(z/w) t_{mn}(w) t_{mn}(z)$$

with

$$\mathcal{Y}_{mn}(x) = \frac{\mathcal{F}_n^*(x) \mathcal{F}_{-n}^*(x)}{\mathcal{F}_m(x) \mathcal{F}_{-m}(x)} = \frac{\prod_{k=1}^{|m|} \mathcal{U}((-p^{\frac{1}{2}})^{-k} x) \prod_{k'=0}^{|n|-1} \mathcal{U}((-p^{*\frac{1}{2}})^{k'} x)}{\prod_{k=0}^{|m|-1} \mathcal{U}((-p^{\frac{1}{2}})^k x) \prod_{k'=1}^{|n|} \mathcal{U}((-p^{*\frac{1}{2}})^{-k'} x)}$$

Important point : For N=2 an explicit factorization is available :

$$\mathcal{Y}_{mn}(x) = \frac{g_{mn}(x^2)}{g_{mn}(x^{-2})} \quad \text{with} \quad g_{mn}(z) = g_{corr}(z) \left(\prod_{k=1}^{|m|-1} g^{(k)}(z) \right)^2 \left(\prod_{k=1}^{|n|-1} g^{*(k)}(z) \right)^{-2}$$

where

$$g^{(k)}(z) = \exp \left(\sum_{\ell=1}^{\infty} \frac{(1 - p^{-k\ell})(1 - (p^k q^2)^\ell)}{1 + q^{2\ell}} \frac{z^\ell}{\ell} \right)$$

$$g^{*(k)}(z) = g^{(k)}(z) \Big|_{p \rightarrow p^*}$$

$$g_{corr}(z) = g^{(|m|)}(z) (g^{*(|n|)}(z))^{-1}$$

Note the overall square !

Hence scaling limit can be defined :

Defining the scaling limit as $p = 1 + \varepsilon$ and $q = 1 + \eta\varepsilon$ with $\varepsilon \rightarrow 0$, we observe that

$$g^{(k)}(z) = 1 - \varepsilon^2 \frac{k(k+2\eta)}{2} \frac{z}{(1-z)^2} + o(\varepsilon^2)$$

$$g^{*(k)}(z) = 1 - \varepsilon^2 \frac{k(1-2\eta c)(k(1-2\eta c)+2\eta)}{2} \frac{z}{(1-z)^2} + o(\varepsilon^2)$$

leading to

$$g_{mn}(z) = 1 - \varepsilon^2 \left\{ \beta_m - (1 - 2\eta c)^2 \beta_n - 2\eta^2 c(1 - 2\eta c)n(n-2) \right\} \frac{z}{(1-z)^2} + o(\varepsilon^2)$$

where

$$\beta_\ell = \frac{|\ell|(|\ell|-1)(2|\ell|-1)}{6} + \eta\ell(\ell-2).$$

Also true for $N > 2$, with no full square expression for g_{mn} and more complicated factor for scaling limit. But scaling limit nevertheless gives **exact structure function for quantum Virasoro algebra**. Hence quadratic algebras characterized as **DEFORMED VIRASORO ALGEBRAS**. But DVA of **Shiraishi Kubo Awata Odake** **elusive due to square!** And delicate issue with central extension to get exact scaling limit (**BUT!** **1st term of g anyway fully controls centerless part of limit**).

Additional exchange relations on surface \mathcal{S}_{mn}

$$t_{mn}(z) t_{rs}(w) = \frac{\mathcal{F}_n^*(z/w)}{\mathcal{F}_{-m}(z/w)} \frac{\mathcal{F}_{-m}((-p^{*\frac{1}{2}})^s z/w)}{\mathcal{F}_n^*((-p^{*\frac{1}{2}})^s z/w)} t_{rs}(w) t_{mn}(z)$$

Yields coefficients Y if $s=n$, no requirement on r .

Abelianity relations now follow from examining possible cancellations of terms in above products. Essentially « vertical » cancellation between m -products, and separately n -products, except when $m = \pm n$ where cross-cancellations may occur as already seen on example of extended centre and weak extended center.

Remark 6 : First immediate abelianity condition : **for t_{nm} at $n=0, |m| > 1$ and t_{nm}^* at $n=0, |m| > 1$.**
Not extended center : Only on-shell (closure relation obeyed). Requires only closure. Liouville formula = stronger statement (off-shell).

II-2 THE CRITICALITY RELATIONS

Particular case already seen $m=-n$, more developed here.

General case :

- for $|m|, |n| > 1$:

$$c = \frac{N}{nm}(\lambda'm - \lambda n), \quad -p^{\frac{1}{2}} = q^{-N\lambda/m}, \quad -p^{*\frac{1}{2}} = q^{-N\lambda'/n} \quad (3.19)$$

where $\lambda, \lambda' \in \mathbb{Z} \setminus \{0\}$ and $\lambda + \lambda' = 1$.

- for $|n| = 1, |m| > 1$:

$$c = Nn(1 - \lambda(m+n)), \quad -p^{\frac{1}{2}} = q^{-N\lambda}, \quad -p^{*\frac{1}{2}} = q^{-Nn(1-\lambda m)} \quad (3.20)$$

where $\lambda \in \mathbb{Z}/2$ or $\lambda \in \mathbb{Z}/u$, u being any divisor of m or $m+n$.

- for $|m| = 1, |n| > 1$:

$$c = Nm(\lambda'(n+m) - 1), \quad -p^{\frac{1}{2}} = q^{-Nm(1-\lambda'n)}, \quad -p^{*\frac{1}{2}} = q^{-N\lambda'} \quad (3.21)$$

where $\lambda' \in \mathbb{Z}/2$ or $\lambda' \in \mathbb{Z}/u'$, u' being any divisor of n or $n+m$.

- for $m = n = \pm 1$, formulas (3.20)–(3.21) also hold with $\lambda, \lambda' \in \mathbb{Z}/2$ and $\lambda + \lambda' = 1$.

- for $m+n=0$ with $n > 0$ and odd:

$$c = \frac{N}{n}, \quad -p^{\frac{1}{2}} = q^{-\frac{n-1}{2n}N}, \quad -p^{*\frac{1}{2}} = q^{-\frac{n+1}{2n}N}. \quad (3.22)$$

Remark 7 : Similar result holds for generators in Remark 4 up to overall sign. Same result holds for t^* generators since same exchange algebra.

II-3 THE POISSON STRUCTURES

Obtained in usual way. Fix closure condition \mathcal{S}_{mn} . Then set quasi-abelianity relation as :

$$p^{1-\epsilon} = q^{\alpha N \ell}$$

and expand as

$$\{t(z), t(w)\}_\ell = \lim_{\epsilon \rightarrow 0} \frac{1}{\epsilon} (t(z)t(w) - t(w)t(z)) \quad \text{One gets :}$$

$$\{t(z), t(w)\}_\ell = f_\ell(z/w) t(z)t(w) \quad f_\ell(x) = 2N\ell(\ln q)(2I(x) - I(qx) - I(q^{-1}x) - (x \leftrightarrow x^{-1}))$$

$I(x)$ is given by the following expressions depending on the different cases of proposition

- for $|m|, |n| > 1$ ($\alpha = 2/m$):

$$I(x) = \sum_{s=-1}^{\infty} \left(w \frac{x^2 q^{2Nsw/m}}{1 - x^2 q^{2Nsw/m}} + w' \frac{x^2 q^{2Nsw'/n}}{1 - x^2 q^{2Nsw'/n}} \right) + \frac{1}{2} (w + w') \frac{x^2}{(1 - x^2)}$$

where $w = \gcd(\ell, m)$ and $w' = \gcd(\ell', n)$ (w and w' being taken positive).

- for $|n| = 1, |m| > 1$ ($\alpha = 1$) and ℓ even:

$$I(x) = \frac{1}{2} m(m+1) \left[\sum_{s=1}^{\infty} \frac{x^2 q^{2Ns}}{1 - x^2 q^{2Ns}} + \frac{1}{2} \frac{x^2}{1 - x^2} \right] \quad (4.15)$$

while for ℓ odd

$$I(x) = \left[\frac{m}{2} \right] \left(\left[\frac{m}{2} \right] + 1 \right) \left(\sum_{s=1}^{\infty} \frac{x^2 q^{2Ns}}{1 - x^2 q^{2Ns}} + \frac{1}{2} \frac{x^2}{1 - x^2} \right) + \left[\frac{m+1}{2} \right]^2 \sum_{s=0}^{\infty} \frac{x^2 q^{N(2s+1)}}{1 - x^2 q^{N(2s+1)}} \quad (4.16)$$

$\lfloor x \rfloor$ is the floor function (integer part) of x .

- for $|n| = 1, |m| > 1$ ($\alpha = 2/u$ where u is any divisor of k_{sup} or $k_{\text{sup}} + 1$ and g is defined as k_{sup}/u or $(k_{\text{sup}} + 1)/u$ respectively):

$$I(x) = g \left[\eta \sum_{s=1}^{\infty} \left(\frac{x^2 q^{2Ns}}{1 - x^2 q^{2Ns}} + gw \frac{x^2 q^{2Nsw/u}}{1 - x^2 q^{2Nsw/u}} \right) + \frac{1}{2} (gw + \eta) \frac{x^2}{1 - x^2} \right] \quad (4.17)$$

where $\eta = 1$ when $gu = k_{\text{sup}}$ and $\eta = -1$ when $gu = k_{\text{sup}} + 1$, $w = \gcd(\ell, u)$, see the definition of k_{sup} after formula (C.14).

- for $|m| = 1, |n| > 1$ ($\alpha = 1$): the formulas are analogous to (4.15)-(4.16) (ℓ being replaced by ℓ' and m by n), up to a non relevant normalization factor that can be absorbed by a redefinition of ϵ .

II-4 THE DEFORMED VIRASORO ALGEBRAS

Classical limit yields classical DVA in e.g. (4.15), (4.16), **up to normalizations**

Directly realize quantum DVA exchange algebra on closure surface ?

Not true, as mentioned earlier

only S^2 for $m = -2$! Worse for higher values of m .

$\hat{\mathbf{A}}_{q,p,c}(\widehat{\mathfrak{gl}}_N)$: modified elliptic quantum algebra. Already defined above :

$$\begin{aligned} \widehat{R}_{12}(z/w) L_1^{\pm}(z) L_2^{\pm}(w) &= L_2^{\pm}(w) L_1^{\pm}(z) \widehat{R}_{12}^*(z/w), \\ \widehat{R}_{12}(q^{\frac{\epsilon}{2}} z/w) L_1^+(z) L_2^-(w) &= L_2^-(w) L_1^+(z) \widehat{R}_{12}^*(q^{-\frac{\epsilon}{2}} z/w). \end{aligned} \quad \text{this time with unitary R matrix. No relation now}$$

assumed between L^+ and L^- . Seen in [Foda et al. \(1995\)](#). as vertex operator construction of original elliptic algebra.

HENCE NOT EQUIVALENT TO $\mathbf{A}_{\text{qp}}(\mathfrak{gl}(N))_c$. « vertex-operator algebra » from its construction.

Introduce now :

$$t_{mn}(z) = \text{tr} \left((g^{\frac{1}{2}} h g^{\frac{1}{2}})^{-m+1} \mathcal{L}^+ \left((-p^{\frac{1}{2}})^{n+1} q^{-nc-c/2} z \right) (g^{\frac{1}{2}} h g^{\frac{1}{2}})^{-n-1} \mathcal{L}^- (z)^{-1} \right)$$

$$t_{mn}^*(z) = \text{tr} \left((g^{\frac{1}{2}} h g^{\frac{1}{2}})^{n-1} \mathcal{L}^- \left((-p^{\frac{1}{2}})^{-m-1} q^{\frac{c}{2}} z \right)^{-1} (g^{\frac{1}{2}} h g^{\frac{1}{2}})^{m+1} \mathcal{L}^+ (z) \right)$$

Same closure conditions, conditions since exchange relations of generating matrices only differ by scalar factor. Abelianity conditions certainly modified.

For N=2, on closure surface n=-1, m=2: exchange algebra realizes exactly quantum DVA

Recall : p,q (SKAO) = q², p⁻¹ (alternative elliptic algebra parameters, c(p,q) expressed by closure relation)

III OPEN ISSUES

RE : DVA

Interpretation of new algebra ; in particular Hopf structure ?

Connection with MacDonalD polynomials ? Seen in vertex operator construction of SKAO DVA, how about other DVA's ? in particular those obtained from $A_{qp}(\mathfrak{gl}(N))_c$ with « squared » structure functions.

Central extensions ? Not directly gotten in abstract construction from algebra relations. Explicit realizations and/or resolution of cocycle condition (partially done in '98). Important e.g. for consistent linear scaling limit

RE : OTHER STRUCTURES

N>2 leads to possible constructions of q-W_{N-1} algebras. (proposed in e.g. [Feigin-Frenkel, AKOS](#)) Requires to define consistent higher-spin generators as (?) traces of multilinear objects. First attempts in '98.

Dynamical algebras $B_{qp\lambda}(\mathfrak{gl}(N))_c$: promising ?